

1+1次元有限密度QCDの ハミルトン形式 による数値解析

西村健太郎 (広島大学)

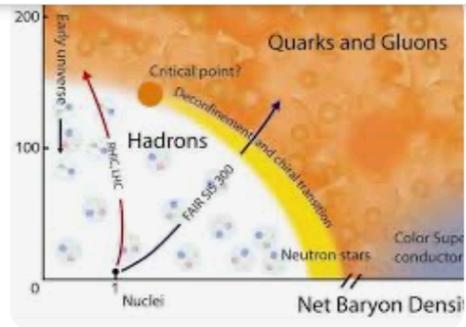
共同研究者

早田智也 (慶應義塾大学) 日高義将 (京都大学)

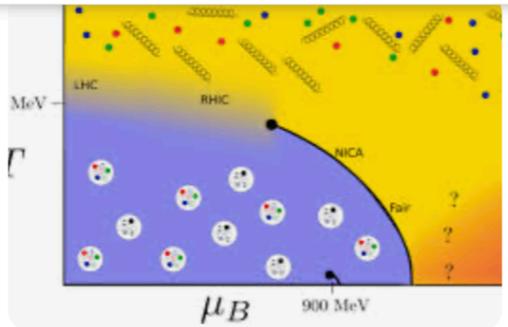
熱場の量子論 2024/9/11

[JHEP 07 \(2024\) 106](#)

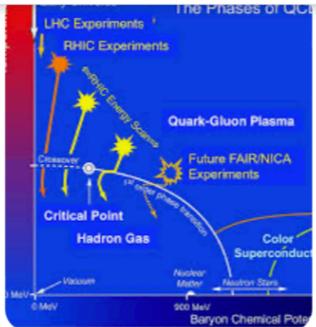
ResearchGate A possible sketch of the QCD pha...



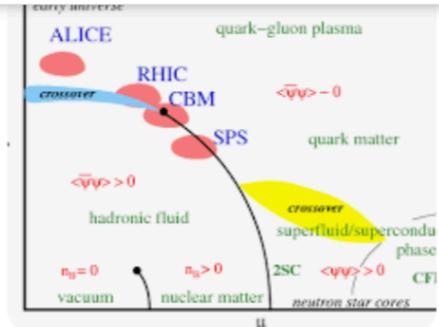
SpringerLink Overview of the QCD phase diagram...



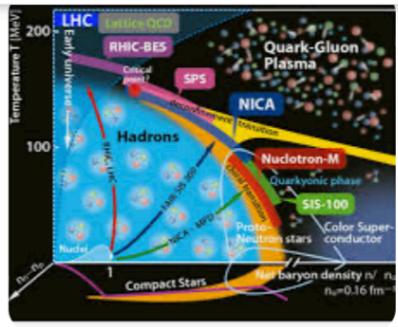
CERN Document Server Probing the QCD pha...



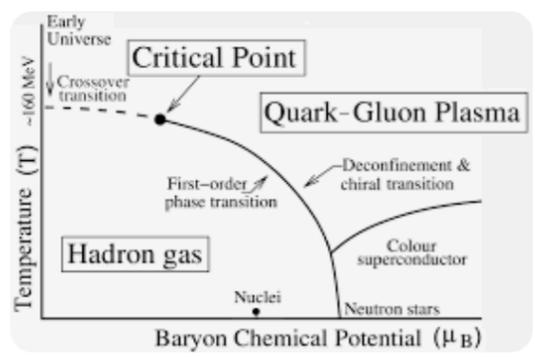
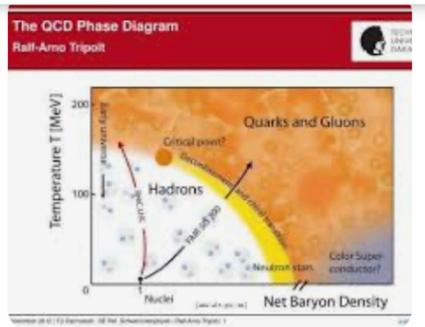
ResearchGate 2: Schematic view of QCD pha...



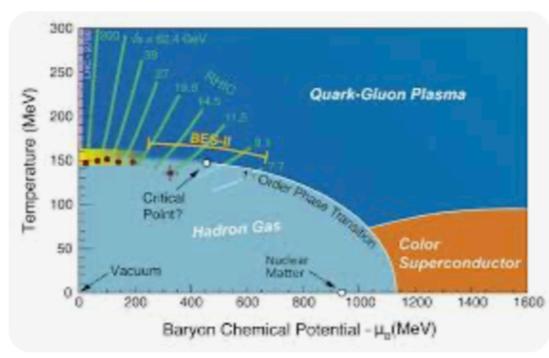
ResearchGate Color Online) A schematic o...



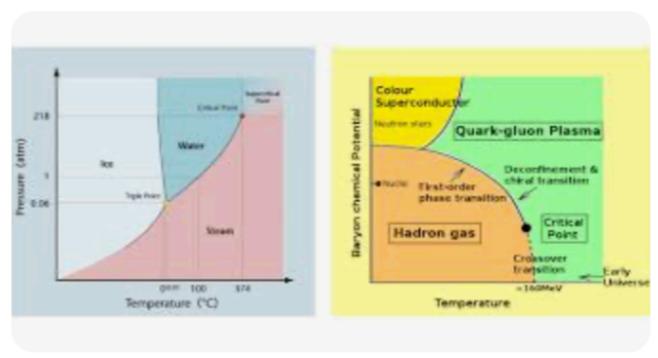
YUMPU The QCD Phase Diagram - T...



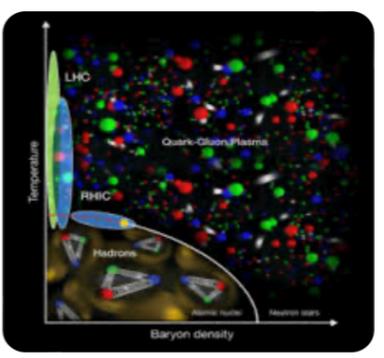
Gauss Centre for Supercomputing



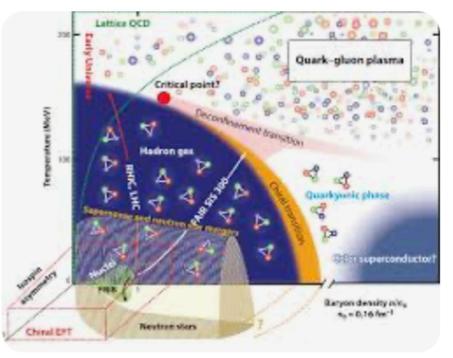
CERN Courier



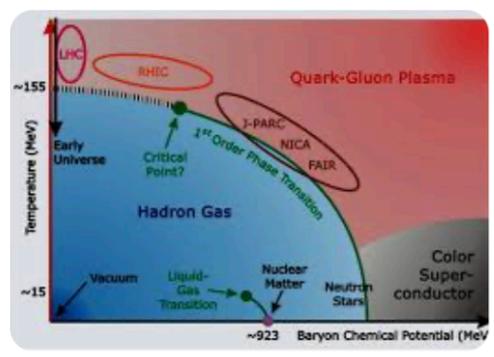
Particles and friends - WordPress.com



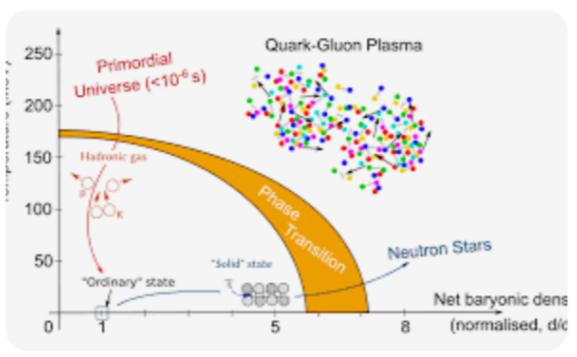
GdR QCD



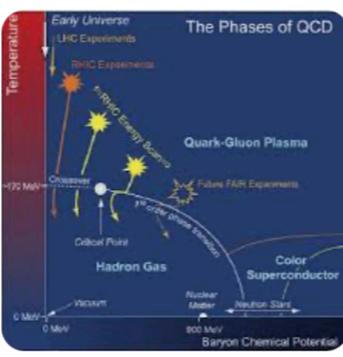
ResearchGate



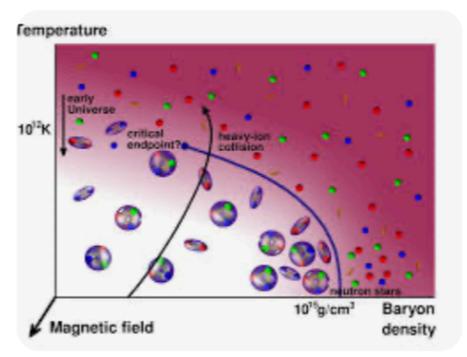
www.gsi.de GSI - Hot and Dense QCD Matter



