

冷却原子気体における 交流スピン伝導率

Yuta Sekino, Hiroyuki Tajima, & Shun Uchino, arXiv:2103.02418 (accepted by PRResearch)
Hiroyuki Tajima, Yuta Sekino, & Shun Uchino, PRB 105, 064508 (2022)

関野 裕太 (理研CPR, 理研iTHEMS)



iTHEMS

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田島裕之 (東大理) 内野瞬 (原研先端研)

Outline of this talk

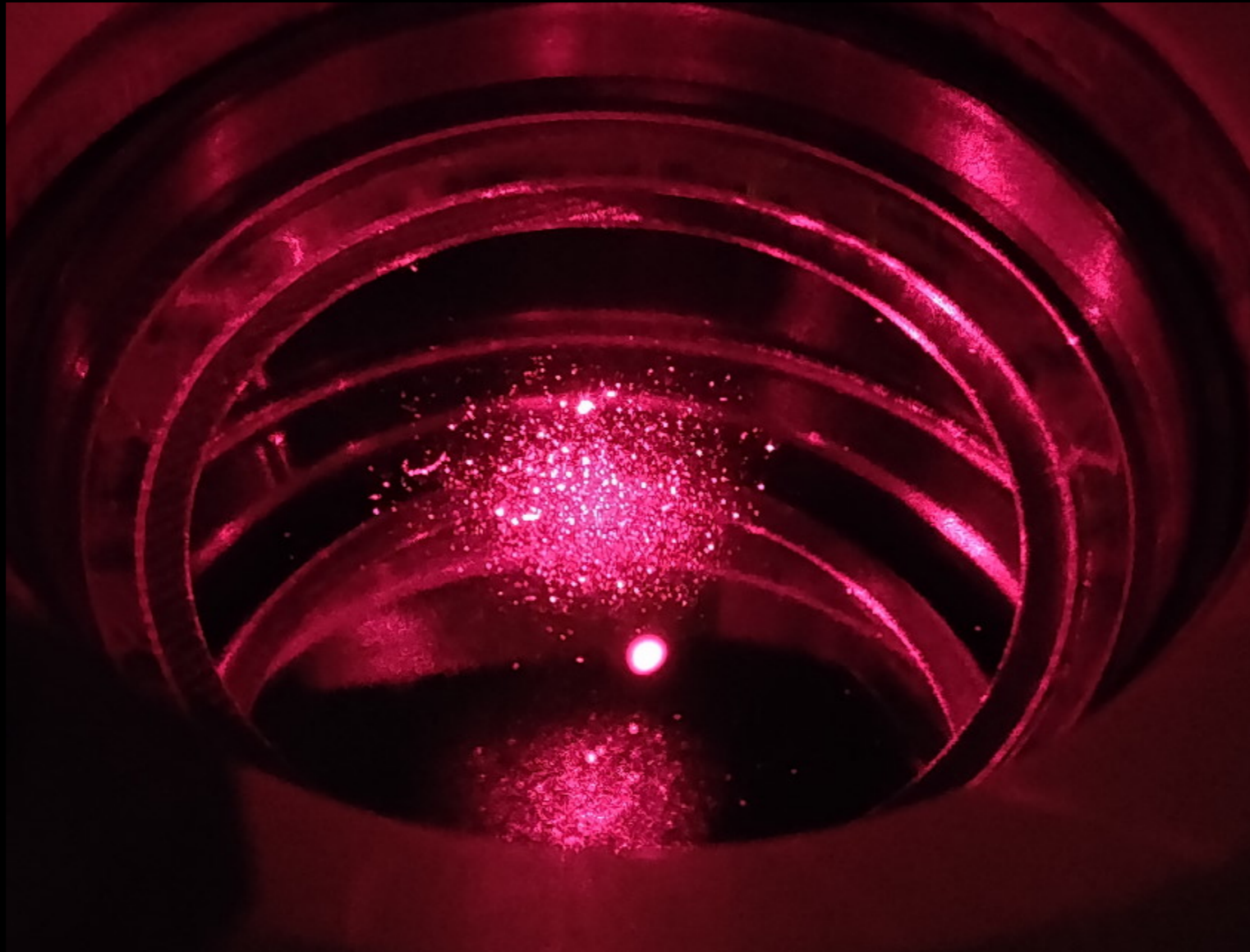
1. Introduction
2. Optical spin conductivity
 - Proposal of measurement
3. Theoretical studies
 - Formalism
 - Fermi superfluids
 - Tomonaga-Luttinger liquid
4. Summary

Outline of this talk

1. Introduction
2. Optical spin conductivity
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3. Theoretical studies
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4. Summary

Ultracold atoms

- ▶ Very pure & highly controllable atomic gases
- ▶ Ideal research platform for quantum many-body phenomena



High controllability

1. Quantum statistics & spin degrees of freedom

Bose atom: ${}^7\text{Li}$, ${}^{23}\text{Na}$, ${}^{39}\text{K}$, ...

Fermi atom: ${}^6\text{Li}$, ${}^{40}\text{K}$, ...

Hyperfine states

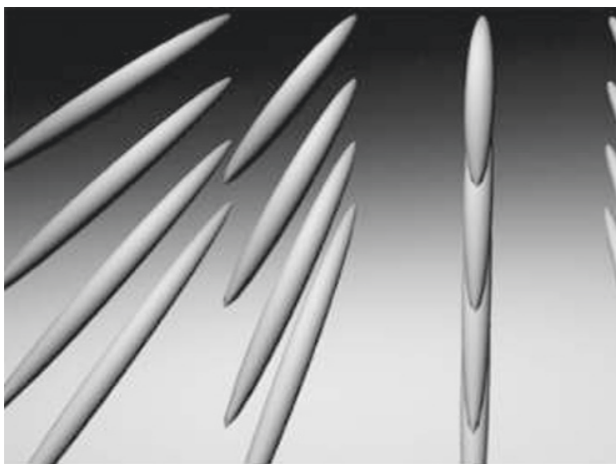


Spin

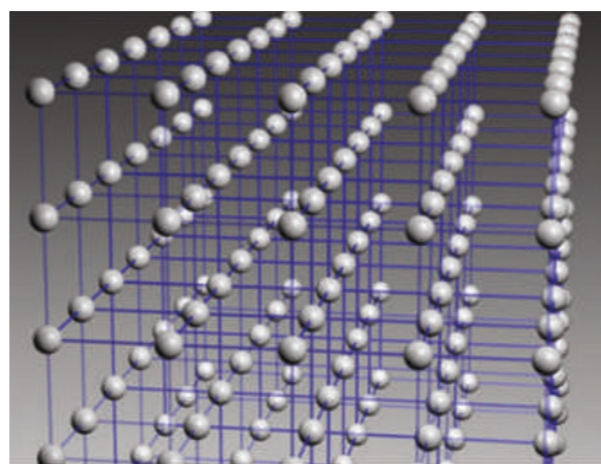
2. Spatial geometry of gas

3D, 2D, 1D, Lattice, ...

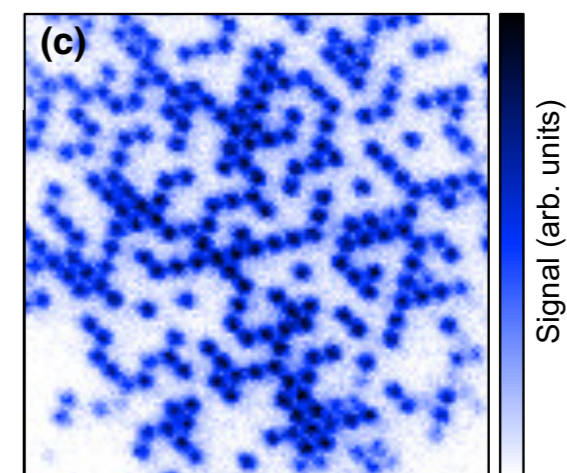
1D



Cubic lattice



Triangular lattice



Bloch, Nat. Phys. (2005)

Yang et al.(Virginia), PRX
Quantum (2021)

High controllability

1. Quantum statistics & spin degrees of freedom

Bose atom: ${}^7\text{Li}$, ${}^{23}\text{Na}$, ${}^{39}\text{K}$, ...

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Hyperfine states



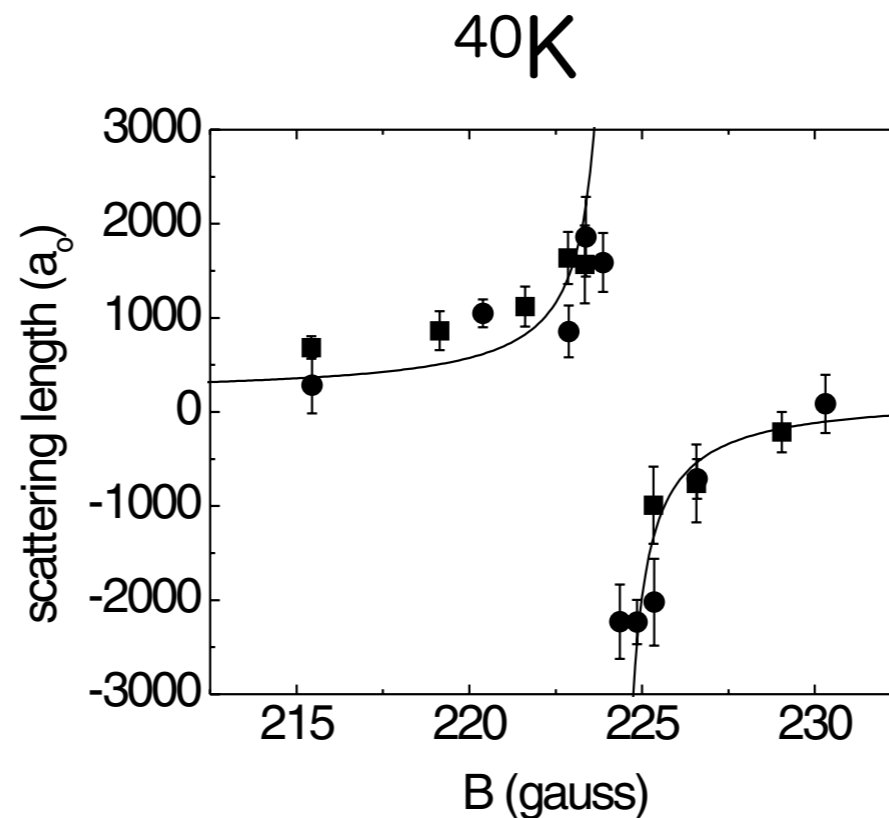
Spin

2. Spatial geometry of gas

3D, 2D, 1D, Lattice, ...

3. Interaction between atoms

Feshbach resonance, Lattice depth, ...



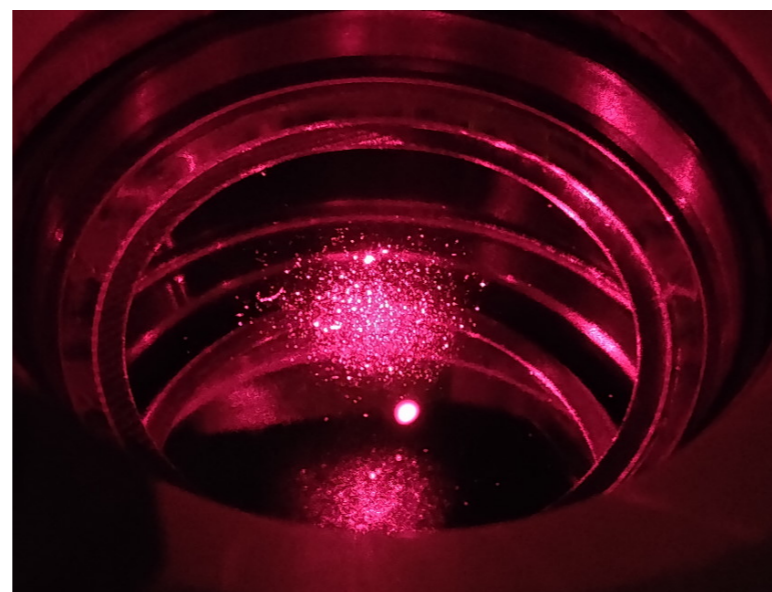
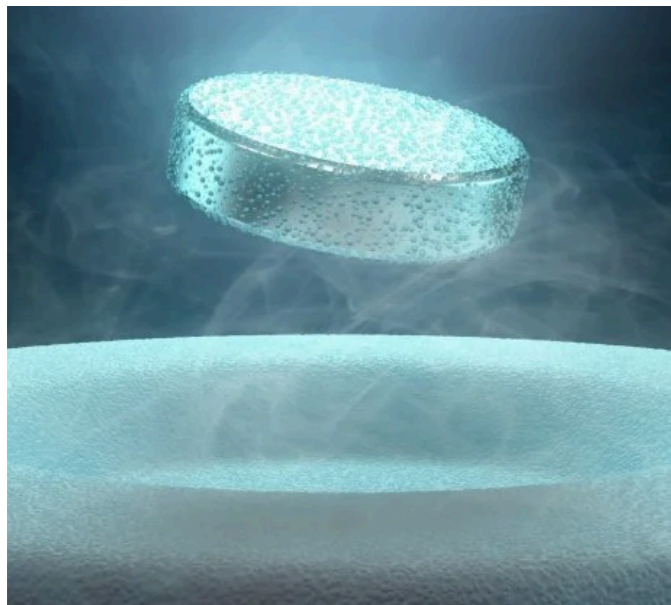
Regal & Jin (JILA), PRL(2003)

