

Properties of a driven-dissipative non-equilibrium Fermi superfluid

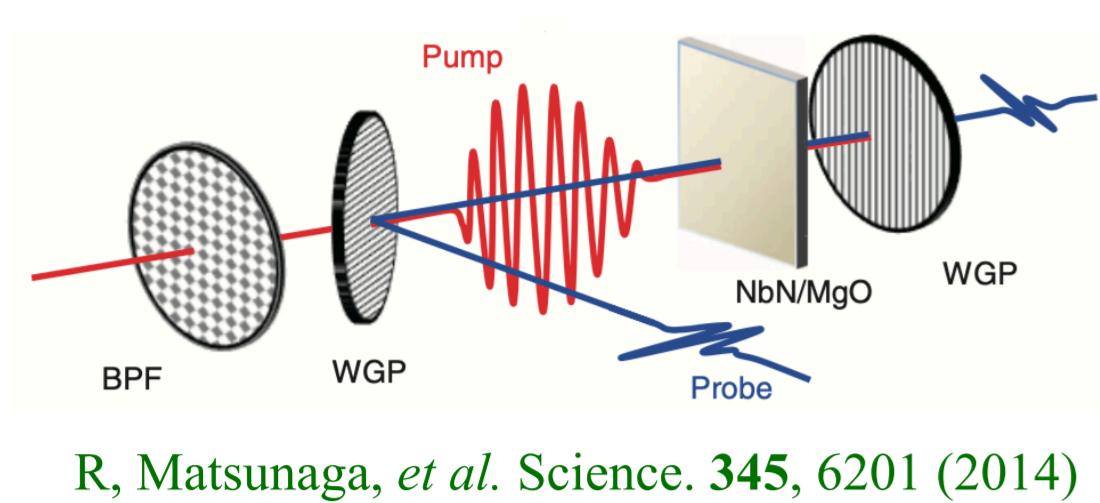
Taira Kawamura¹, Ryo Hanai² and Yoji Ohashi¹

¹Department of Physics, Keio University

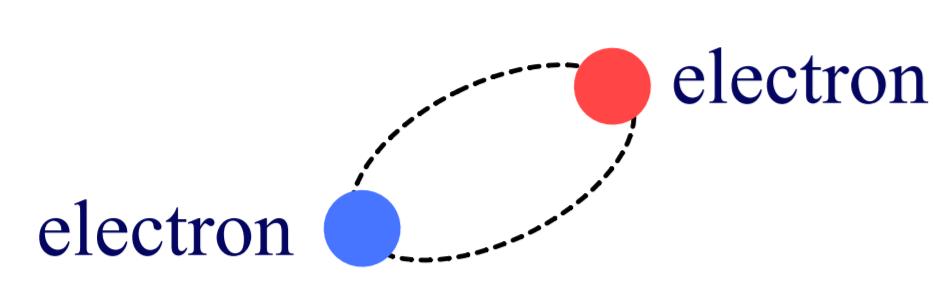
²James Franck Institute and Department of Physics, University of Chicago