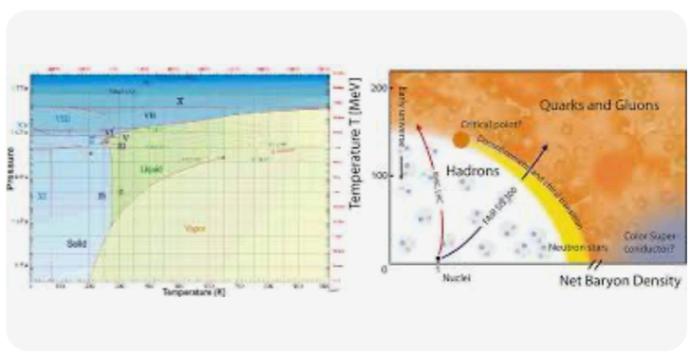
CERN Document Server



Semantic Scholar



IOPscience



Gauss Centre for Supercomputing

Sign problem

- Monte Carlo simulation can not be applied to QCD at finite density.
 - Three-color QCD with isospin chemical potential:
[Kogut and Sinclair \(2002\)](#).....
 - QCD at finite baryon density with $N_c=2$ and $N_f=(\text{even number})$:
[Hands, Kim and Skullerud \(2006\)](#), [Iida and Itou \(2022\)](#),
- Finite density QCD
 - Interior of the Neutron stars [See talk by T. Kojo](#)
 - Compressed matter by the stopping in low-energy heavy ion collision experiment
- **It is important to theoretically investigate QCD at finite density!**

Hamiltonian formalism

- Lattice Hamiltonian on a spatial lattice, and time is continuous.
[Kogut and Susskind \(1975\)](#)
- Minimization of the free energy without Monte Carlo → No sign problem
 - We use good variational ansatz and efficient algorithm for searching the ground state.
- There are many applications of Hamiltonian formalism.
 - SU(2) or SU(3) lattice gauge theory in (1+1) dim: [Kuhn et.al., \(2015\)](#); [Silvi et.al., \(2017\)](#); [Banuls et.al., \(2017\)](#); [Sala et.al., \(2018\)](#)
 - Schwinger model: [Silvi et.al., \(2017\)](#); [Rigobello et.al., \(2023\)](#)
- Numerical calculation of thermodynamic quantities 1+1 dim QCD in finite baryon chemical potential.
 - To our knowledge, comprehensive study of dense QCD₂ has not been done yet.

Matrix Product State and Density Matrix Renormalization Group

- **Matrix Product State (MPS):** [Perez-Garcia, et.al., \(2007\)](#) ; [McCulloch, et.al., \(2007\)](#)

- This is frequently used for quantum spin systems in 1+1 dimension

- **Density Matrix Renormalization Group (DMRG)** [White \(1992\)](#)

- Efficient algorithm for searching the ground state represented by MPS

- **iTensor = A wonderful package that performs MPS and DMRG!**

<https://itensor.github.io/ITensors.jl/stable/ITensorType.html>

Also see reference by [Fishman, White and Stoudenmire \(2022\)](#)

Hamiltonian Formalism for Lattice gauge theory in (1+1)D

- **Kogut-Susskind Hamiltonian:** $H = H_E + H_{\text{hop}} + H_m$ [Kogut and Susskind \(1975\)](#)

$$H_E = \frac{ag_0^2}{2} \sum_{n=1}^{N-1} E_i^2(n),$$

$$H_{\text{hop}} = \frac{1}{2a} \sum_{n=1}^{N-1} (\chi^\dagger(n+1)U(n)\chi(n) + \chi^\dagger(n)U^\dagger(n)\chi(n+1)),$$

$$H_m = m_0 \sum_{n=1}^N (-1)^n \chi^\dagger(n)\chi(n)$$

- **Quark (Baryon) number:** $N_q = \sum_{n=1}^N \chi^\dagger(n)\chi(n)$

- **Full Hamiltonian:** $H_{\text{tot}} \equiv H_E + H_{\text{hop}} + H_m - \mu_{0B}N_B$

Dimensionless Hamiltonian

- **Full Hamiltonian w/ n_B normalized by g_0 :** $\tilde{H}_{\text{tot}} \equiv H_{\text{tot}}/g_0$

$$H_E/g_0 = J \sum_{n=1}^{N-1} E_i^2(n) \equiv \tilde{H}_E,$$

$$H_{\text{hop}}/g_0 = \epsilon \sum_{n=1}^{N-1} (\chi^\dagger(n+1)U(n)\chi(n) + \chi^\dagger(n)U^\dagger(n)\chi(n+1)) \equiv \tilde{H}_{\text{hop}},$$

$$H_m/g_0 = m \sum_{n=1}^N (-1)^n \chi^\dagger(n)\chi(n) \equiv \tilde{H}_m,$$

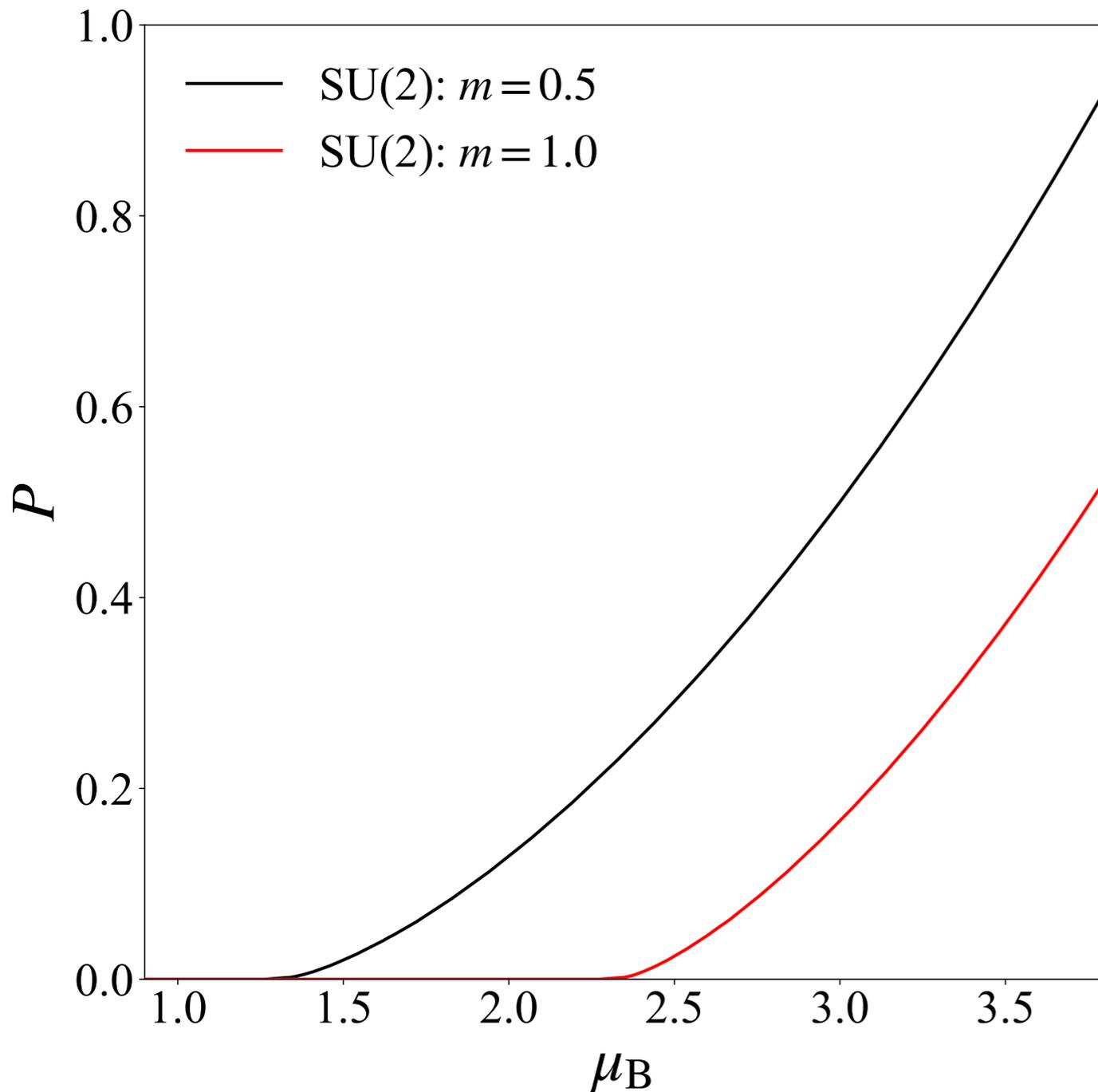
$$\mu_{0B}N_B/g_0 = \mu_B \frac{1}{N_c} \sum_{n=1}^N \chi^\dagger(n)\chi(n) \equiv \mu_B N_B.$$

- **Parameters:** $J = ag_0/2$, $\epsilon = 1/(2ag_0)$, $m = m_0/g_0$, $\mu = \mu_0/g_0$

- Since I can solve Gauss law constraint in open boundary condition, $U(n)$ can be eliminated and represented by Staggered fermions.
- After Wigner-Jordan transformation, our Hamiltonian is represented by spins, or Pauli matrices.

SU(2) results

Pressure



- $N=160$, $N_c=2$, and $w=2$.