Major research directions

Ideal platform to study quantum many-body phenomena

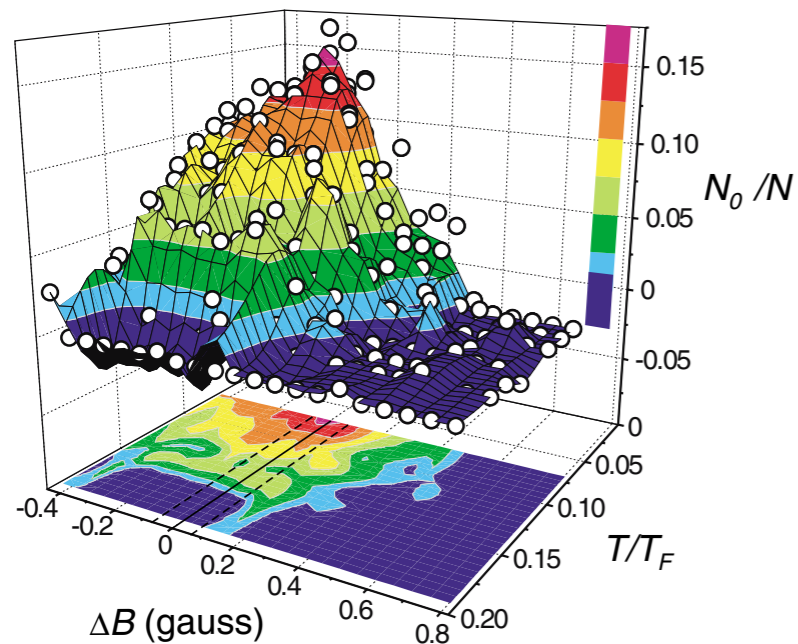
1. Novel quantum phenomena
2. Quantum computation
3. Analog quantum simulation:
cold-atomic systems equivalent/similar to other interesting systems

Superconductors ↔ Fermi atoms in superfluid phase ↔ Neutron superfluid in neutron stars



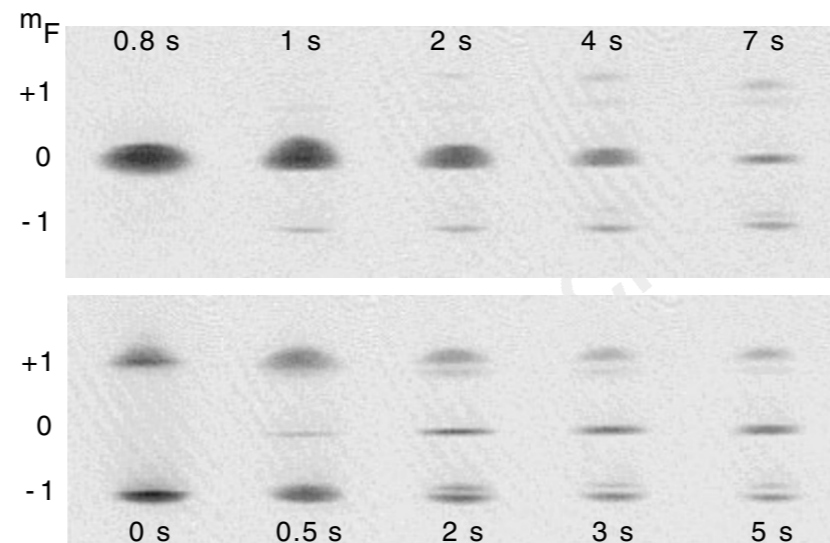
Various many-body states with spin

Fermi superfluid
(BCS-BEC crossover)



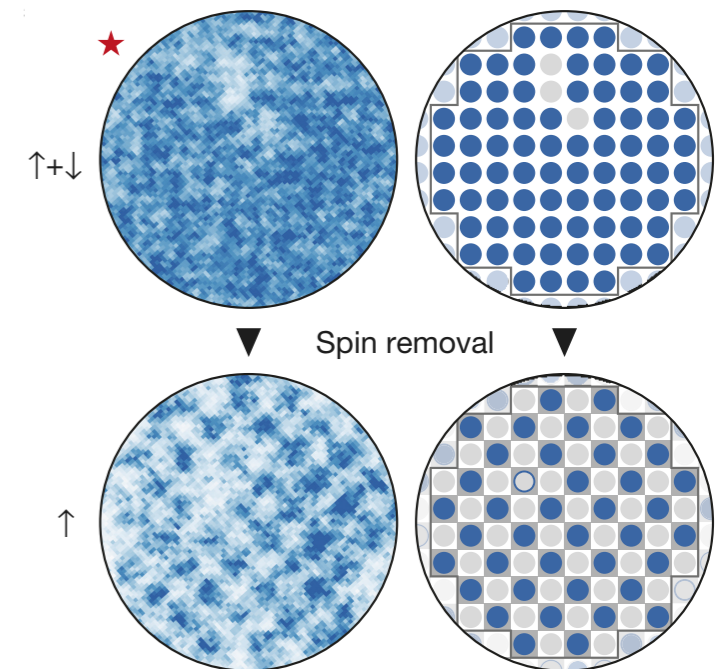
Regal et al., (JILA) PRL (2004)

Spinor Bose-Einstein
Condensate (BEC)



Stenger et al., (MIT) Nature (1998)

Heisenberg
antiferromagnet

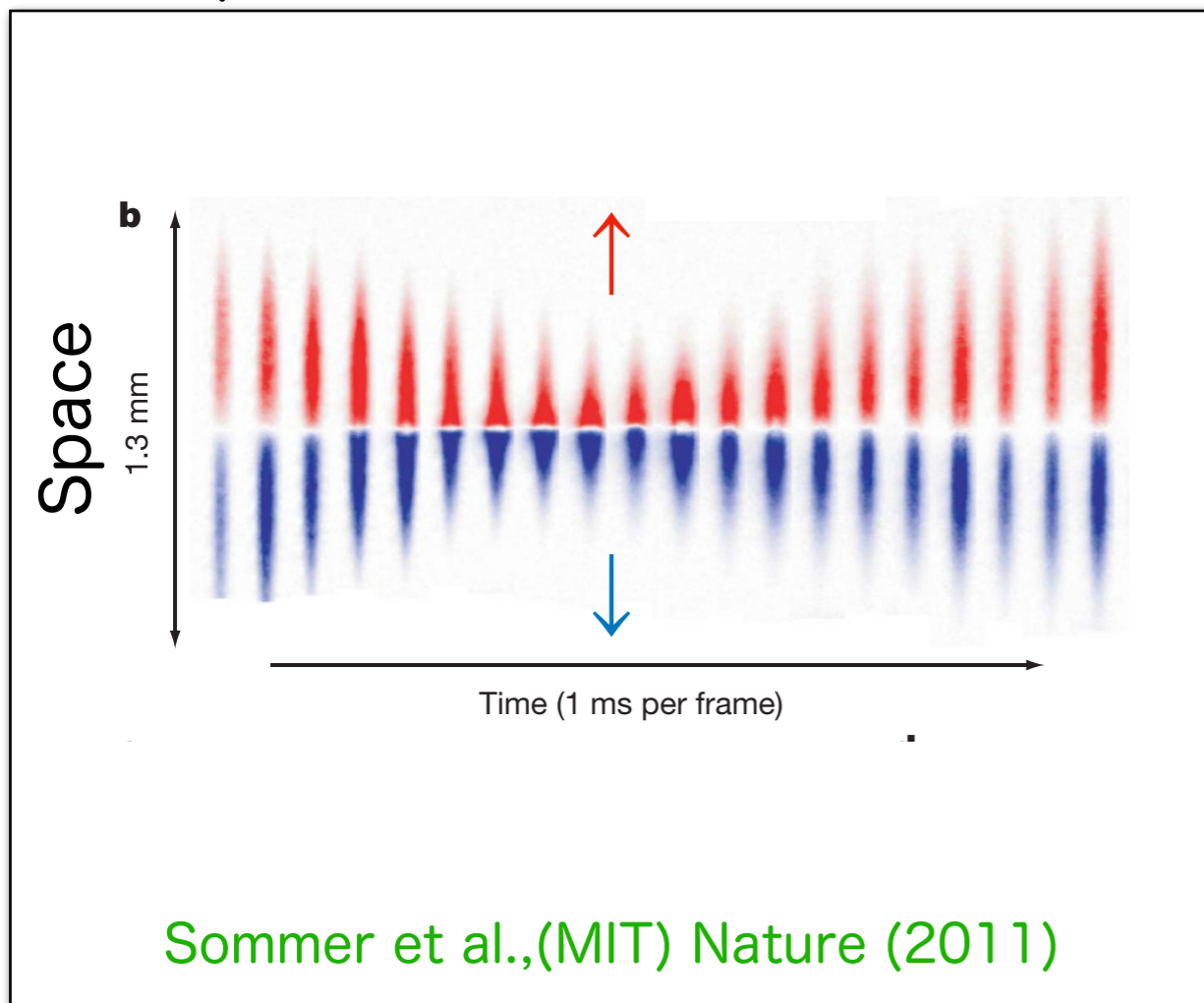


Mazurenko et al.,(Harvard) Nature (2017)

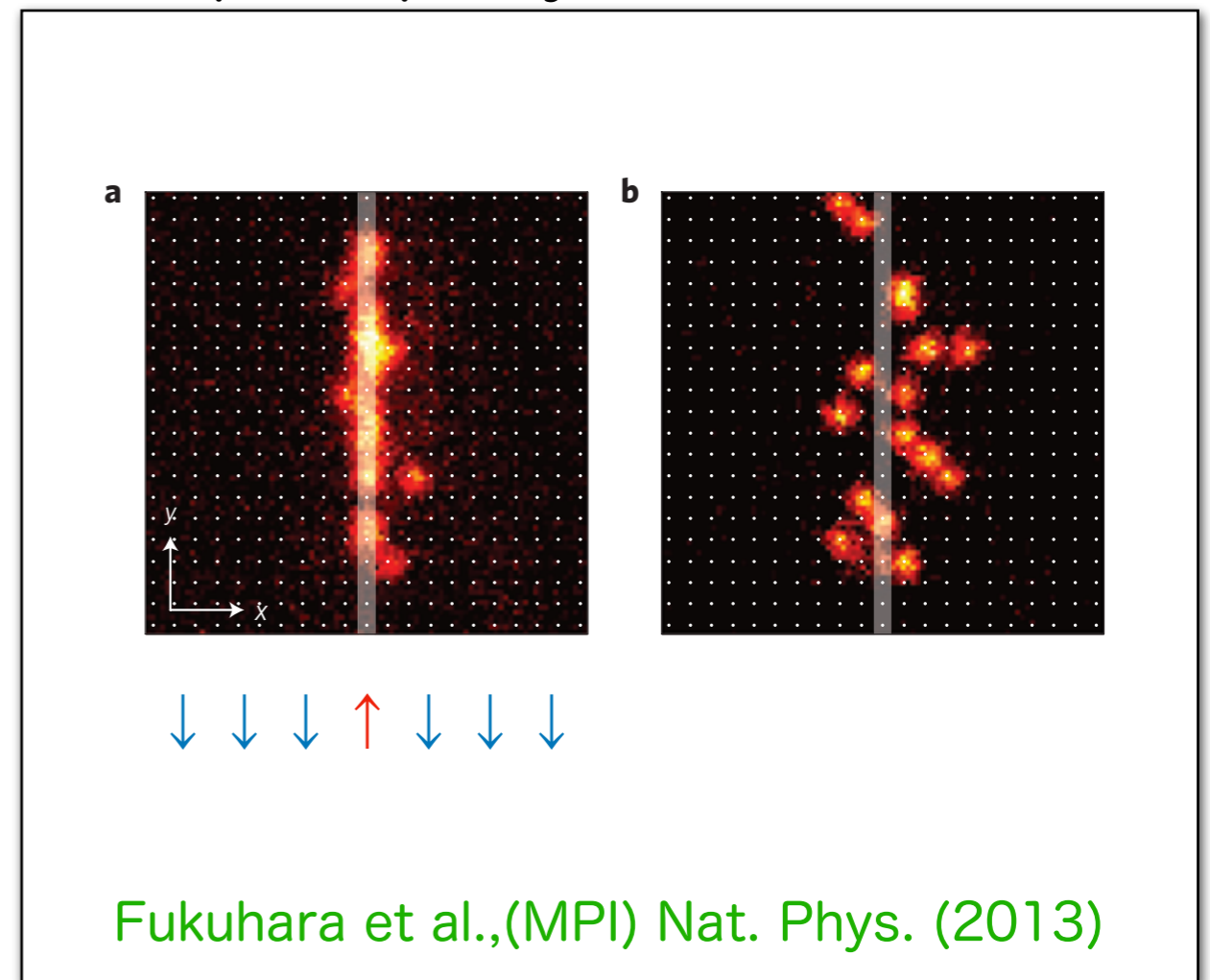
Spin dynamics with cold atoms

- Ideal experimental grounds to study spin dynamics
Spin-resolved manipulation & detection

1. Spin diffusion w/o lattice



2. Spin impurity on lattice



Spin dynamics with cold atoms

- Ideal experimental grounds to study spin dynamics
Spin-resolved manipulation & detection

1. Spin diffusion w/o lattice

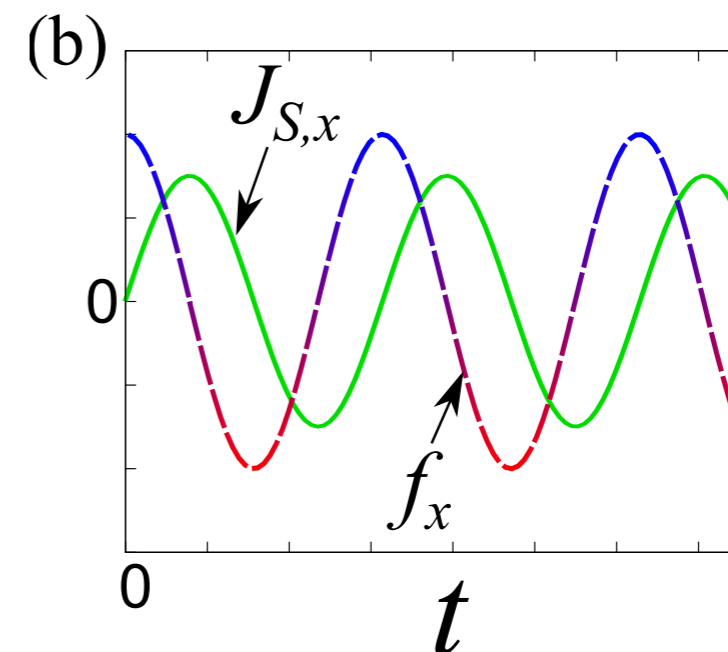
2. Spin impurity on lattice

3. AC spin transport (Our proposal)

YS, Tajima, & Uchino, accepted by PRResearch (2022)