Introduction *non-equilibrium Fermi condensates*

✓ non-equilibrium superconductivity

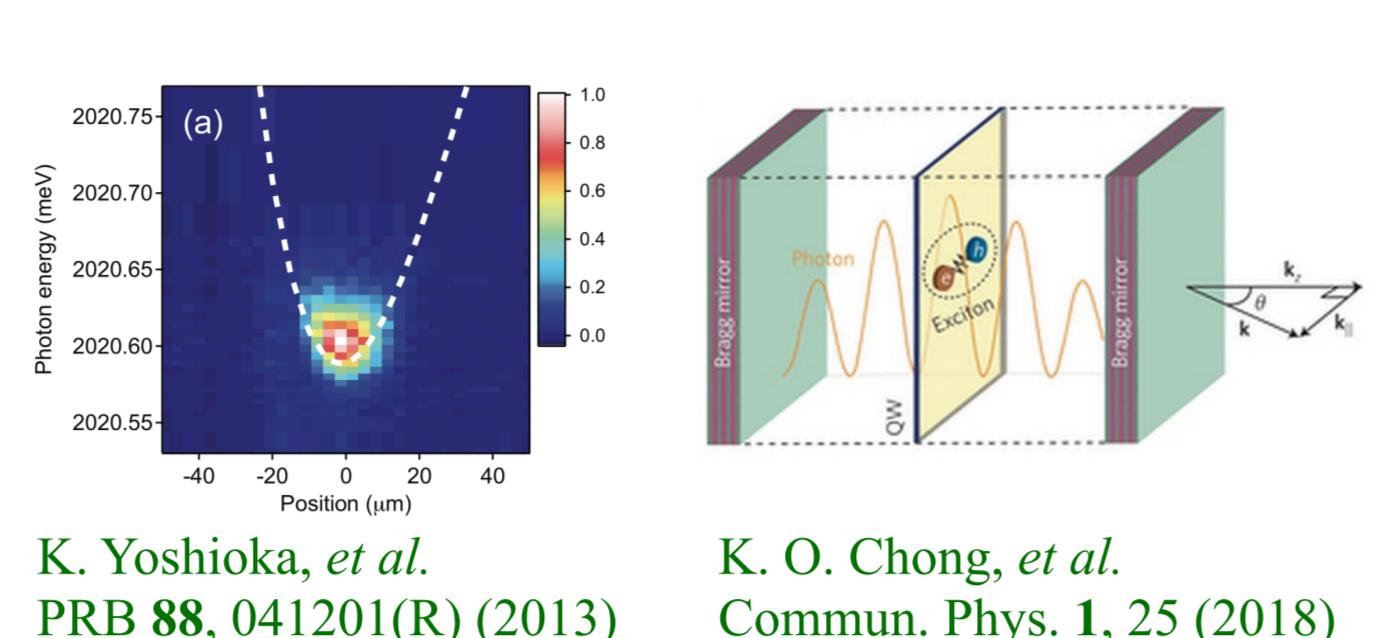


R. Matsunaga, et al. Science. **345**, 6201 (2014)



pair formation out of equilibrium \Rightarrow non-equilibrium + many body effect

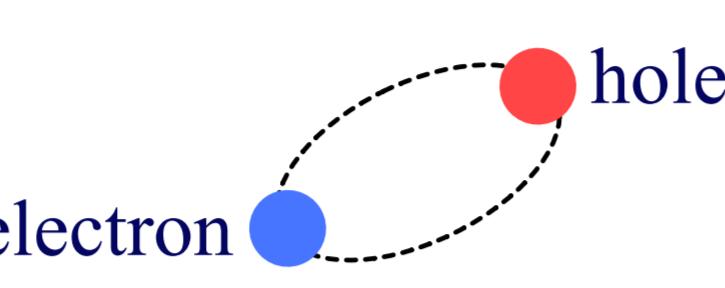
✓ exciton-polariton condensates



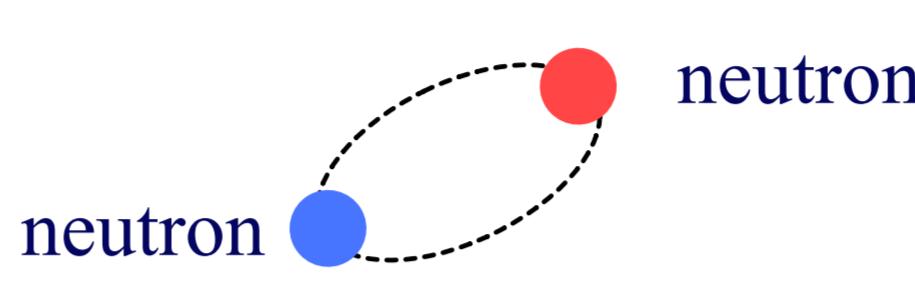
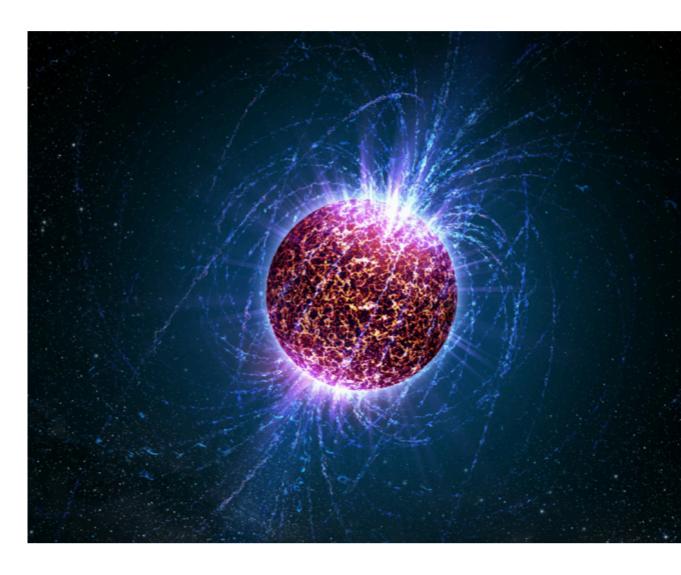
K. Yoshioka, et al. PRB **88**, 041201(R) (2013)