- **Grand potential and Pressure:**

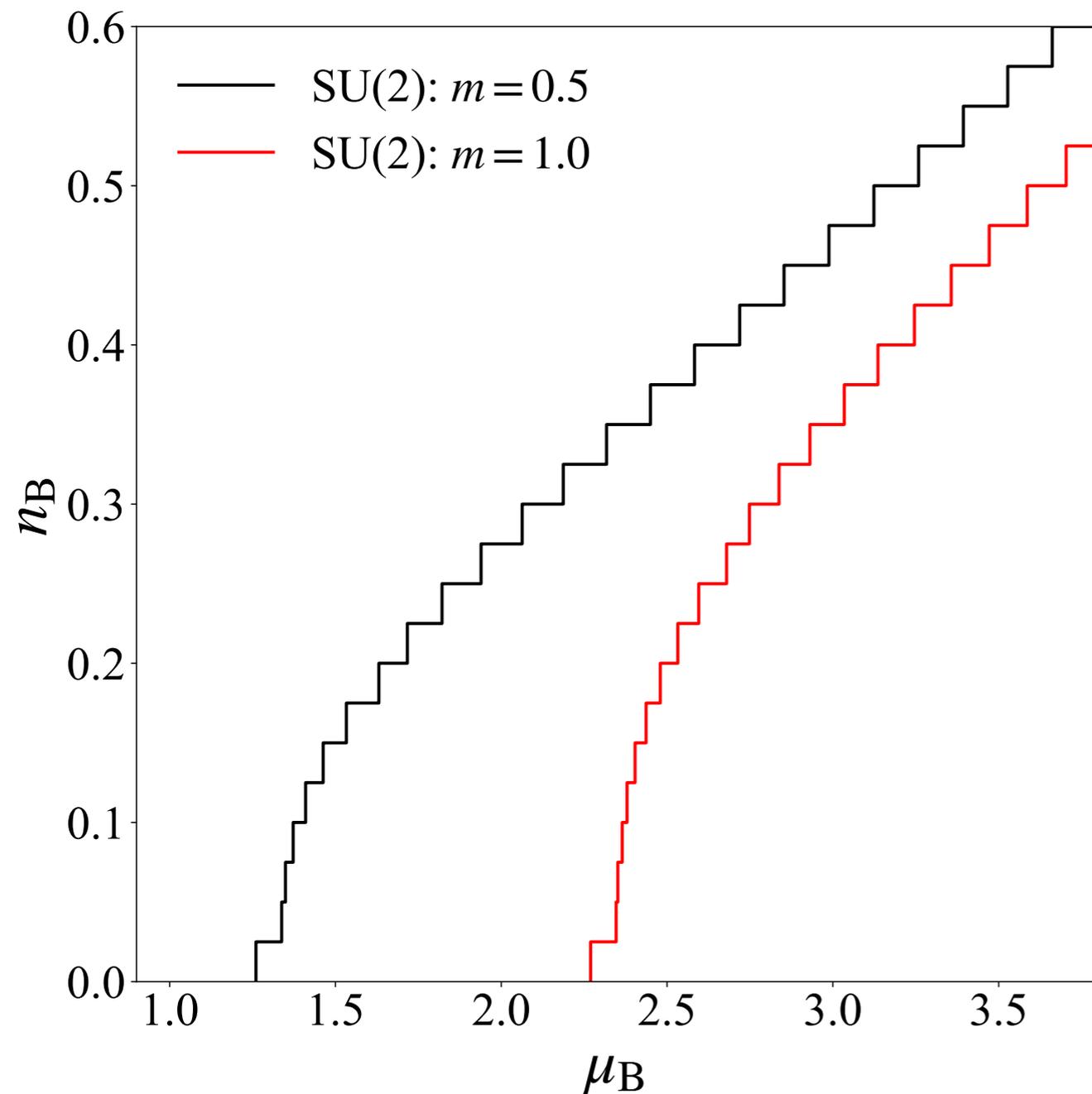
$$J = \langle \text{GS} | (H - \mu_B N_B) | \text{GS} \rangle$$

$$P = -J/V$$

- **First derivative = baryon number:**

$$n_B = \frac{dP(\mu_B)}{d\mu_B} = \frac{1}{V} \langle N_B \rangle$$

Baryon number



- **Threshold value:**

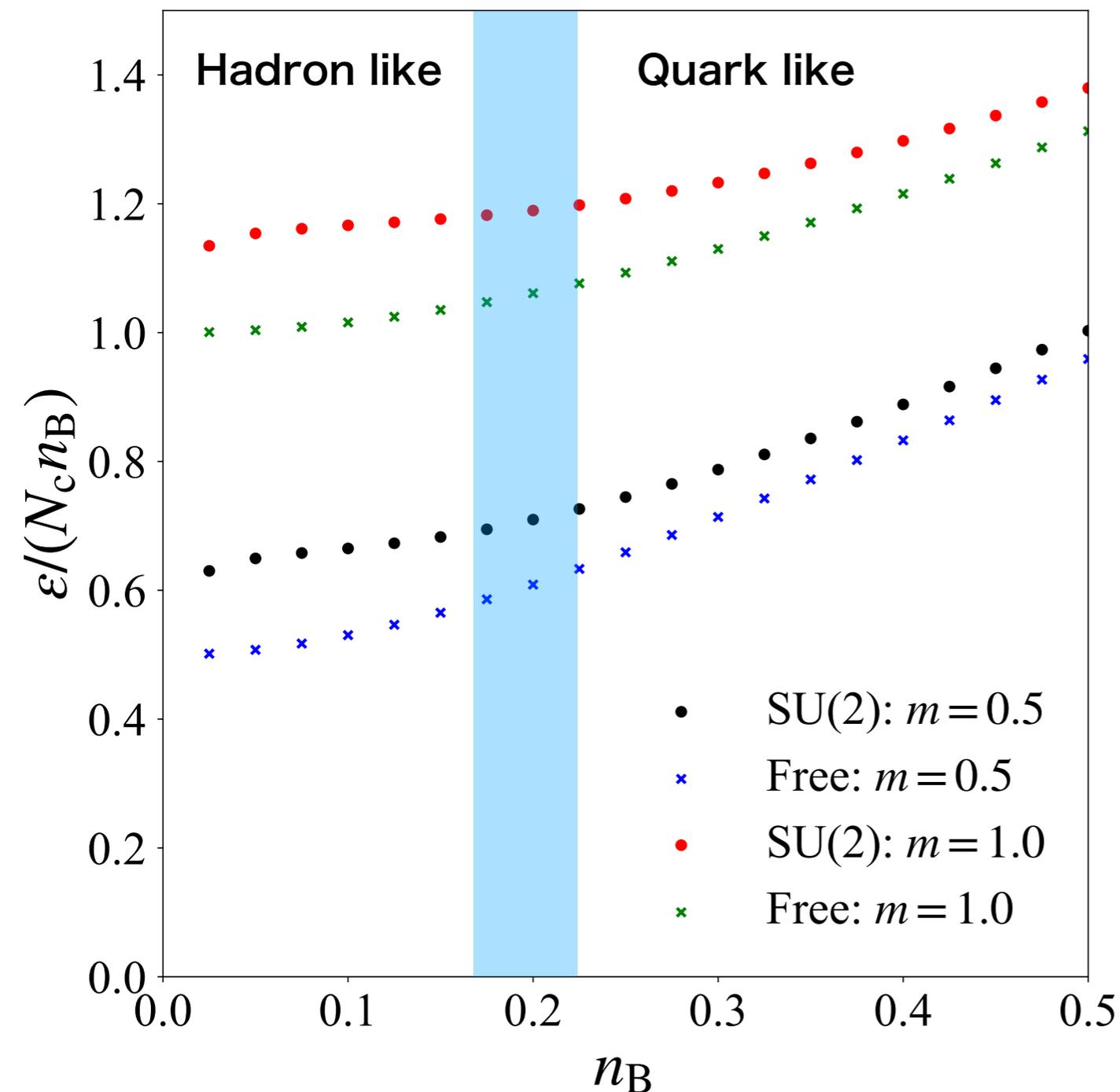
$$\mu_c = \begin{cases} 1.26 > 2 \times 0.5 & (m = 0.5) \\ 2.27 > 2 \times 1.0 & (m = 1.0) \end{cases}$$

- **At high density $g_0/\mu \ll 1$ and $m/\mu \ll 1$, the contribution from g_0 and m can be negligible \rightarrow Free-quark!**