Driven by magnetic or optical fields

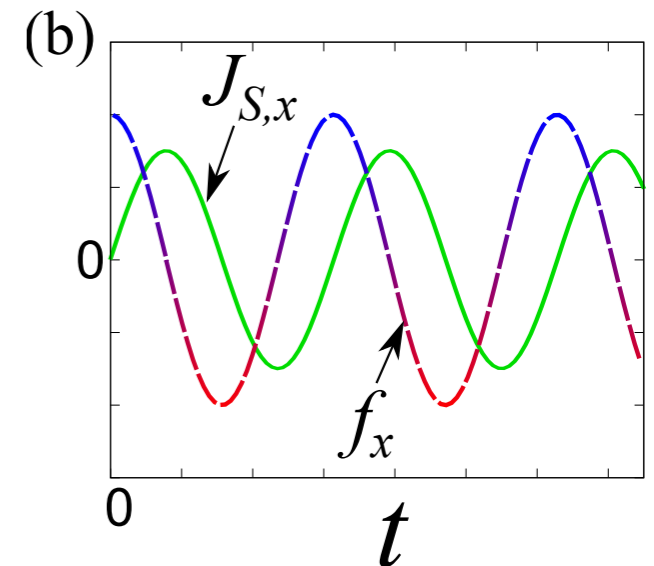


Today's topic

▶ Optical (AC) spin conductivity

$$\sigma_{\alpha,\beta}^{(S)}(\omega) = \tilde{J}_{S,\alpha}(\omega) / \tilde{f}_{\beta}(\omega)$$

$(\alpha, \beta = x, y, z)$



Measurable in cold-atom experiments

YS, Tajima, & Uchino, accepted by PRResearch (2022)

▶ Significance of $\sigma_{\alpha\beta}^{(S)}(\omega)$

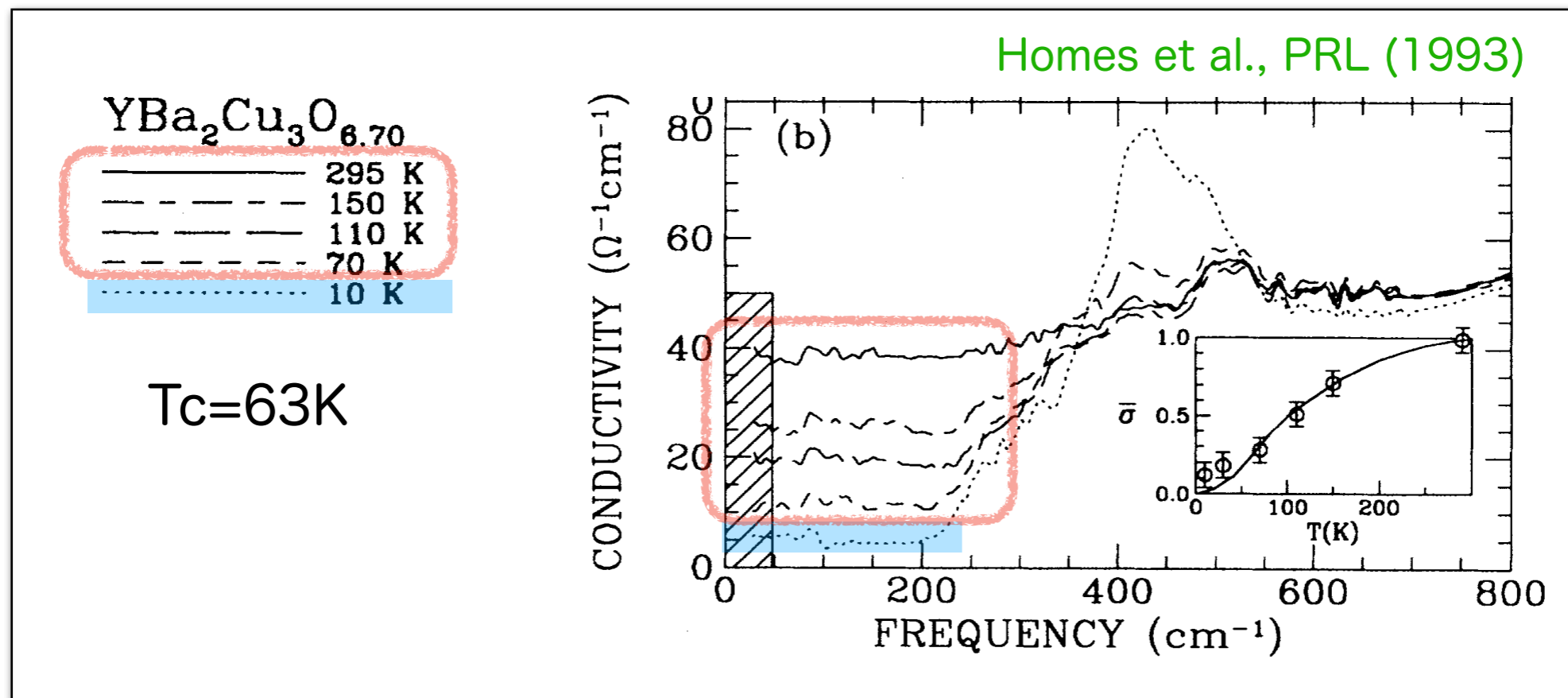
YS, Tajima, & Uchino, accepted by PRResearch (2022)
Tajima, YS, & Uchino, PRB (2022)

1. Elusive in solid-state systems
2. **Powerful probe for quantum many-body states**
BCS-BEC crossover, Tomonaga-Luttinger liquid(TLL), Spinor BEC,...
3. Widely applicable probe for clean systems

2. Optical conductivity for solids

► Powerful probe for exotic electron systems

Superconductor, Pseudogap phase, Non-Fermi liquid, Dirac fermions, ...



Optical spin conductivity would also be a useful probe for nontrivial spin dynamics

Outline of this talk

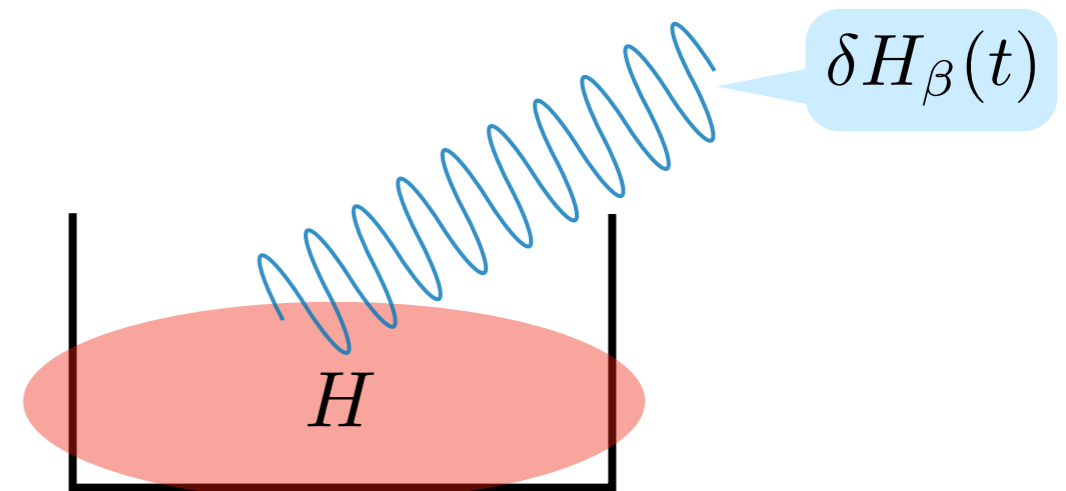
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Measurement scheme: set up

YS, Tajima, & Uchino, arXiv:2103.02418

► Total Hamiltonian:

$$H(t) = H + \delta H_{\beta}(t)$$



Cold atoms with **spin (at least S_z) conserved**

Spin: $S=1/2, 1, 3/2, \dots$

Zeeman field, synthetic gauge field, \dots

trap & lattice potentials, \dots

BCS-BEC crossover, Spinor BEC, ferromagnets/antiferromagnets, \dots

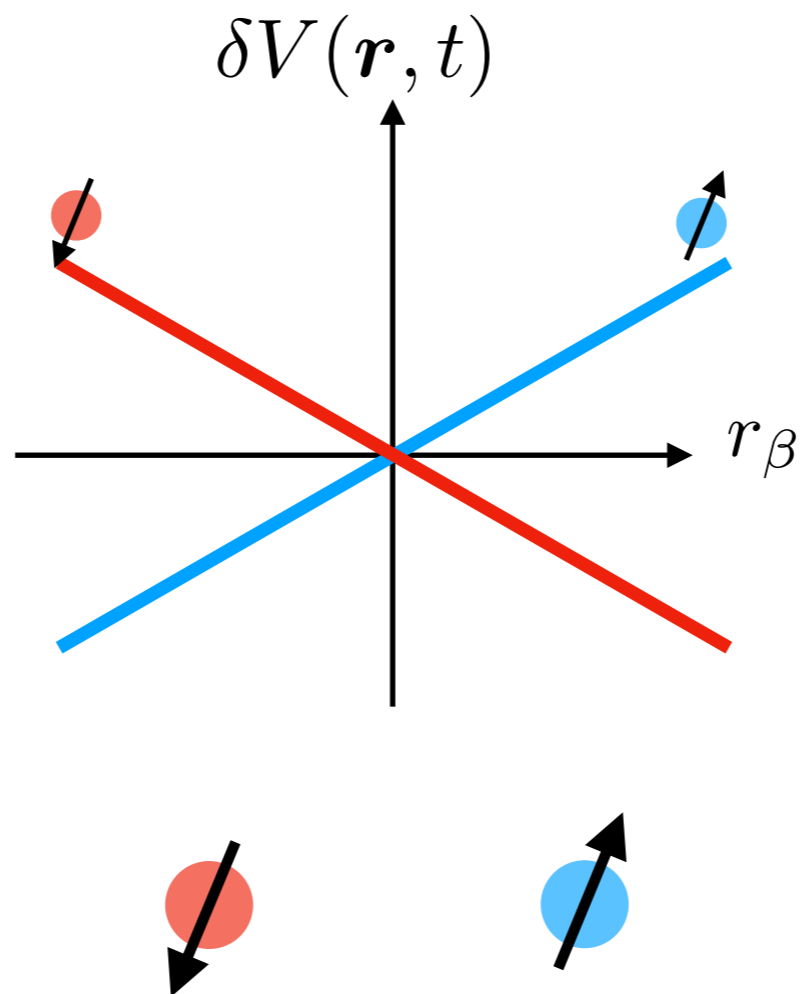
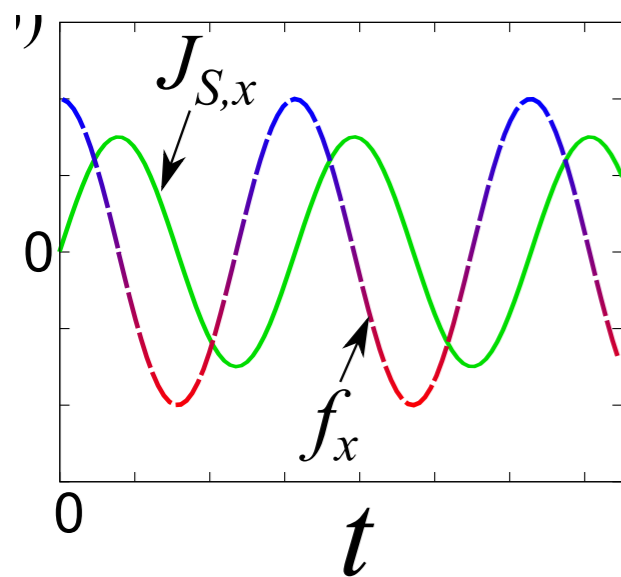
(Extension to spin-nonconserving & nonequilibrium systems is possible)

How to induce AC spin current

YS, Tajima, & Uchino, arXiv:2103.02418

► Time-dependent force coupled to spin density S_z

$$\delta H_\beta(t) = - \int d\mathbf{r} f_\beta(t) r_\beta S_z(\mathbf{r}), \quad (\beta = x, y, z)$$

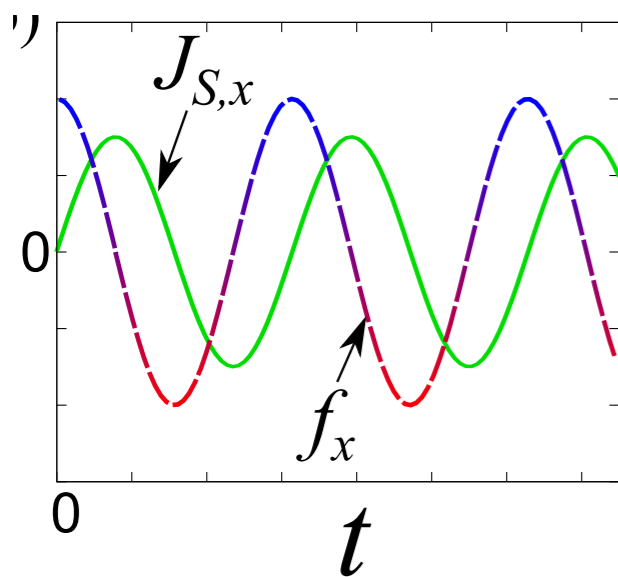


How to induce AC spin current

YS, Tajima, & Uchino, arXiv:2103.02418

► Time-dependent force coupled to spin density S_z

$$\delta H_\beta(t) = - \int d\mathbf{r} \, \underline{f_\beta(t)} r_\beta \underline{S_z(\mathbf{r})}, \quad (\beta = x, y, z)$$



Single-frequency driving force toward $\beta = x, y, z$

$$f_\beta(t) = F_\beta \cos(\omega_0 t)$$

1. Magnetic field gradient

Medley et al., (MIT) PRL (2011); Jotzu et al., (ETH) PRL (2015)

2. Optical Stern-Gerlach effect

Taie et.al., (Kyoto) PRL (2010)

How to extract spin conductivity

YS, Tajima, & Uchino, arXiv:2103.02418

1. Spin current: $\langle \mathbf{J}_S(t) \rangle = \frac{d}{dt} \left\langle \int d\mathbf{r} \mathbf{r} S_z(\mathbf{r}, t) \right\rangle \equiv \frac{d}{dt} \langle \mathbf{X}_S(t) \rangle$ (Spin conservation)
2. Spin conductivity: $\langle \tilde{\mathbf{J}}_{S,\alpha}(\omega) \rangle = \sigma_{\alpha\beta}^{(S)}(\omega) \tilde{f}_\beta(\omega)$ (Ohm's law in frequency space)
3. Driving force: $f_\beta(t) = F_\beta \cos(\omega_0 t)$



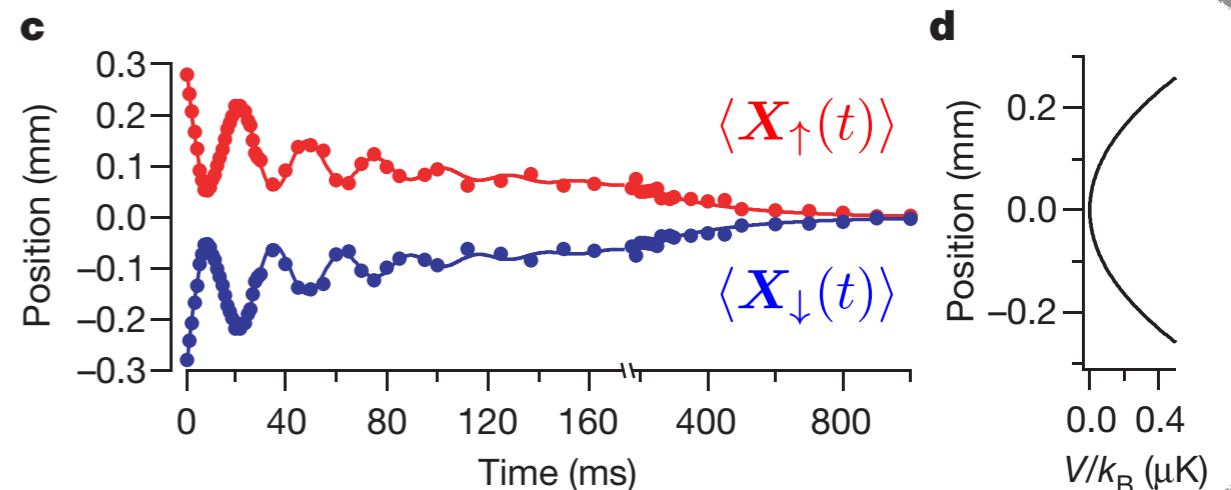
$$\frac{\langle \delta X_{S,\alpha}(t) \rangle}{F_\beta} = -\frac{\text{Im} \sigma_{\alpha\beta}^{(S)}(\omega_0)}{\omega_0} \cos(\omega_0 t) + \frac{\text{Re} \sigma_{\alpha\beta}^{(S)}(\omega_0)}{\omega_0} \sin(\omega_0 t)$$

Measurement of $\langle X_{S,\alpha}(t) \rangle \longrightarrow \sigma_{\alpha\beta}^{(S)}(\omega = \omega_0)$

e.g. Experiment on spin diffusion (w/o $f_\beta(t)$)

$$\langle \mathbf{X}_S(t) \rangle = \langle \mathbf{X}_\uparrow(t) \rangle - \langle \mathbf{X}_\downarrow(t) \rangle$$

Sommer et al.,(MIT) Nature (2011)



Outline of this talk

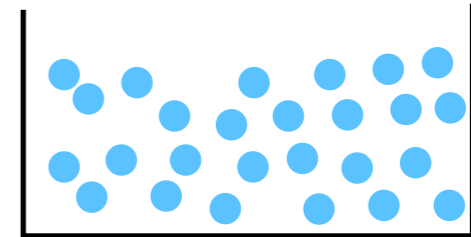
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Theoretical studies

► Our works on optical spin conductivity in homogeneous gases

YS, Tajima, & Uchino, arXiv:2103.02418

Tajima, YS, & Uchino, PRB (2022)



1. ~~1D~~ $S=1/2$ superfluid Fermi gas with spin gap
2. ~~1D~~ 1D p-wave Fermi superfluid with topological phase transition
3. ~~1D~~ Tomonaga-Luttinger liquid
4. $S=1$ polar BEC with gapped or gapless spin modes

General relations

Enss & Haussmann PRL (2012) Enss, Euro. Phys. J. Special Topics (2013)

YS, Tajima, & Uchino, arXiv:2103.02418 ($\hbar = k_B = 1$)

1. Kubo formula

$$\sigma_{\alpha\beta}^{(S)}(\omega) = \frac{i}{\omega^+} \left(\delta_{\alpha\beta} \sum_{s_z} \frac{s_z^2 N_{s_z}}{m} + \chi_{\alpha\beta}(\omega) \right) \quad \alpha, \beta \in \{x, y, z\} \quad \omega^+ \equiv \omega + 0^+$$

Magnetic quantum #: $s_z = -S, -S + 1, \dots, S$

Spin-current response func.: $\chi_{\alpha\beta}(\omega) = -i \int_{-\infty}^{\infty} dt e^{i\omega^+ t} \theta(t) \langle [J_{S,\alpha}(t), J_{S,\beta}(0)] \rangle_{\text{eq}}$

Particle # in the s_z channel: N_{s_z}

2. f-sum rule

$$\int_{-\infty}^{\infty} \frac{d\omega}{\pi} \text{Re} \sigma_{\alpha\beta}^{(S)}(\omega) = \delta_{\alpha\beta} \sum_{s_z} \frac{s_z^2 N_{s_z}}{m}.$$

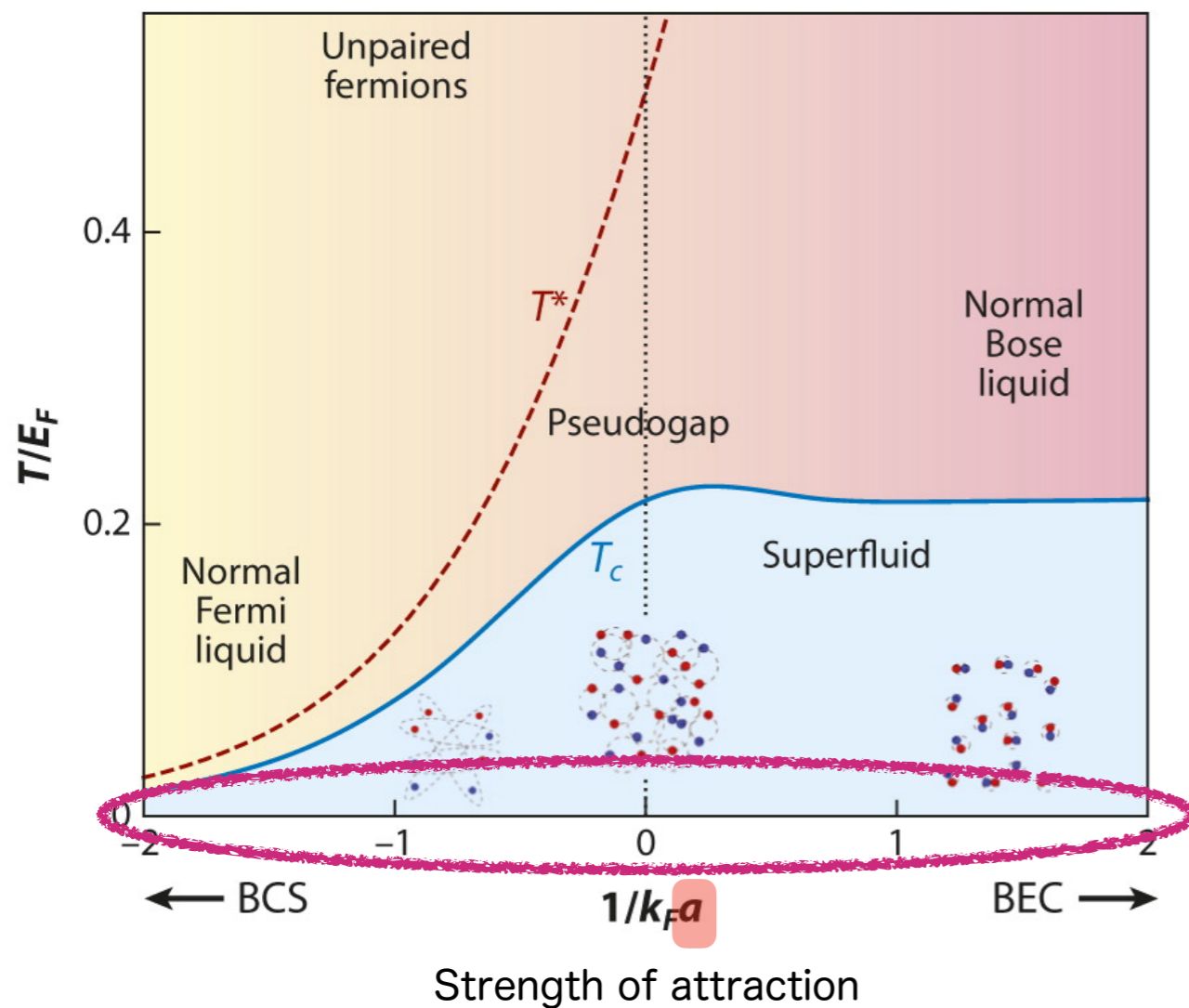
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Superfluid Fermi gas

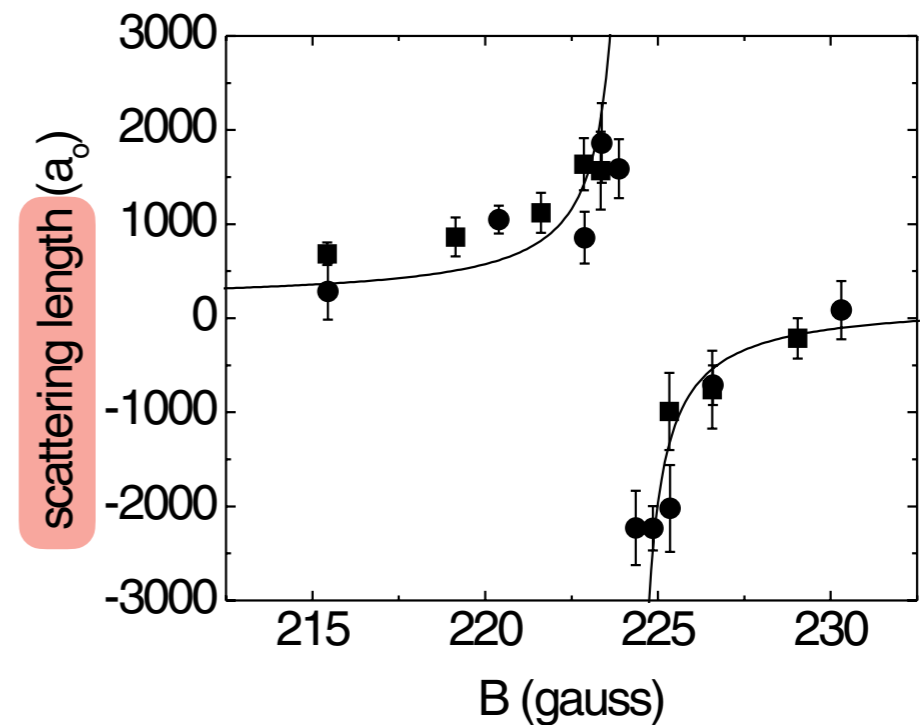
$$H = \int dx \left[\sum_{\sigma=\uparrow,\downarrow} \psi_{\sigma}^{\dagger} \left(-\frac{\nabla^2}{2m} - \mu \right) \psi_{\sigma} - g \psi_{\uparrow}^{\dagger} \psi_{\downarrow}^{\dagger} \psi_{\downarrow} \psi_{\uparrow} \right]$$

$g > 0$: S-wave attraction



Randeria & Taylor (2014)

BCS-Leggett mean-field theory

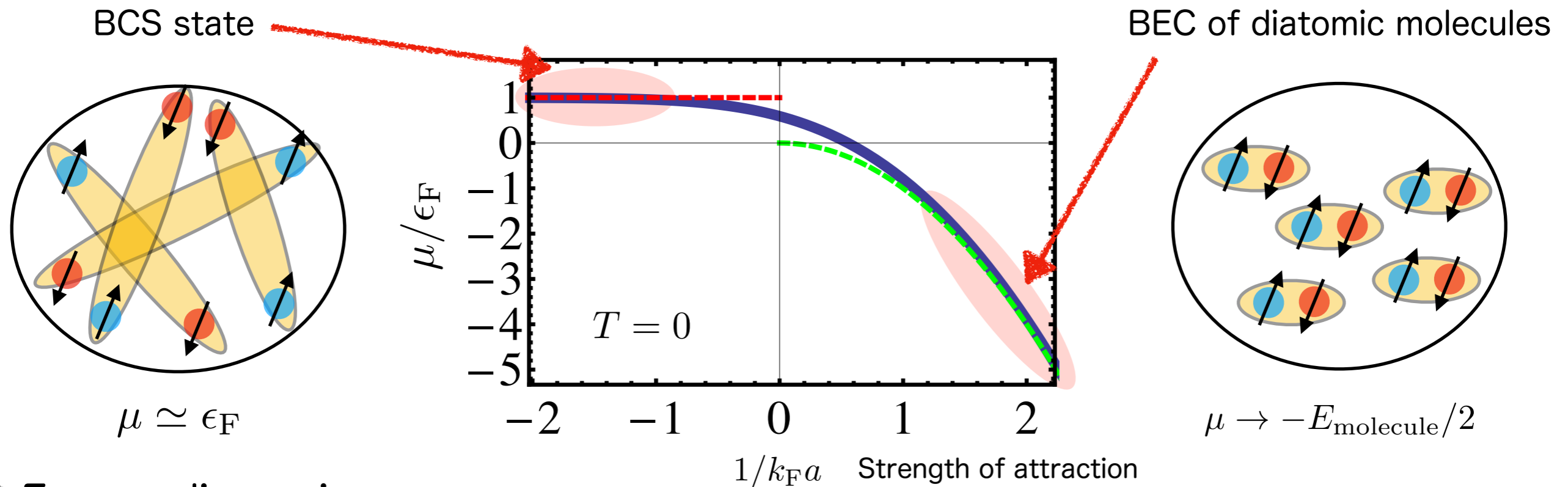


Regal & Jin (JILA), PRL(2003)

Eagles (1969); Leggett (1980)