$$n_B = N_c \mu / \pi = \mu_B / \pi$$

- **How does it behave at low density?**

From hadron to quark

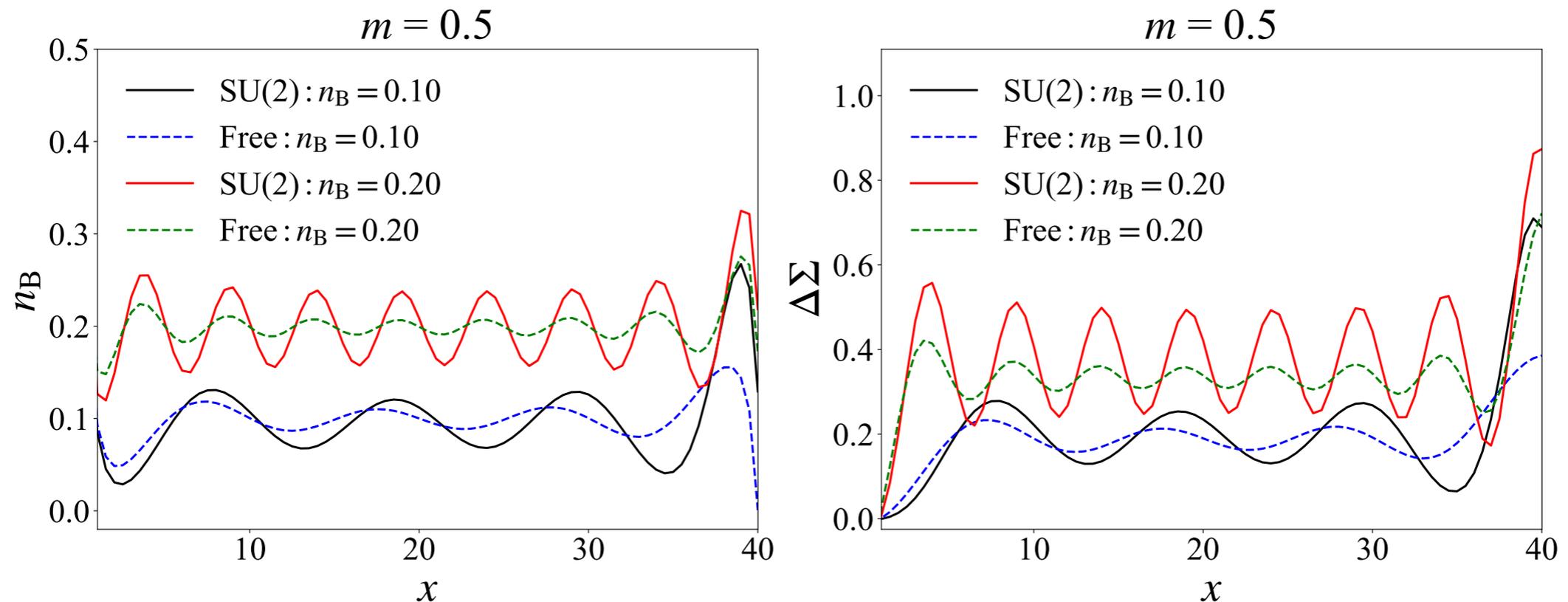


- Energy per one quark
- At high density, ϵ/n_q behaves like the free theory.
- At low density, ϵ/n_q is larger due to the confining energy.
- The behavior changes to free theory around $n_B \sim 0.2$.
- If interaction can be negligible, the baryons will degenerate to the lowest energy state.

$$\epsilon/(N_c n_B) = \text{const}$$

Inhomogeneous phase ($m=0.5$)

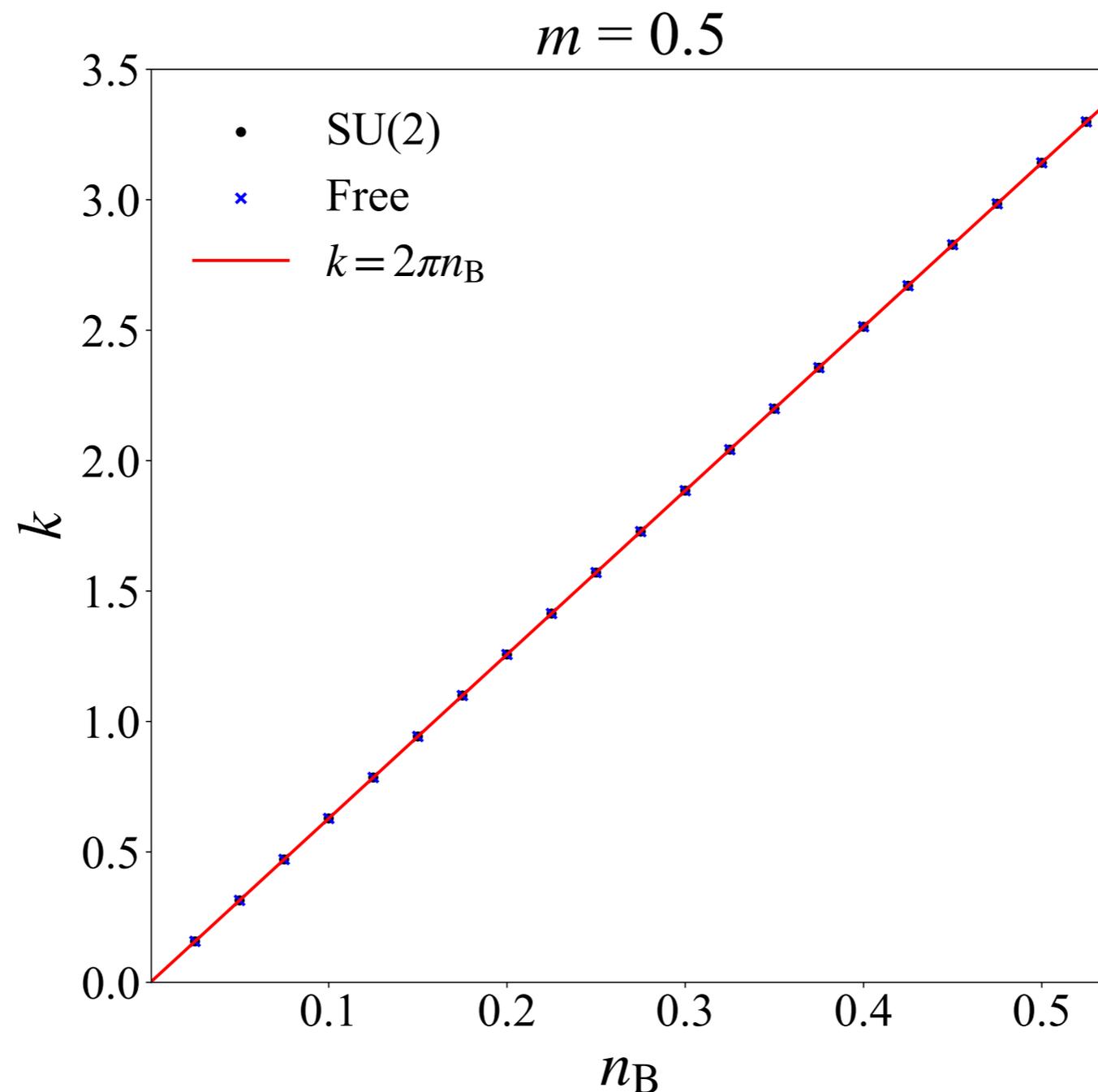
- Can the inhomogeneous state realize beyond mean field approximation?
 - Gross-Neveu, chiral Gross-Neveu, and QCD in (1+1) dim w/ Mean field approximation:
[Schon and Thies, \(2000\)](#); [Thies and Urlichs, \(2003\)](#); [Kojo, \(2012\)](#)
- SSB of continuous symmetry is prohibited.
 - The spatial modulation is induced by the open boundary condition.



- In free theory, the modulations vanish proportionally to $1/L$.

Wave number (period)

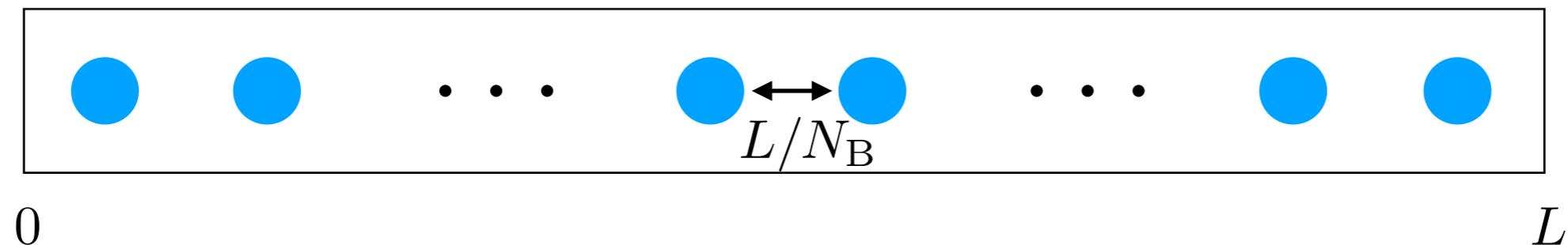
- Fourier transform on $n_B(x)$ and wave number of the largest modulation.



Baryon crystal and chiral density wave

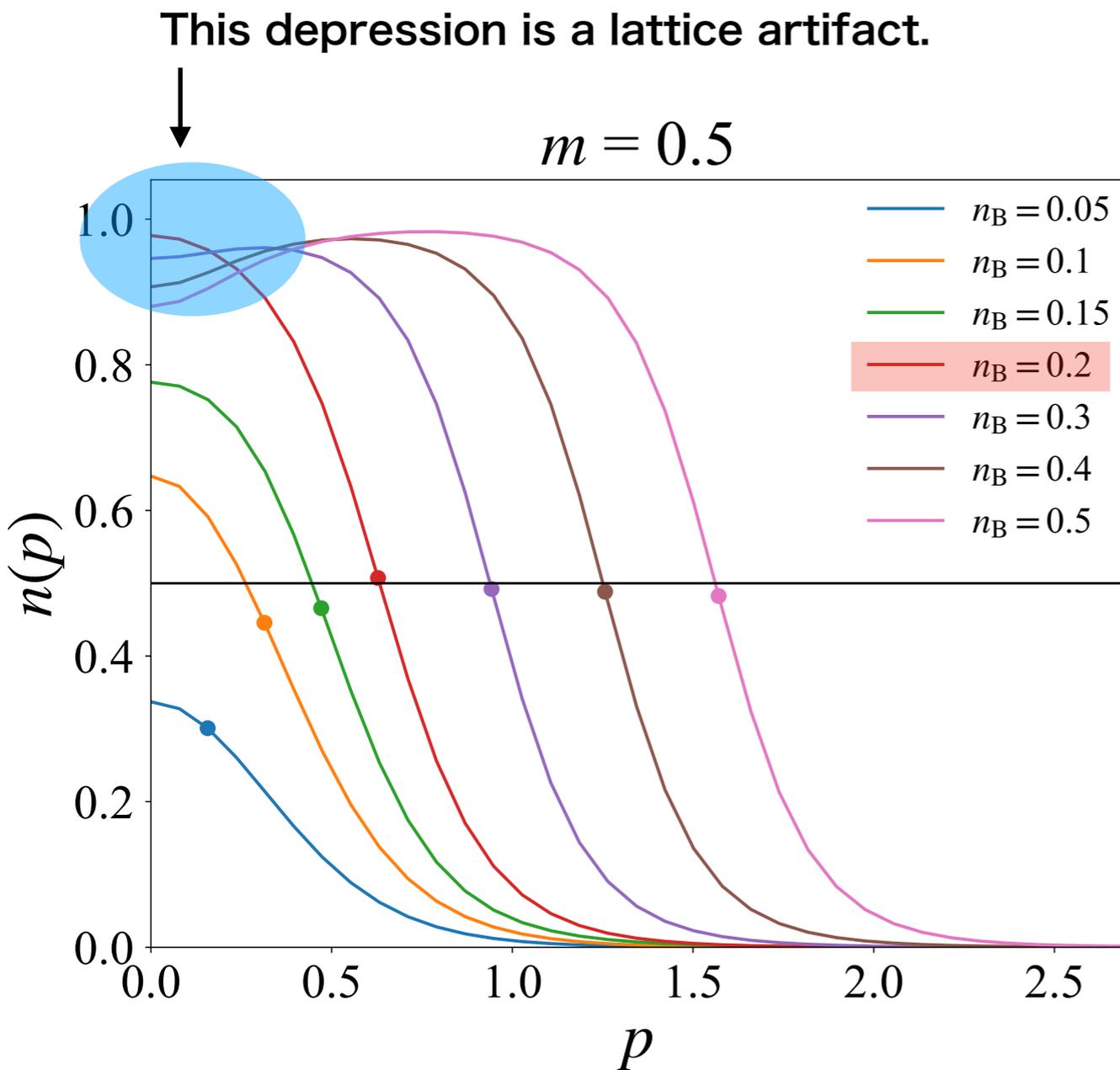
- If I assume repulsive force, the baryons are periodically aligned.