BCS-BEC crossover

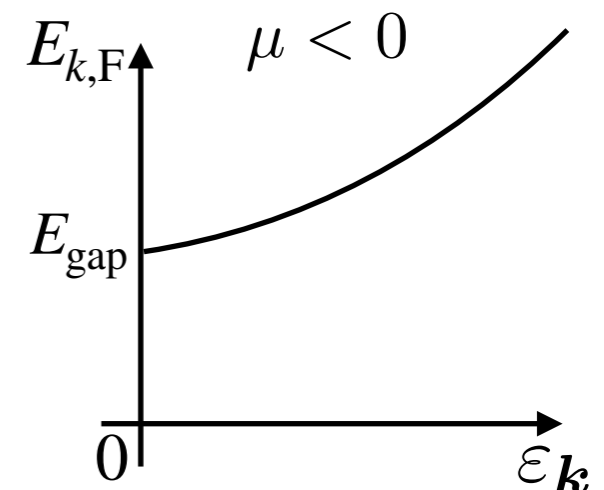
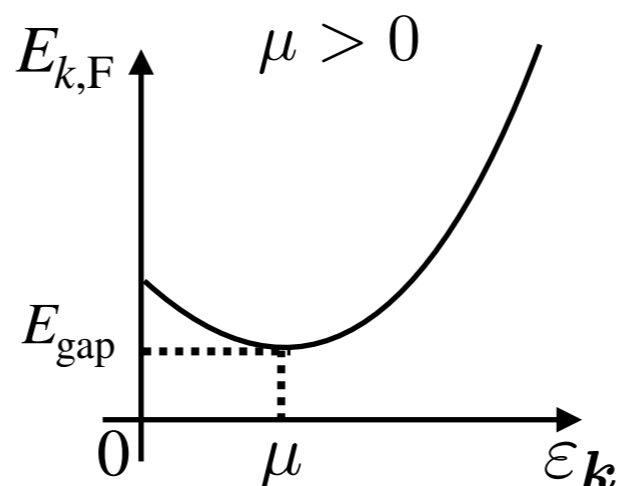
► Chemical potential @ T=0



► Energy dispersion

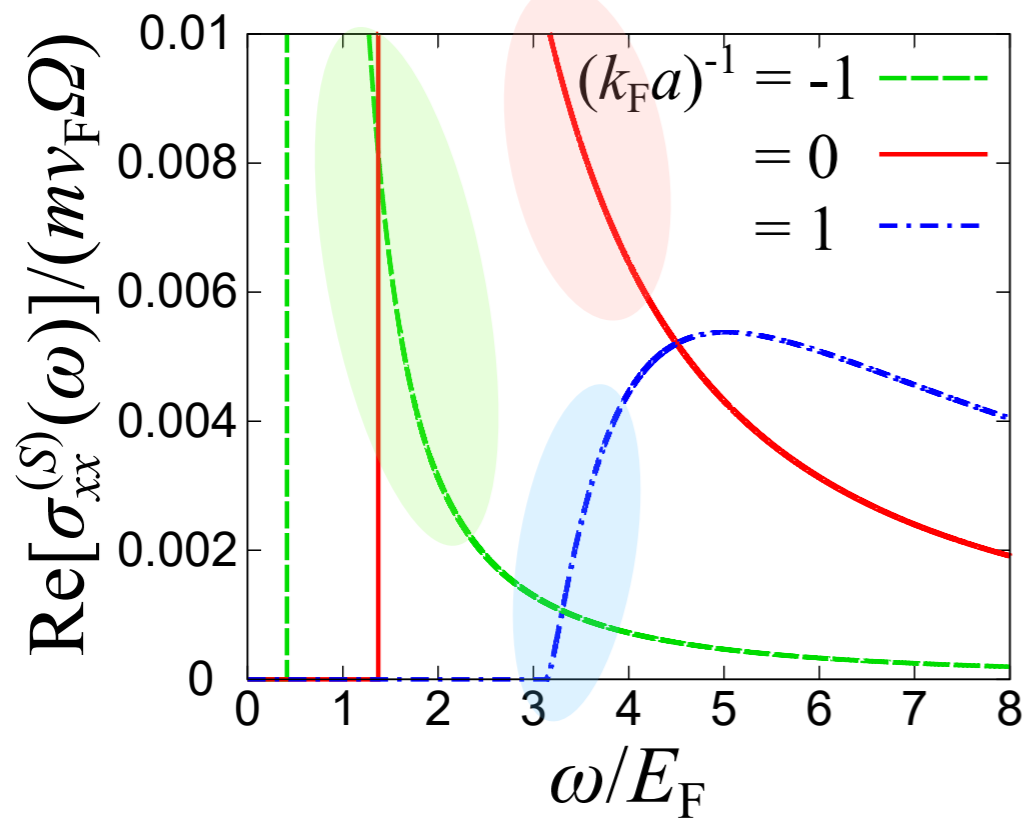
$$E_{\mathbf{k},F} = \sqrt{(\epsilon_{\mathbf{k}} - \mu)^2 + \Delta^2}$$

$$(\epsilon_{\mathbf{k}} = \mathbf{k}^2/2m)$$



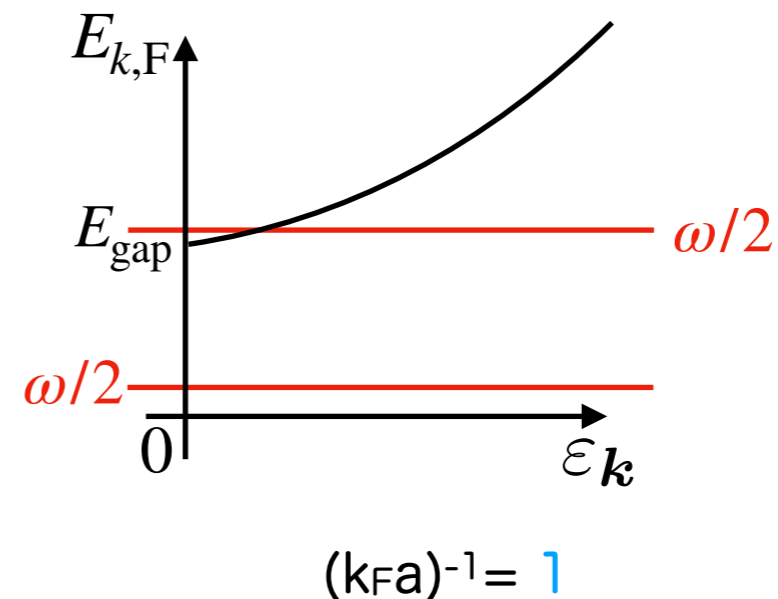
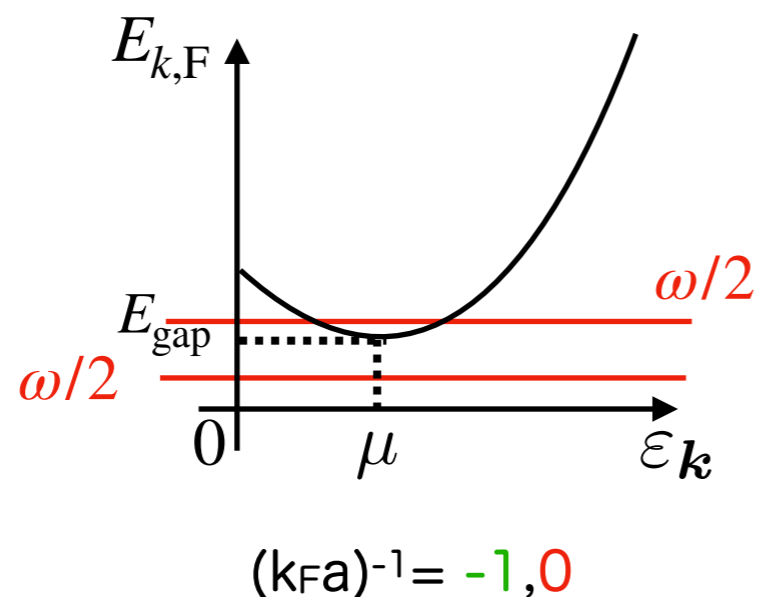
Result for a Fermi superfluid

YS, Tajima, & Uchino, arXiv:2103.02418



$$\text{Re} \sigma_{xx}^{(S)}(\omega) \propto \sum_{\mathbf{k}} k_x^2 \delta(|\omega| - 2E_{\mathbf{k},F})$$

1. Spin is insulated for small ω
2. Behaviors for $\omega \rightarrow 2E_{\text{gap}} + 0$
 - $\mu > 0$ [$(k_F a)^{-1} = -1, 0$] \rightarrow coherence peak
 - $\mu < 0$ [$(k_F a)^{-1} = 1$] \rightarrow decay



Topological Fermi superfluid

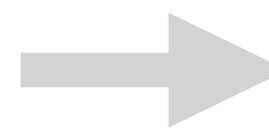
Tajima, YS, & Uchino, PRB (2022)

► Fermi atoms in quasi 1D

$$H = \sum_{k,\sigma} \xi_k c_{k,\sigma}^\dagger c_{k,\sigma} + V, \quad \xi_k = k^2/(2m) - \mu$$

► P-wave Feshbach resonance in the $\uparrow - \downarrow$ channel:

$$V = -U \sum_{k,k',q} k k' c_{k+q/2,\uparrow}^\dagger c_{-k+q/2,\downarrow}^\dagger c_{-k'+q/2,\downarrow} c_{k'+q/2,\uparrow}$$



Triplet pairing
 $S = 1, S_z = 0$

► BdG Hamiltonian

$$H_{\text{MF}} = \sum_k \Psi_k^\dagger H_{\text{BdG}}(k) \Psi_k \quad \Psi_k = \begin{pmatrix} c_{k,\uparrow} \\ c_{-k,\downarrow}^\dagger \end{pmatrix}$$

$$H_{\text{BdG}}(k) = \boldsymbol{\sigma} \cdot \mathbf{R}(k) = -\sigma_x \Delta(k) + \sigma_z \xi_k,$$

$$\Delta(k) = kD \quad (D > 0)$$

Class BDI with winding # $\nu \in \mathbb{Z}$

		TRS	PHS	SLS	$d=1$
Standard (Wigner-Dyson)	A (unitary)	0	0	0	-
	AI (orthogonal)	+1	0	0	-
	AII (symplectic)	-1	0	0	-
Chiral (sublattice)	AIII (chiral unitary)	0	0	1	\mathbb{Z}
	BDI (chiral orthogonal)	+1	+1	1	\mathbb{Z}
	CII (chiral symplectic)	-1	-1	1	\mathbb{Z}
BdG	D	0	+1	0	\mathbb{Z}_2
	C	0	-1	0	-
	DIII	-1	+1	1	\mathbb{Z}_2
	CI	+1	-1	1	-

Schnyder, Ryu, Furusaki & Ludwig, PRB (2008)

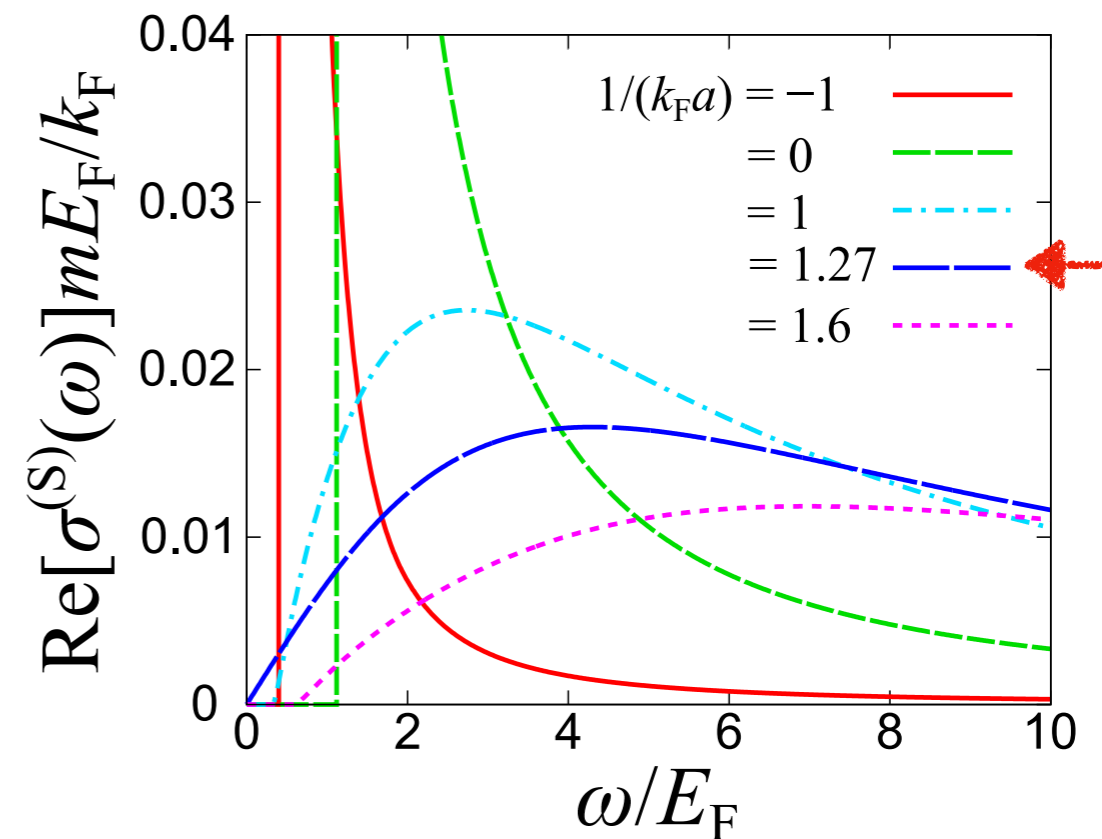
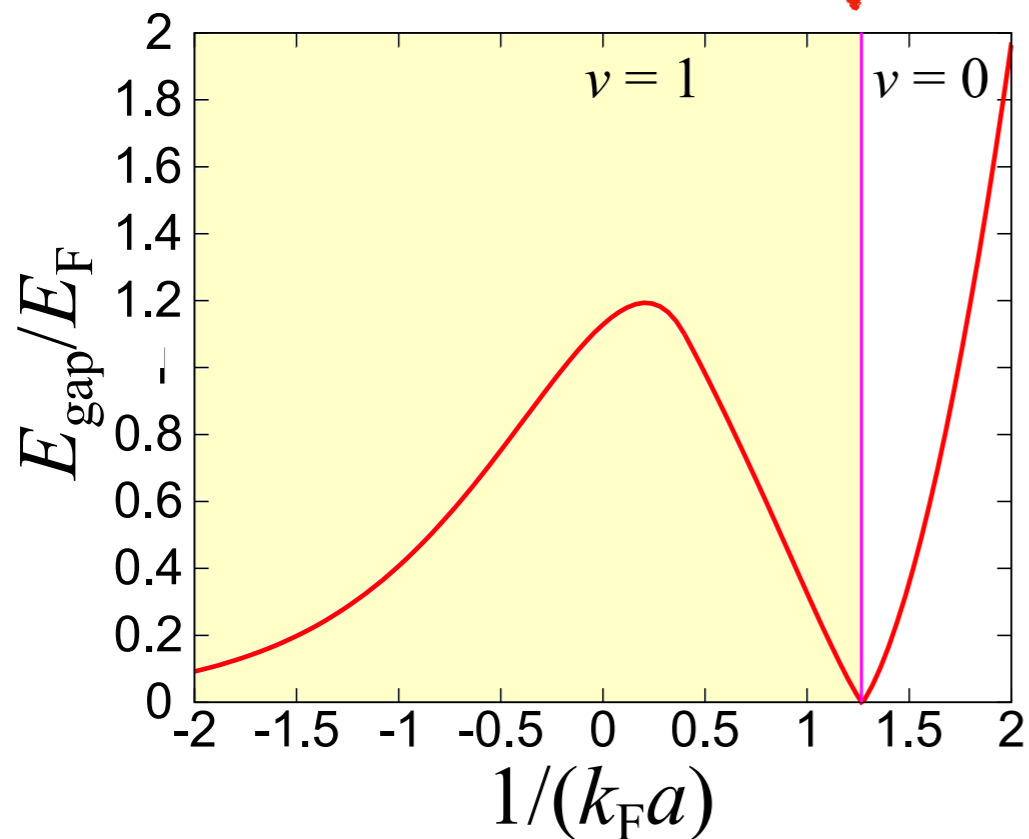
Spectrum of spin conductivity

Tajima, YS, & Uchino, PRB (2022)

Topological phase transition



Closing of the spectral gap

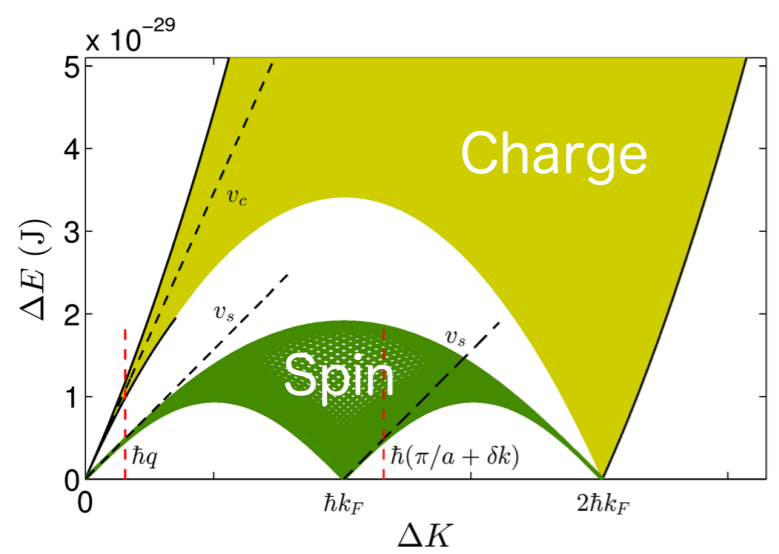
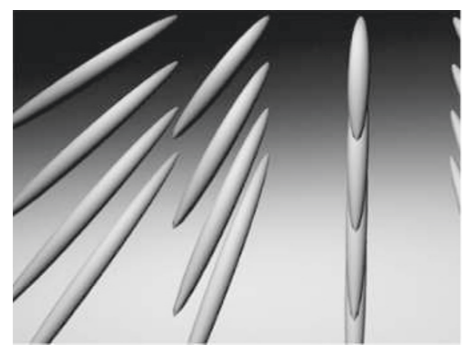


$$\text{Re} \sigma^{(S)}(\omega) \propto \sum_k k^2 \delta(|\omega| - 2E_k)$$

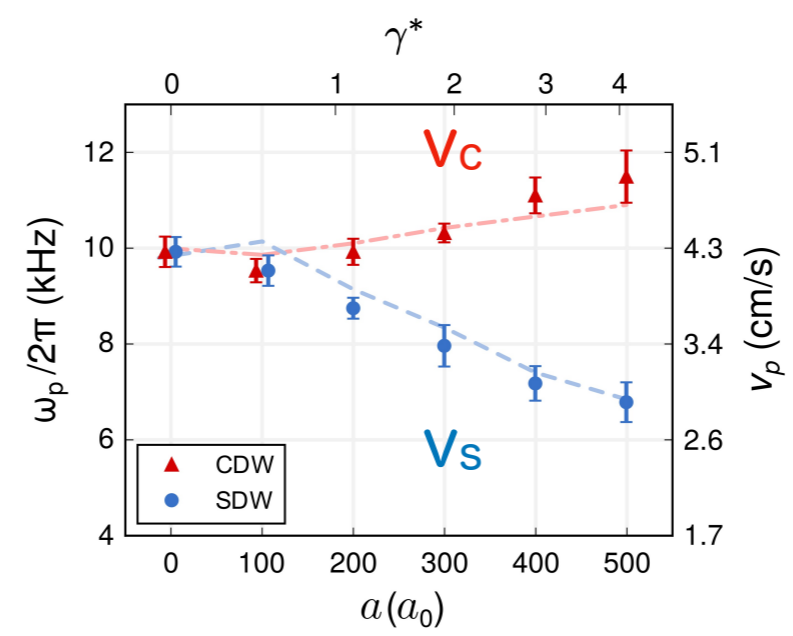
Tomonaga-Luttinger liquid

► One-dimensional systems with spin-charge separation

Described by 4 parameters v_c, v_s, K_c, K_s



He et al., PRL (2020)



Measurement of v_c & v_s
 Senaratne et al. (Rice), to be published in Science (2022)

K_s can be experimentally determined by spin conductivity at low frequency

YS, Tajima, & Uchino, arXiv:2103.02418

$$\text{Re } \sigma^{(S)}(\omega) \propto \omega^{4K_S - 5}$$

(Memory function method)

cf. Charge conductivity

Giamarchi, PRB (1991); (1992)

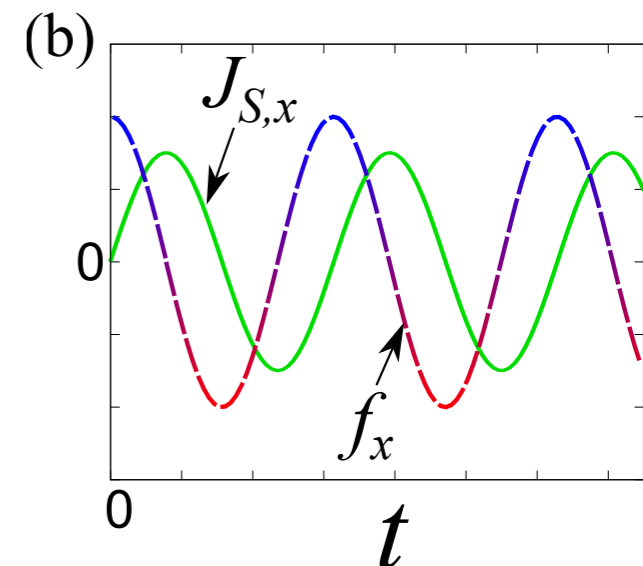
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Summary of this talk

- ▶ Optical (AC) spin conductivity
Measurable in cold-atom experiments

$$\sigma_{\alpha,\beta}^{(S)}(\omega) = \tilde{J}_{S,\alpha}(\omega) / \tilde{f}_{\beta}(\omega) \quad (\alpha, \beta = x, y, z)$$



- ▶ Significance of $\sigma_{\alpha\beta}^{(S)}(\omega)$
 - YS, Tajima, & Uchino, accepted by PRResearch (2022)
 - Tajima, YS, & Uchino, PRB (2022)

1. Elusive in solid-state systems
2. Powerful probe for quantum many-body states
BCS-BEC crossover, Tomonaga-Luttinger liquid(TLL), Spinor BEC,...
3. Widely applicable probe for clean systems

- ▶ Future perspective: Pseudogap of the unitary Fermi gas ?

Backup Slides

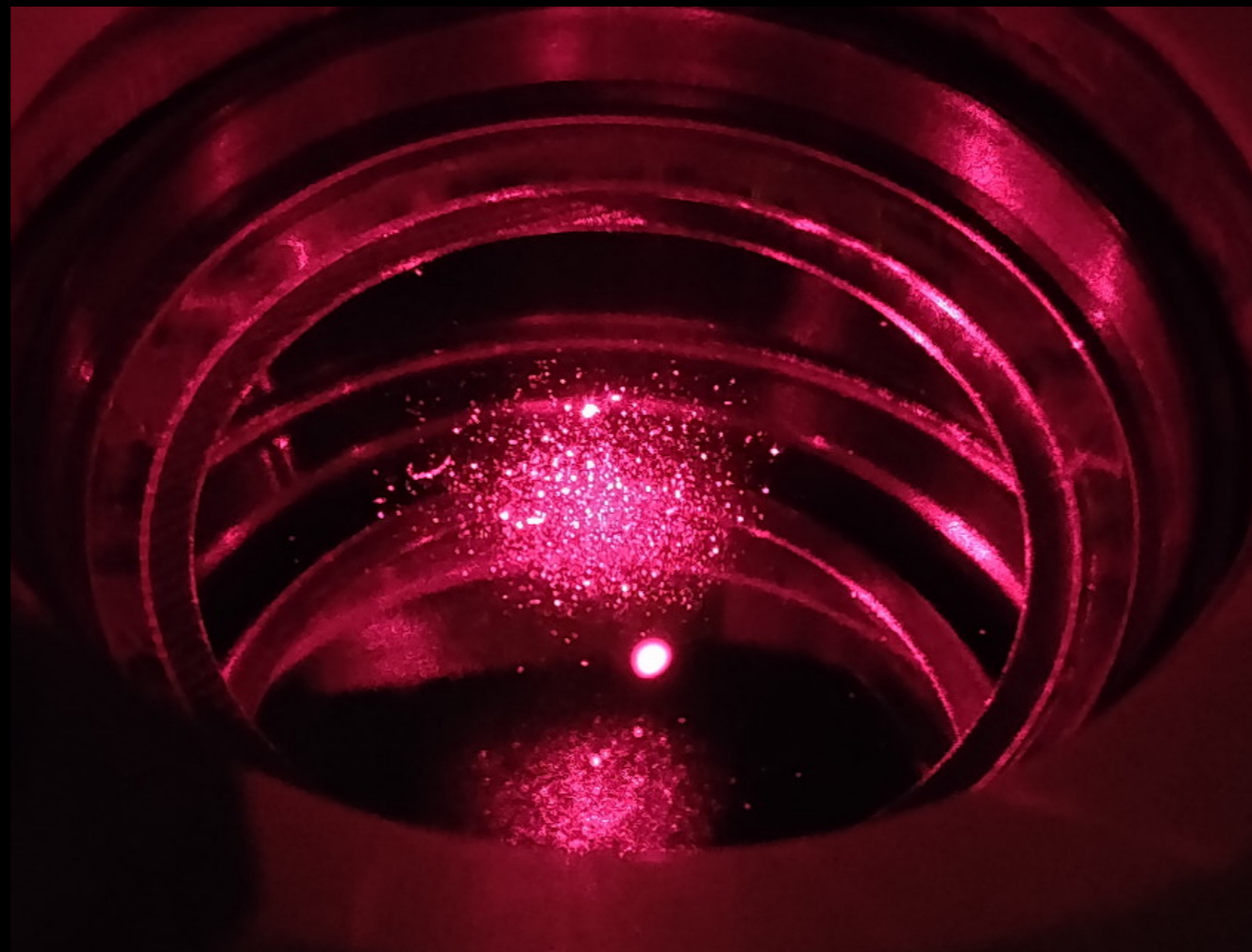
Ultracold atoms

► Very pure & highly controllable atomic gases

$$n = 10^{13} - 10^{15} \text{ cm}^{-3}, T = 10^{-6} - 10^{-8} \text{ K} (\mu\text{K} - \text{nK})$$

Coldest in the Universe !!