cf) Bosonization with large- N_c expansion [Kojo, \(2012\)](#)



- $k=2\pi / (\text{period})=2\pi / (L/N_B)=2\pi n_B$
- Due to the Peierls instability, the particle and hole form the condensate.
- Both of the particle and hole have p_F , then the total momentum is $2p_F$.
- $k=2p_F=2\pi n_B$ - In (1+1) dim, the fermi momentum is πn_B .

Distribution function



- $N=240, w=2.0$

- to reduce the finite volume effect

- As n_B increases, maximum of $n(p)$ increases and forms Fermi sea.

cf) Quarkyonic matter in (3+1) dim

BCS-BEC crossover of ultracold atom gases

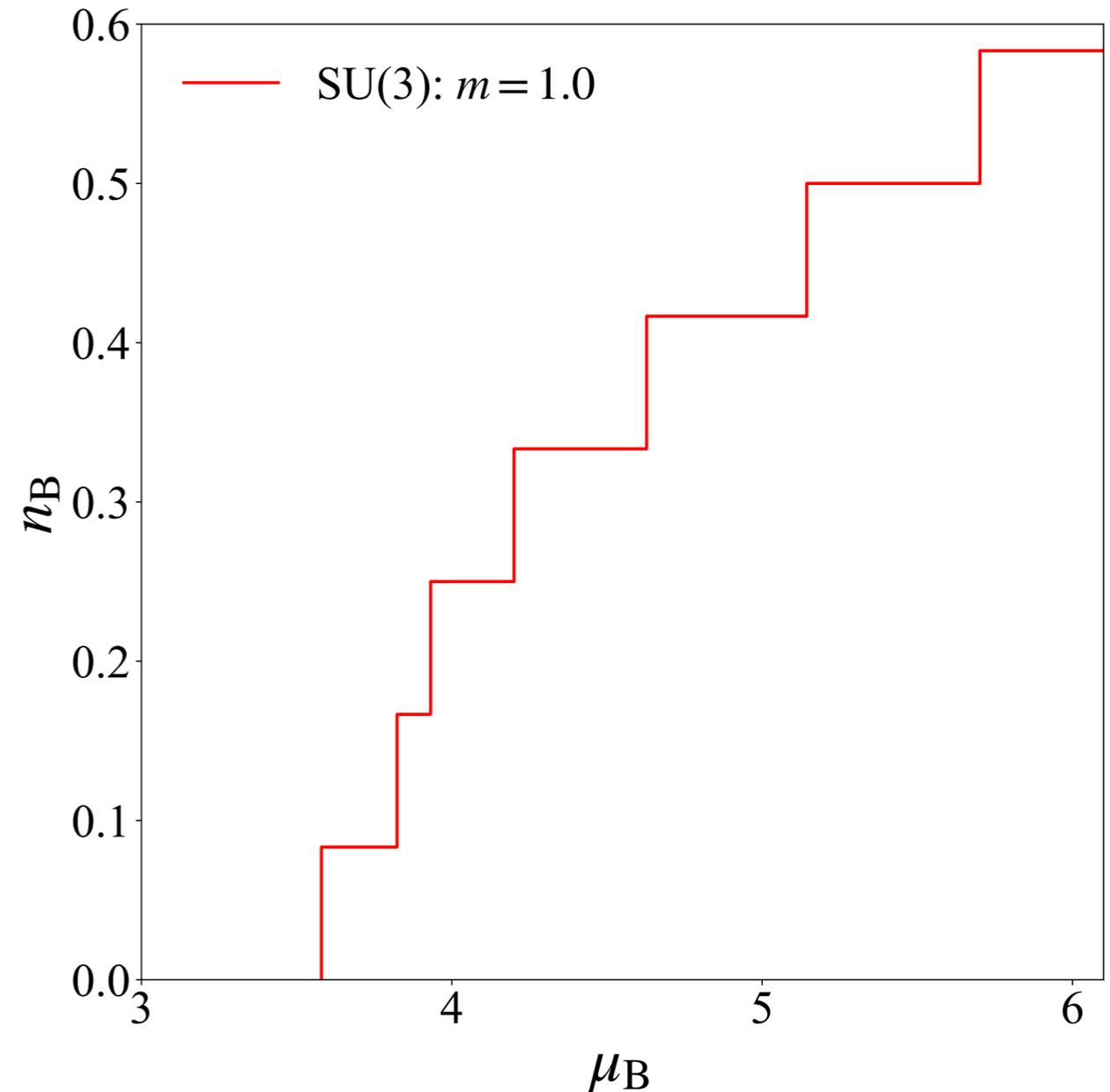
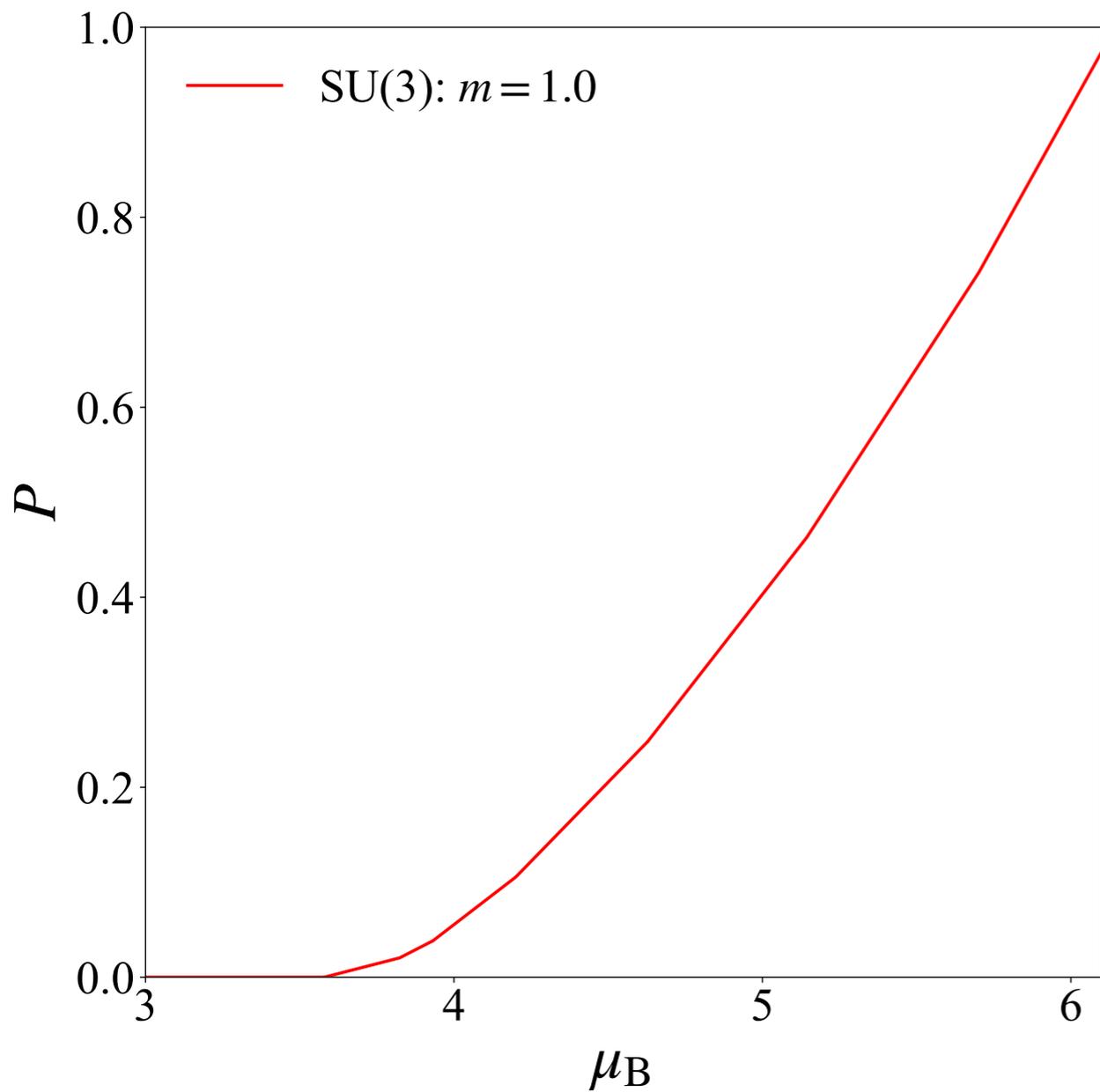
[Kojo, \(2012\)](#); [Astrakharchik et.al., \(2005\)](#); [Regal et.al., \(2005\)](#)