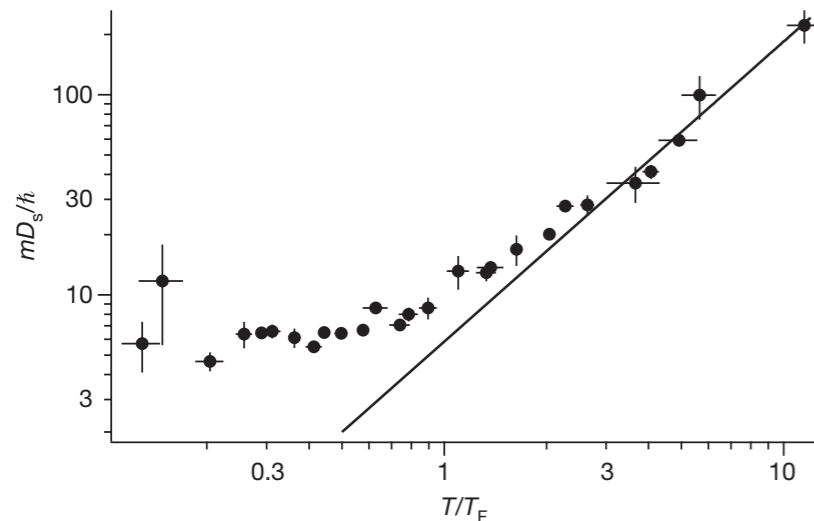
(cf. O₂ in a room: $n = 10^{19} \text{ cm}^{-3}$, $T = 10^3 \text{ K}$)



Transport phenomena

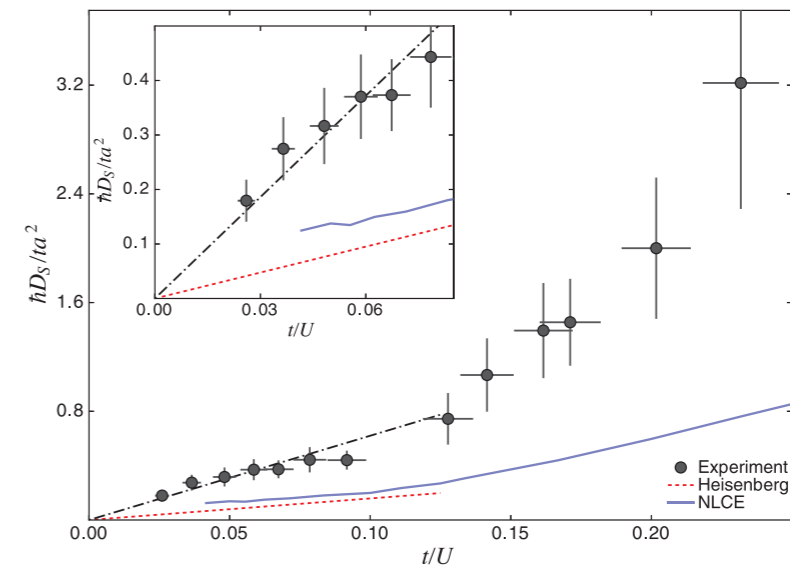
1. Spin diffusivity

Sommer et al.,(MIT) Nature (2011); ...



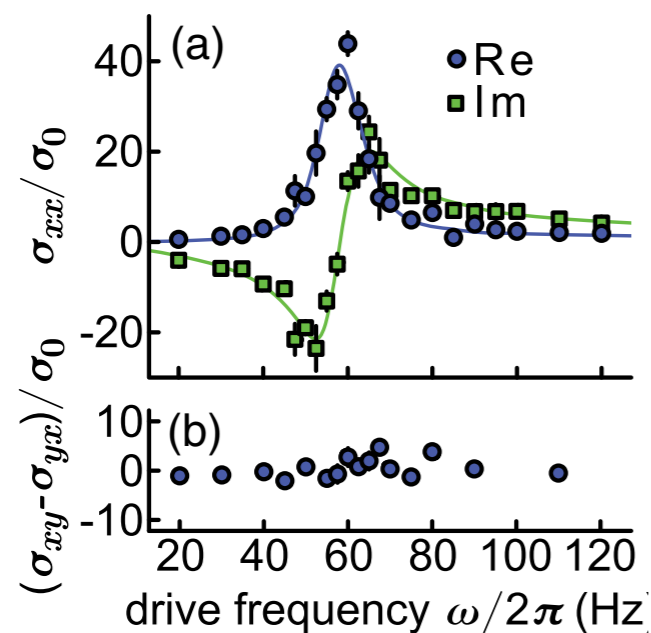
2. DC spin conductivity

Nicholos et al.,(MIT) Nature (2019); ...



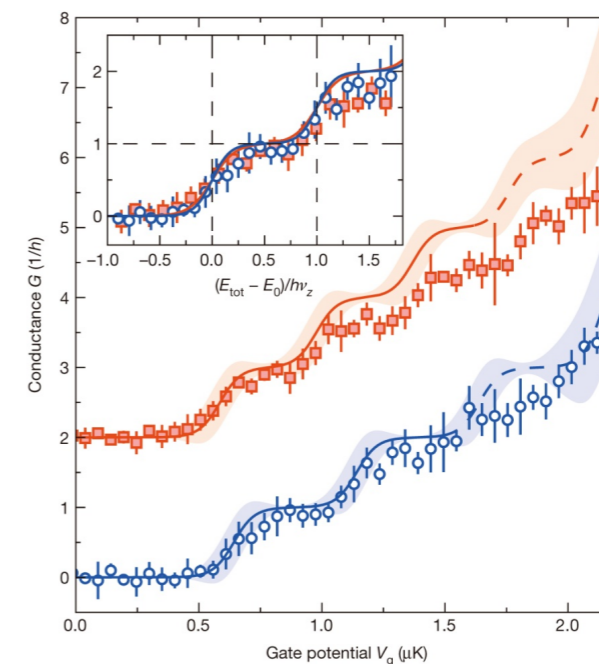
3. Optical (AC) conductivity

Anderson et. al.,[Toronto group] PRL (2019)



4. Quantized conductance

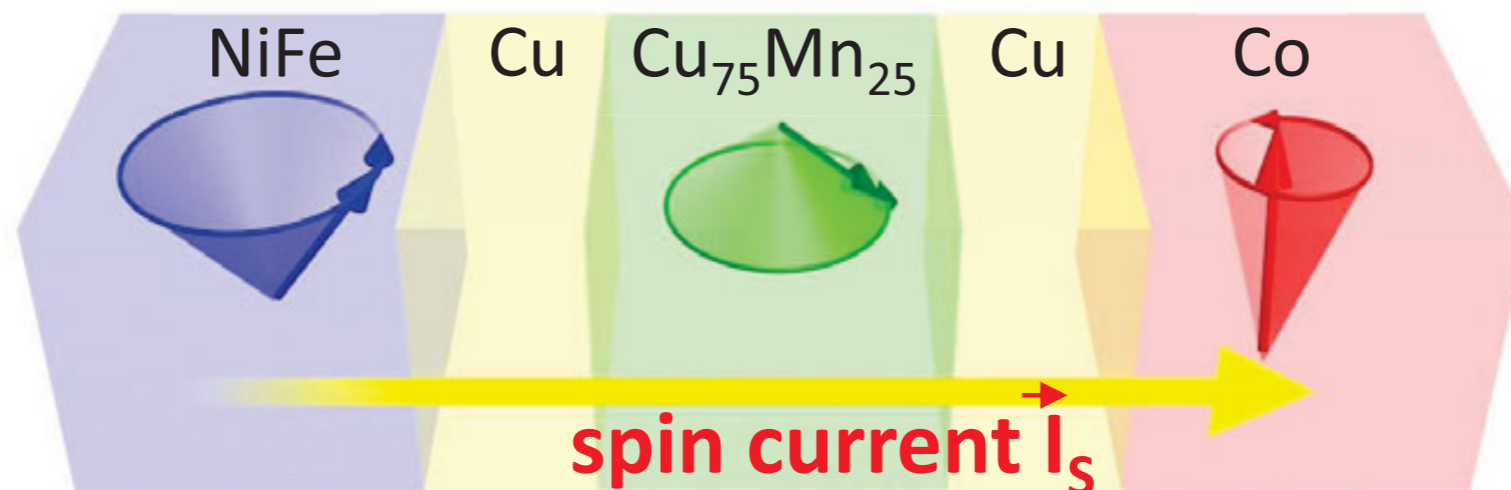
Krinner et al.[EHT], Nature(2015)



1. AC spin transport

► Hot topic in spintronics

Multilayer systems are considered



Ferromagnetic resonance in Multilayer systems [Li et al., PRL \(2016\)](#)

Bulk transport property $\sigma_{\alpha\beta}^{(S)}(\omega)$ is elusive !!

AC spin transport within bulk is accessible with cold atoms

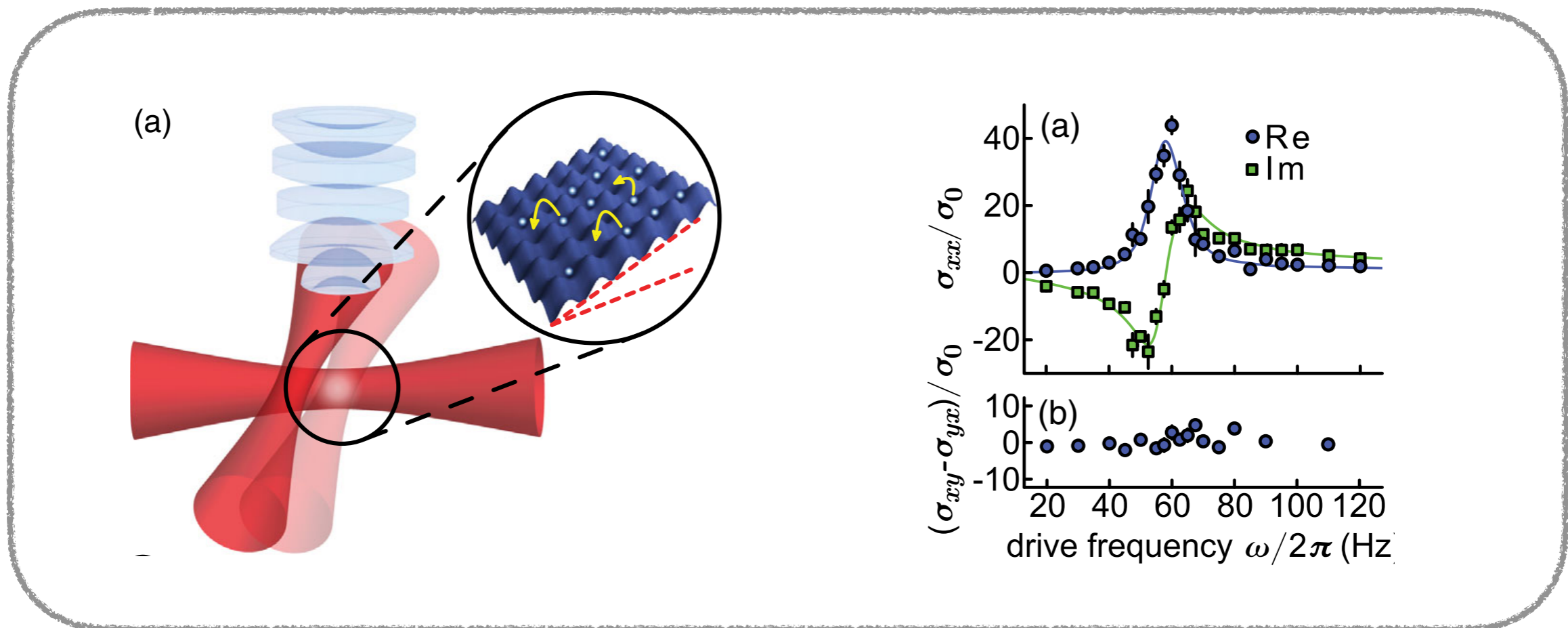
3. Probe for clean gases

► Optical mass conductivity for cold atoms

Proposal: Tokuno & Giamarchi, PRL (2011)

Wu, Taylor, & Zaremba, EPR (2015)

Experiment: Anderson et. al., [Toronto group] PRL (2019)

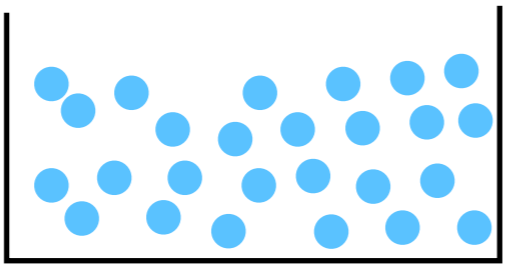
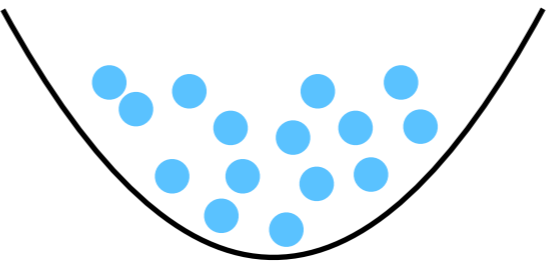
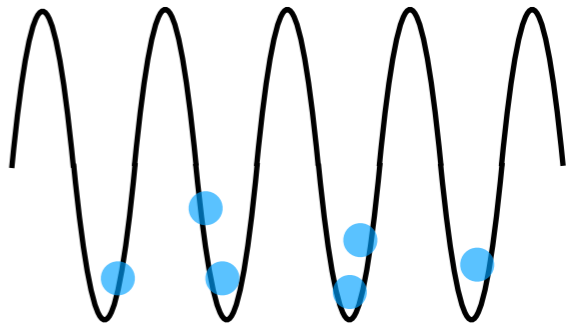


Optical lattice is essential !!!

3. Probe for clean gases

► Generalized Kohn's theorem: [Kohn \(1961\)](#); [Brey et al. \(1989\)](#); [Li et al. \(1991\)](#)

Strong constraint on mass conductivity of clean gases

	Homogeneous	Harmonically trapped	On lattice
Mass conductivity	 $\frac{N}{m} \delta(\omega)$	 $\frac{N}{m} \delta(\omega - \omega_{\text{trap}})$	 <p>Sensitive to quantum states</p>

Spin conductivity is never constrained and works as probes w/o lattice

How to extract spin conductivity

YS, Tajima, & Uchino, arXiv:2103.02418

1. Spin current: $\langle \mathbf{J}_S(t) \rangle = \frac{d}{dt} \left\langle \int d\mathbf{r} \mathbf{r} S_z(\mathbf{r}, t) \right\rangle \equiv \frac{d}{dt} \langle \mathbf{X}_S(t) \rangle$ (Spin conservation)
2. Spin conductivity: $\langle \tilde{\mathbf{J}}_{S,\alpha}(\omega) \rangle = \sigma_{\alpha\beta}^{(S)}(\omega) \tilde{f}_\beta(\omega)$ (Ohm's law in frequency space)
3. Driving force: $f_\beta(t) = F_\beta \cos(\omega_0 t)$



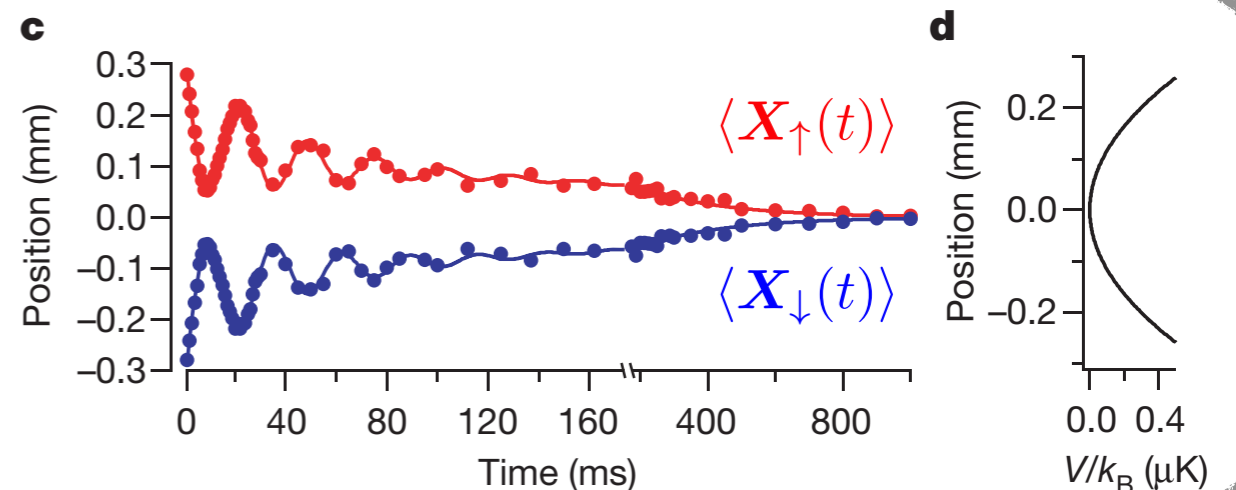
$$\frac{\langle \delta X_{S,\alpha}(t) \rangle}{F_\beta} = -\frac{\text{Im} \sigma_{\alpha\beta}^{(S)}(\omega_0)}{\omega_0} \cos(\omega_0 t) + \frac{\text{Re} \sigma_{\alpha\beta}^{(S)}(\omega_0)}{\omega_0} \sin(\omega_0 t)$$

Measurement of $\langle X_{S,\alpha}(t) \rangle \longrightarrow \sigma_{\alpha\beta}^{(S)}(\omega = \omega_0)$

e.g. Experiment on spin diffusion (w/o $f_\beta(t)$)

$$\langle \mathbf{X}_S(t) \rangle = \langle \mathbf{X}_\uparrow(t) \rangle - \langle \mathbf{X}_\downarrow(t) \rangle$$

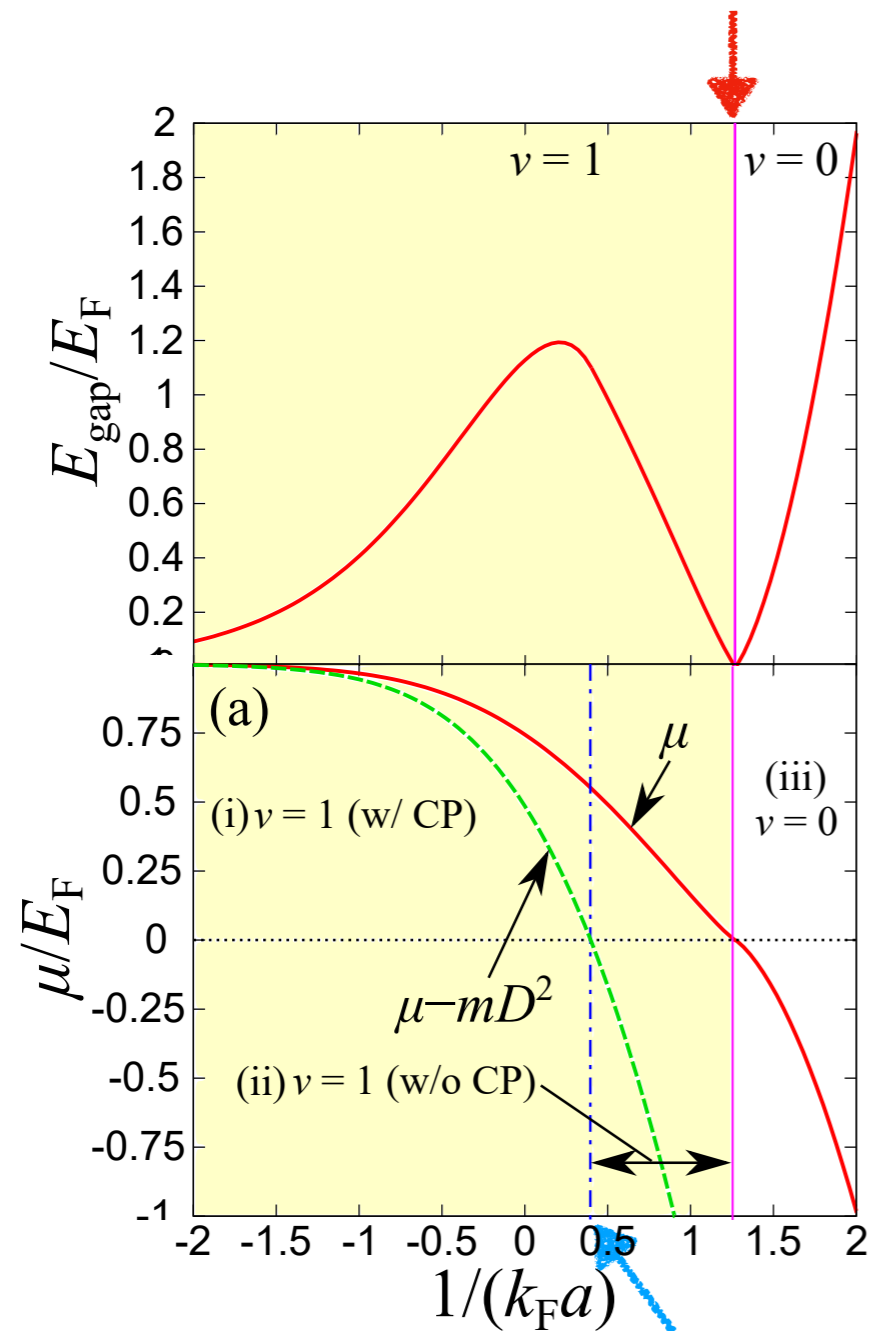
Sommer et al.,(MIT) Nature (2011)



Topological phase transition

Tajima, YS, & Uchino, PRB (2022)

Topological phase transition

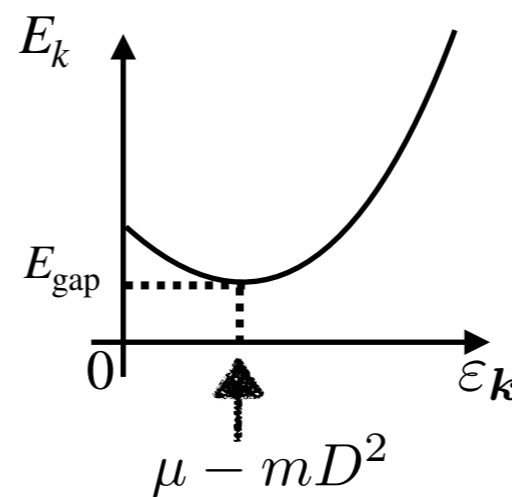


Coherence Peak (CP) disappears

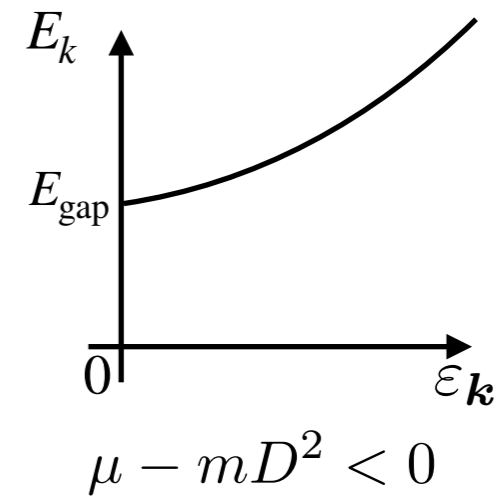
Topological phase transition at $\mu = 0$

$$\nu = 1 \rightarrow \nu = 0$$

Quasiparticle energy $E_k = \sqrt{(\varepsilon_k - \mu)^2 + D^2 k^2}$



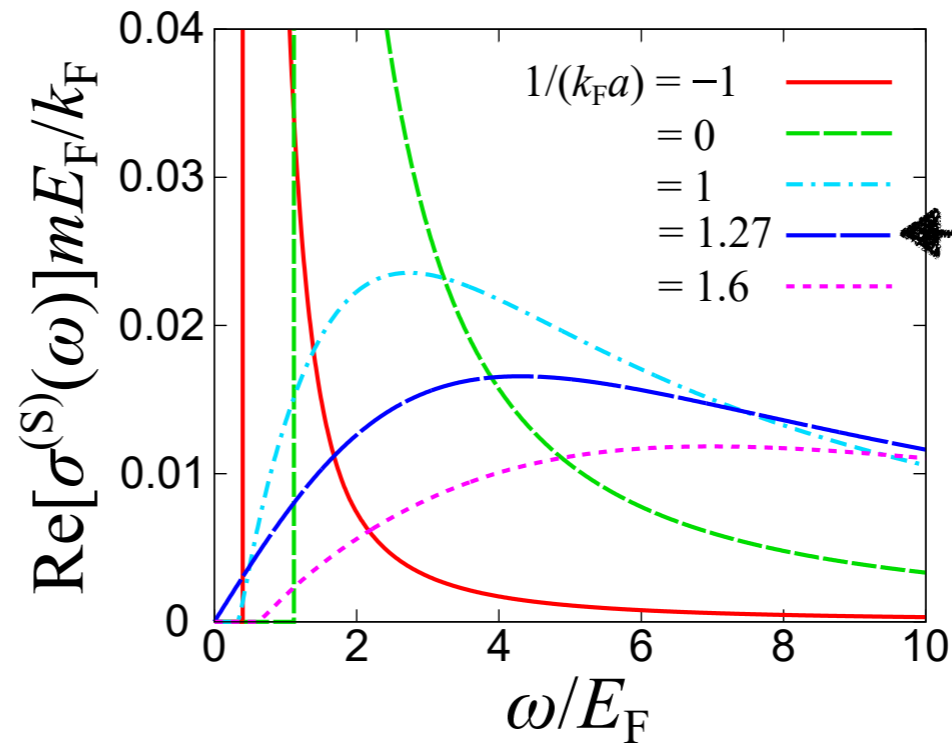
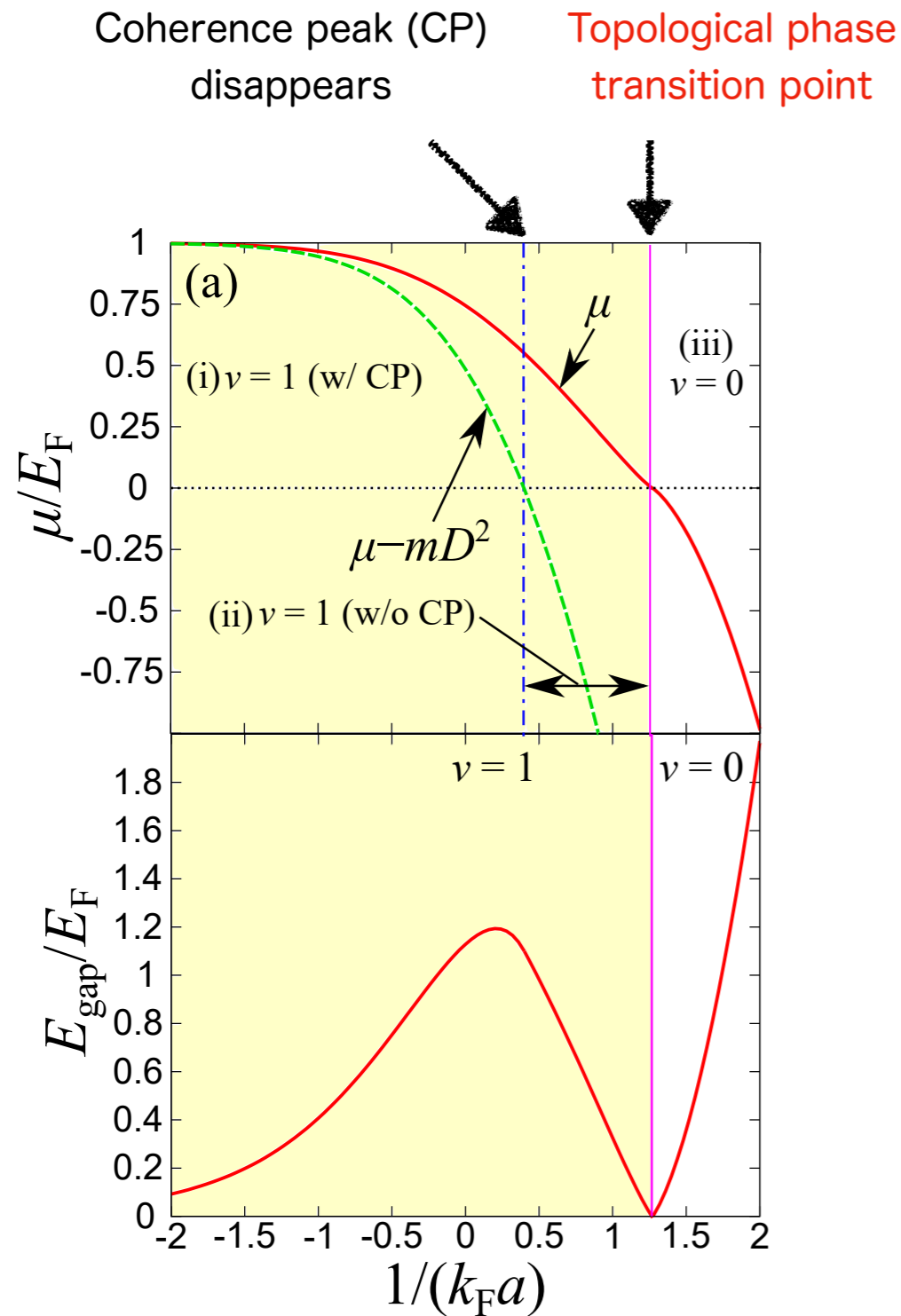
w/ CP



w/o CP

Spectrum of spin conductivity

Tajima, YS, & Uchino, PRB (2022)



$$\text{Re} \sigma^{(S)}(\omega) \propto \sum_k k^2 \delta(|\omega| - 2E_k)$$

Closing of the spectral gap of $\text{Re} \sigma^{(S)}(\omega)$



Topological phase transition