- Fermi sea forms near $n_B \sim 0.2$, which is consistent with behavior of ε/n_q
- Fermi surface is smooth, not sharp.
→ Instability of the Fermi surface

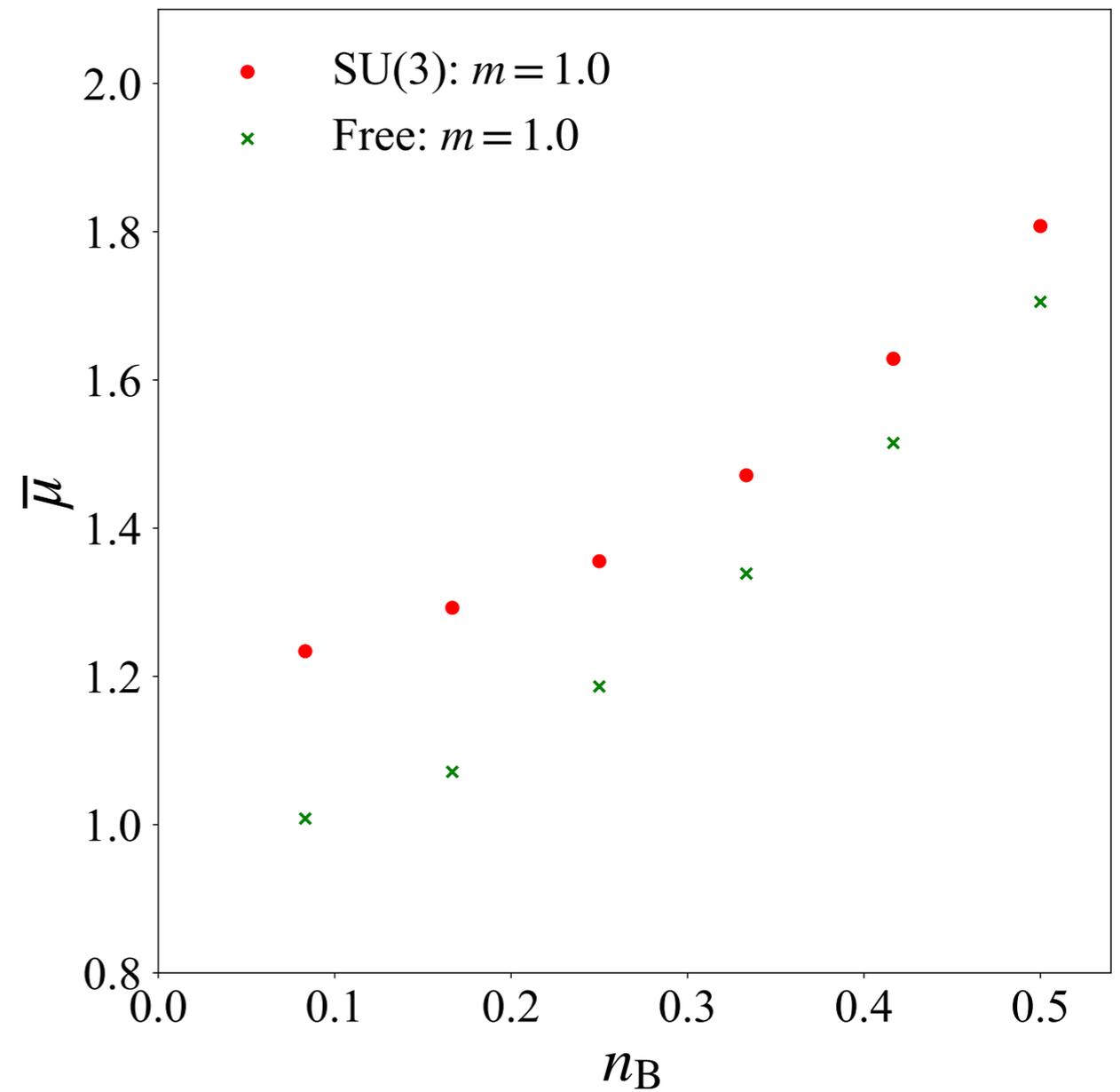
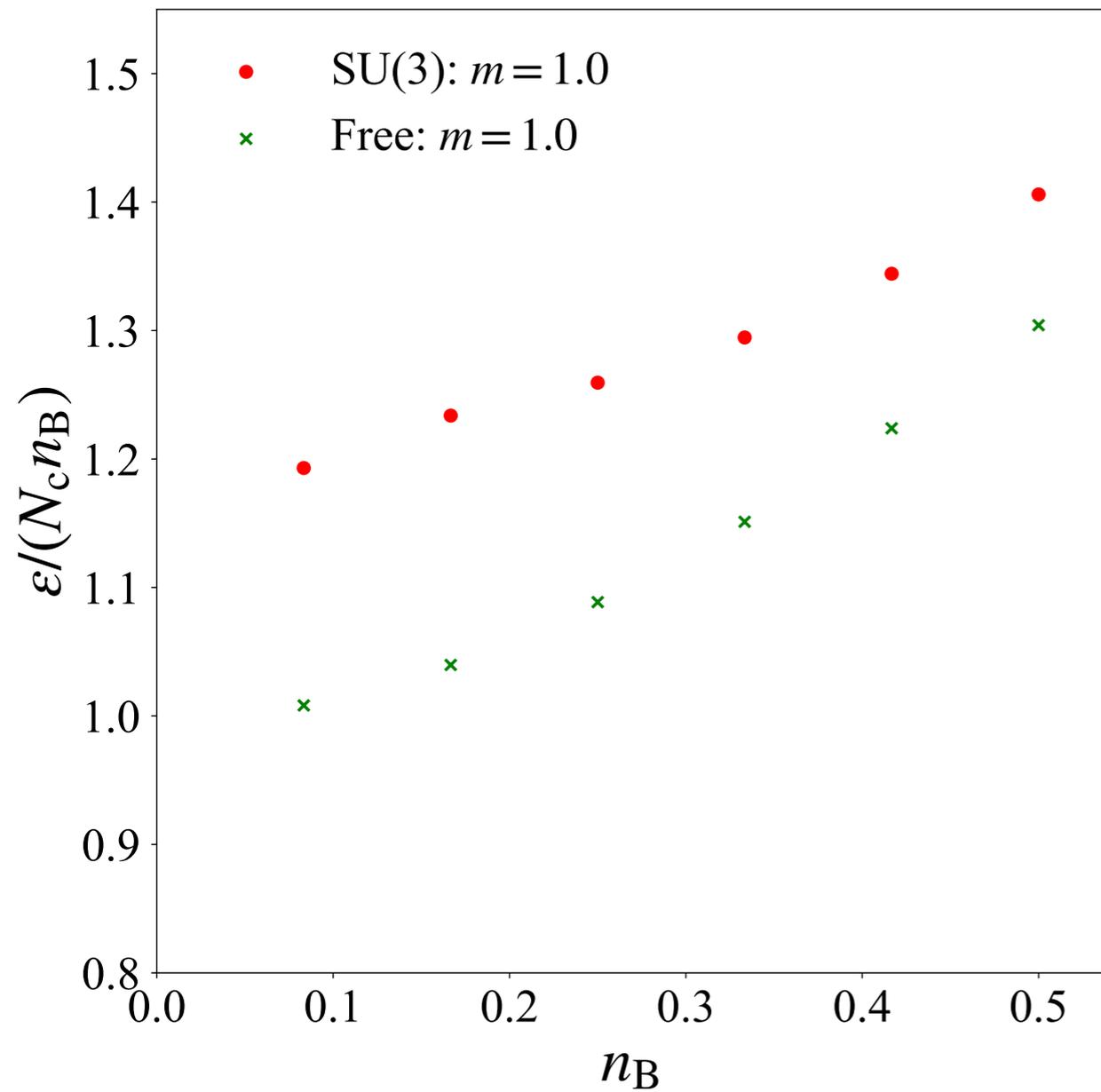
SU(3) results

Pressure and Baryon number

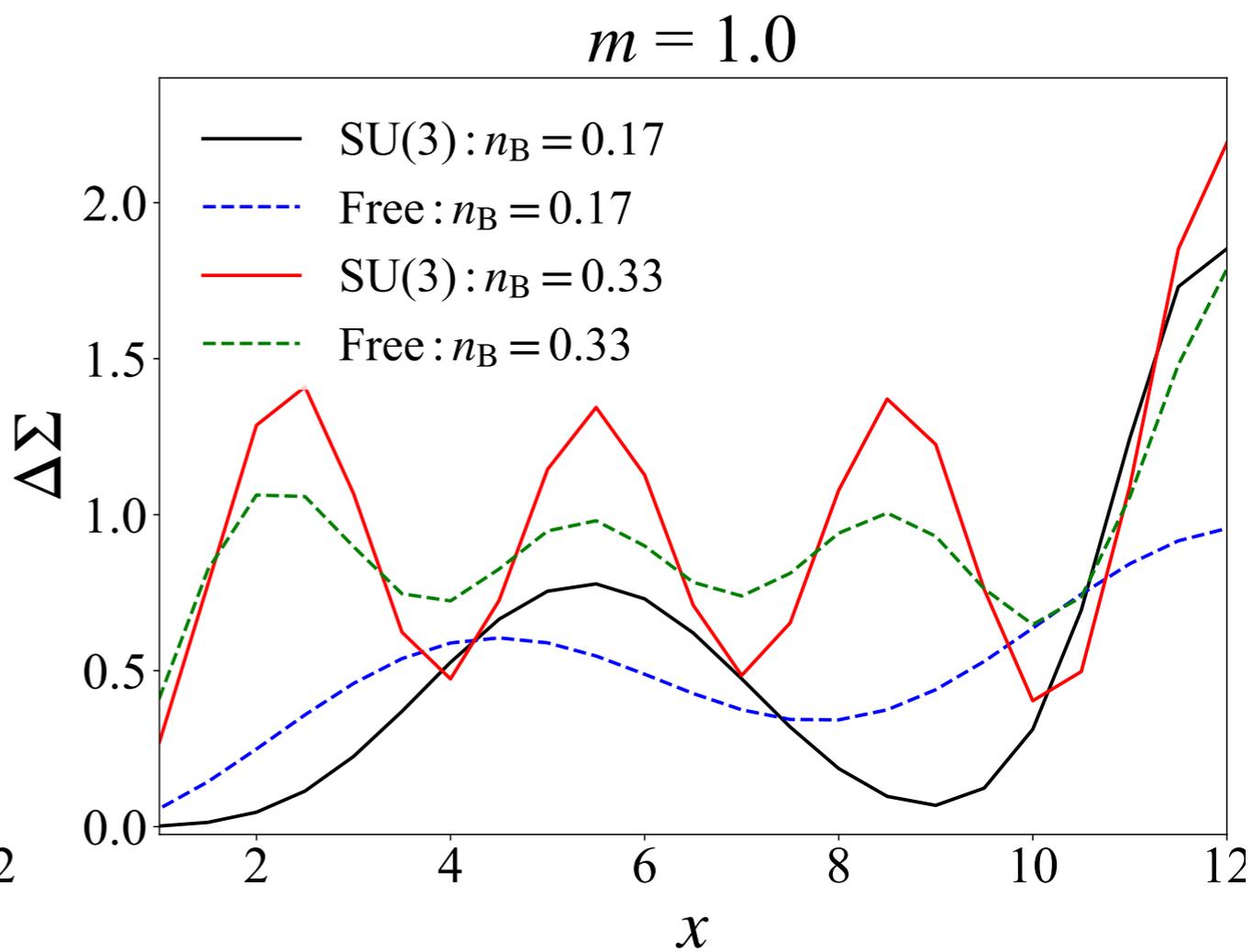
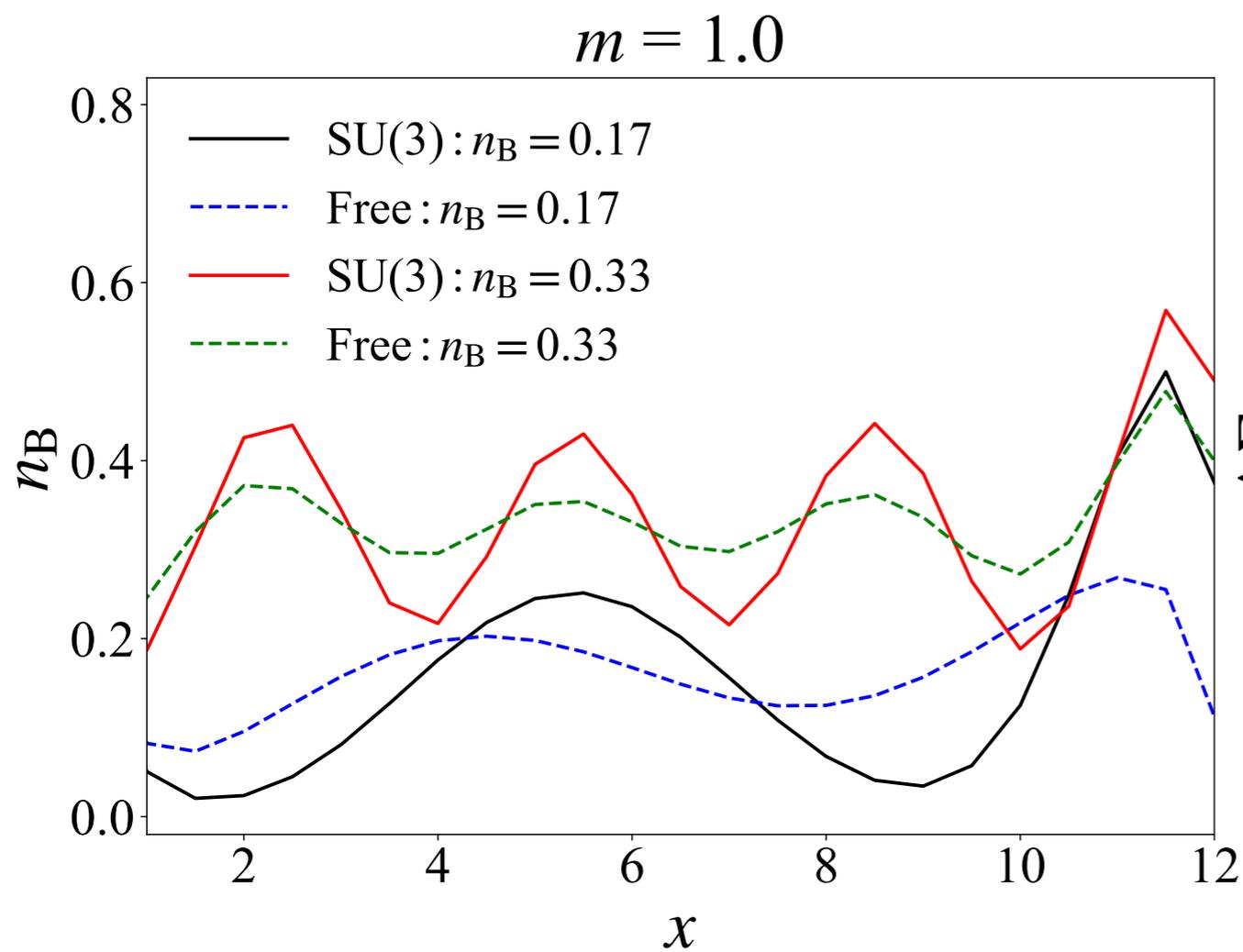
- $N=48$, $w=2$. Threshold value is 3.58.



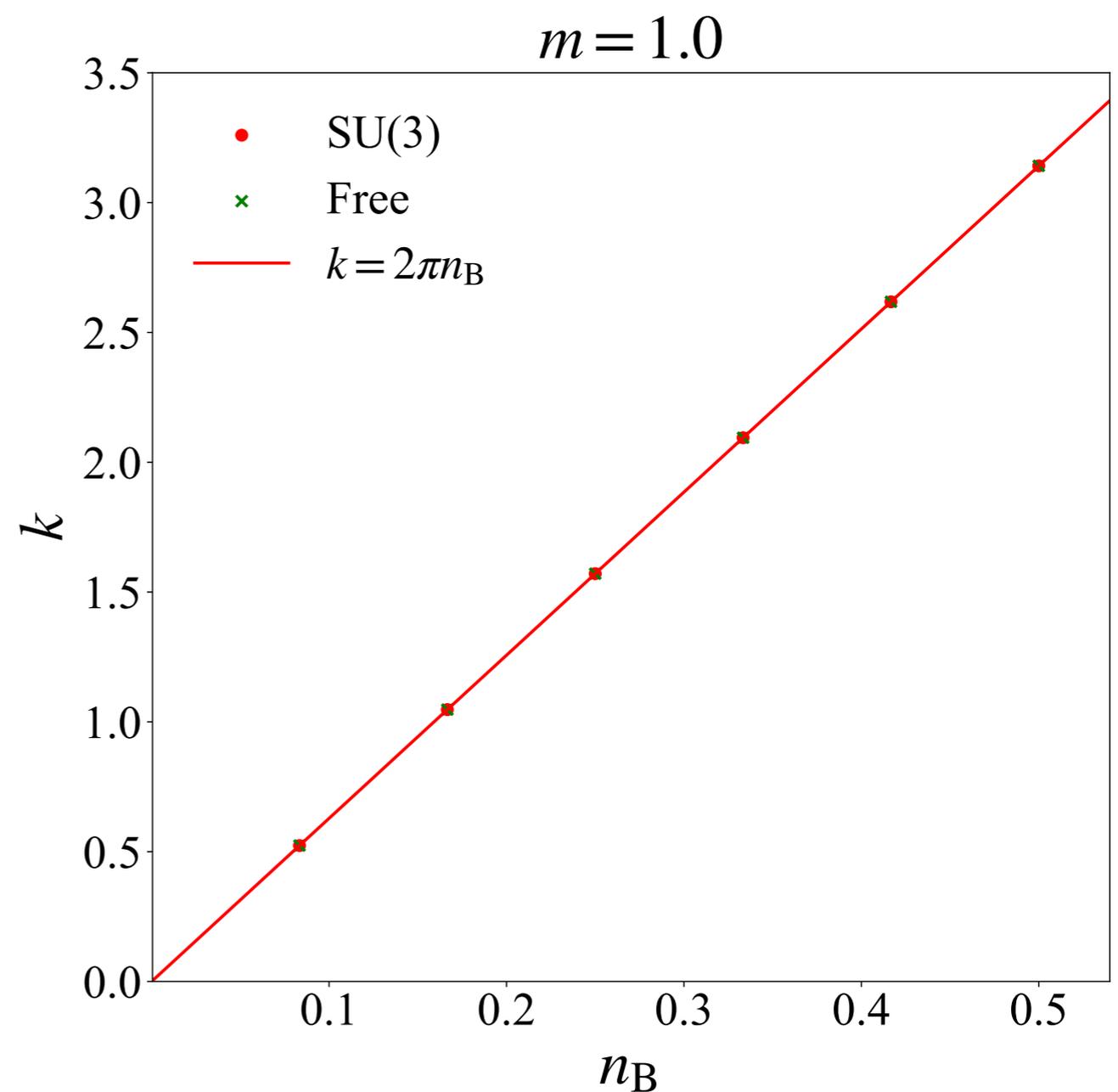
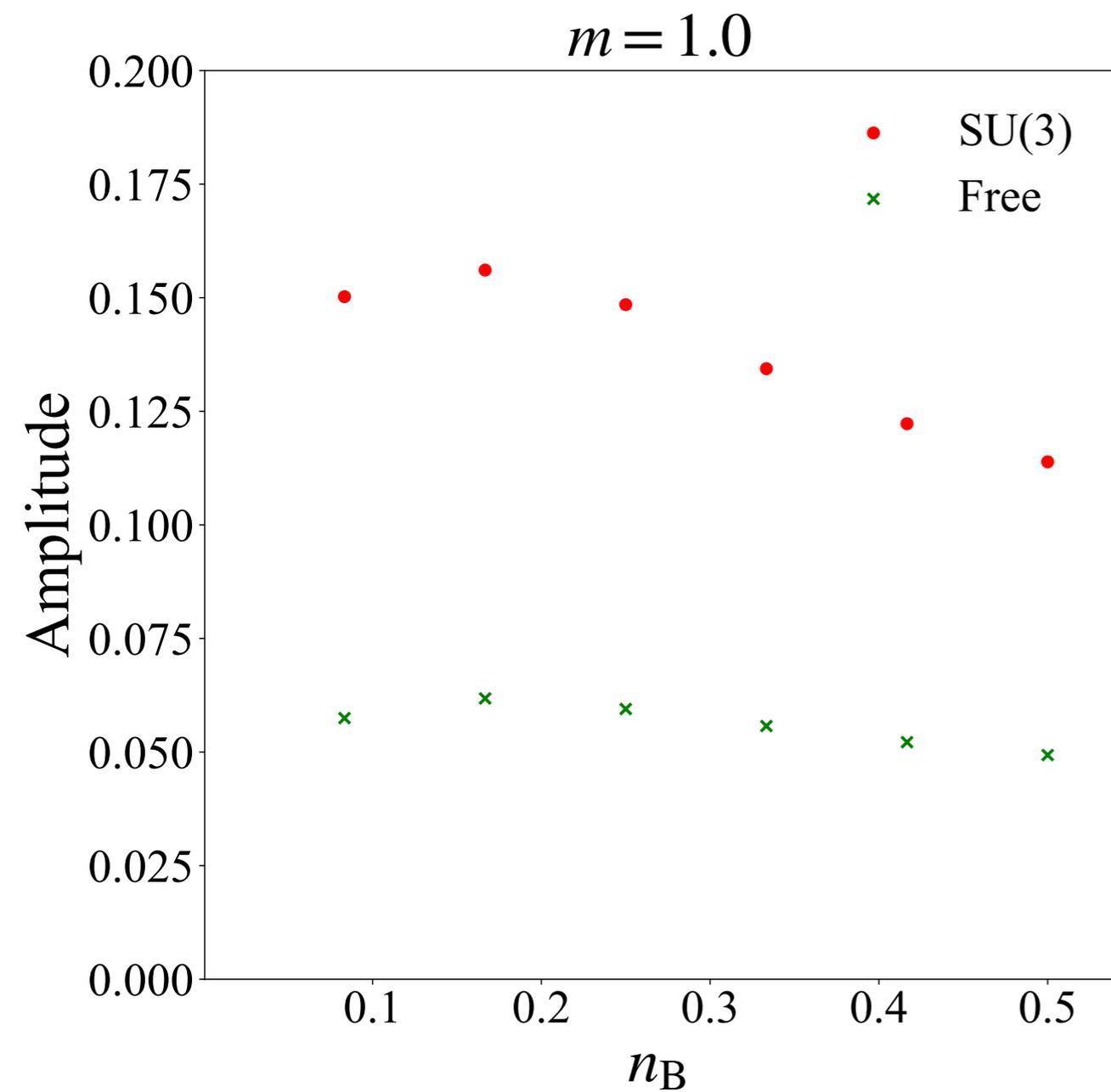
Average of μ_B and ratio



Spatial modulation

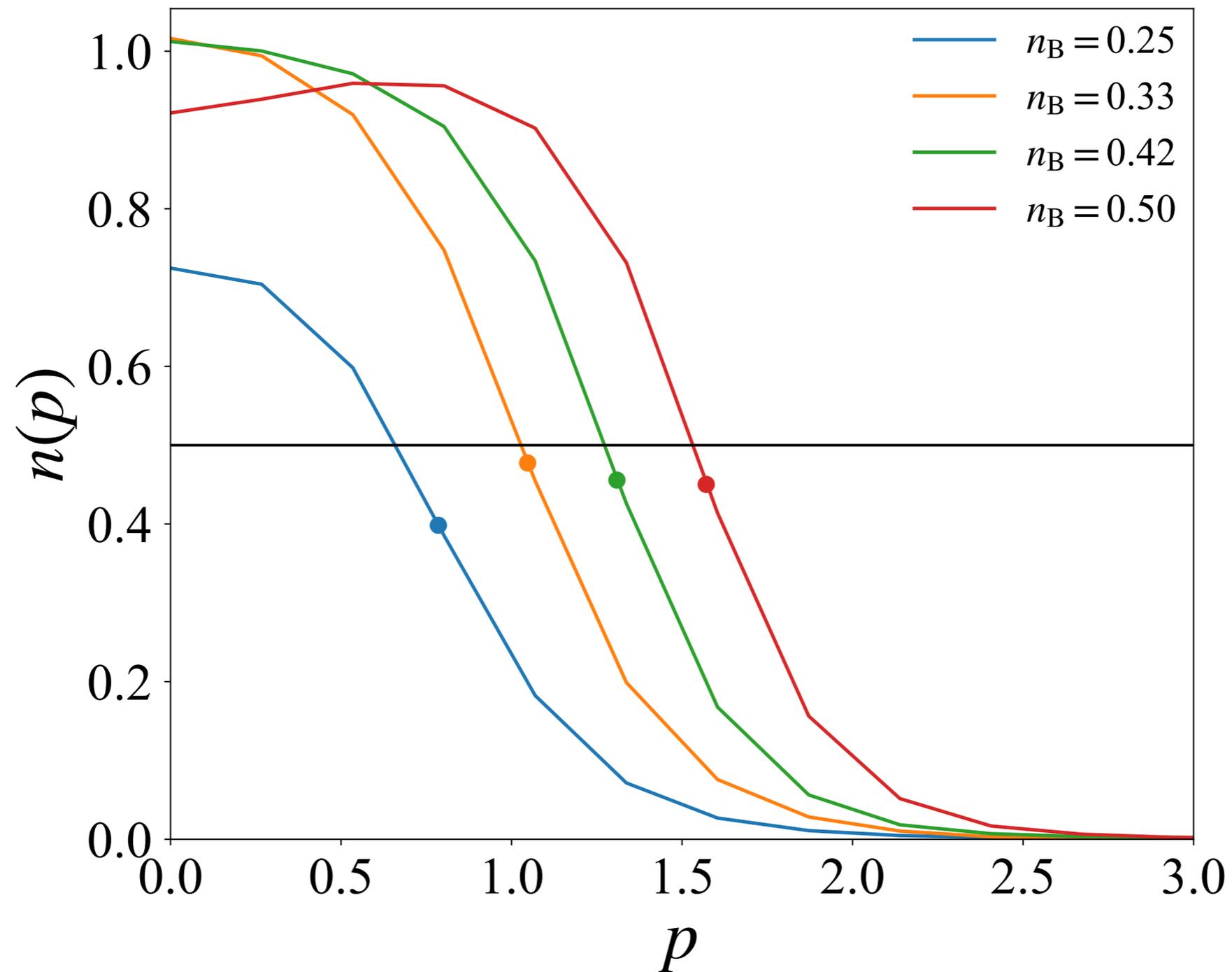


Wavenumber and amplitude



Quark distribution function

$$m = 1.0$$



Summary

- Variational numerical calculation based on Hamiltonian formalism.
 - The transition from baryonic matter to quark matter occurs near $n_B \sim 0.2$.
 - Inhomogeneous state in QCD₂ occurs and is more enhanced than in free theory.
 - Fermi surface forms about $n_B \sim 0.2$.
- Baryon occupation probability? [See talk by T. Kojo](#)
- Real-time dynamics? [Florio, et.al., \(2023\)](#)