K. O. Chong, et al. Commun. Phys. **1**, 25 (2018)

<https://apatrano.wordpress.com/neutron-stars/>



✓ neutron star



Summary

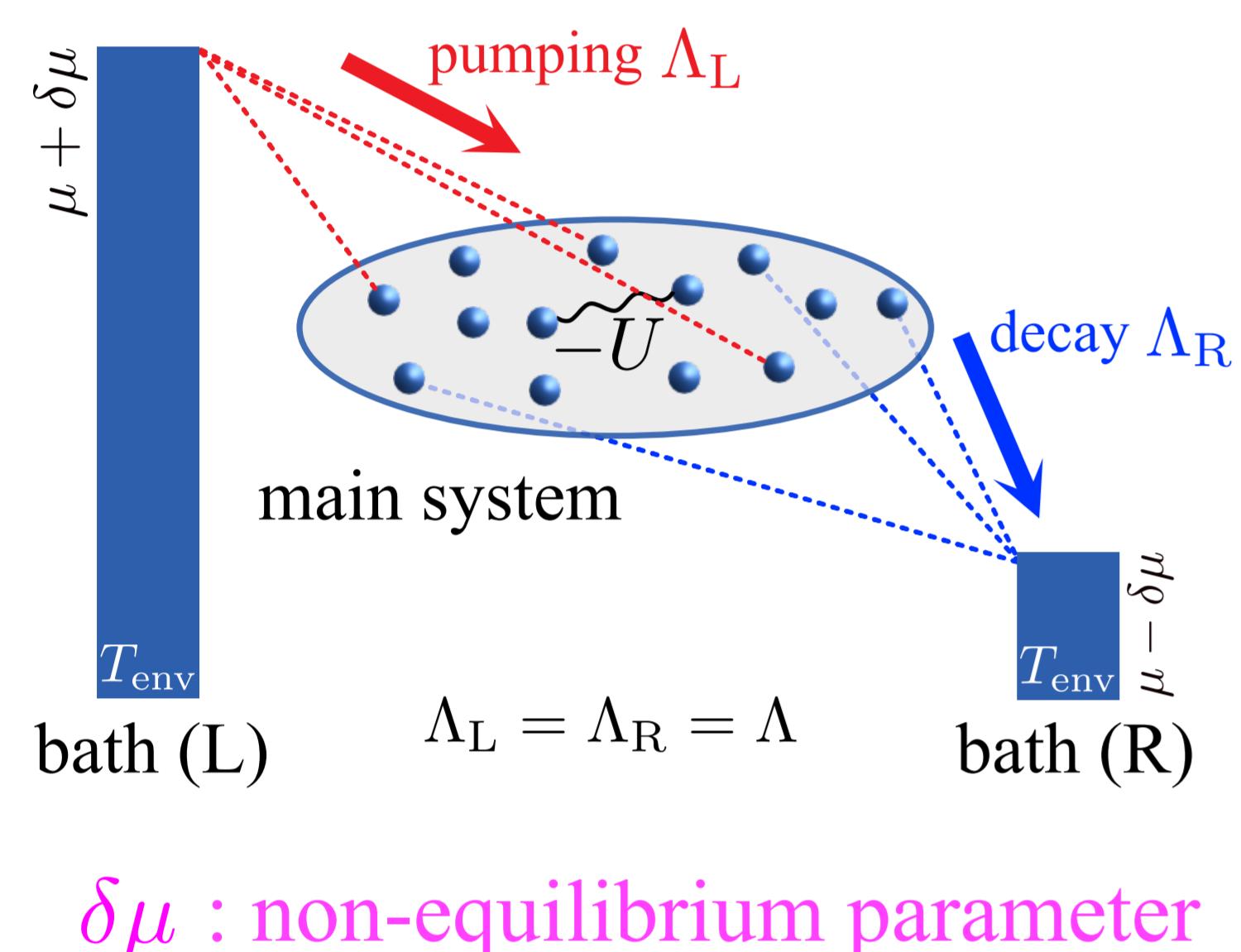
► We studied superfluid states in a driven-dissipative Fermi gas. By using Kadanoff-Baym equation, we derived a quantum kinetic equation, which includes both non-equilibrium and many-body effects. The fixed point of this kinetic equation gives non-equilibrium steady states.

► We found that two different states can be realized as non-equilibrium steady states. One of them is a gapless superfluid state, which is not found in a thermal equilibrium Fermi gas. However, this gapless state turned out to be unstable against superfluid fluctuations from the linear stability analysis.

► We also found that non-uniform Fulde-Ferrell like state can be realized by the non-equilibrium effects. The stability analysis of this state is our future work.

Formalism

✓ model



$$\mathcal{H}_{\text{tot}} = \mathcal{H}_{\text{sys}} + \mathcal{H}_{\text{env}} + \mathcal{H}_{\text{mix}}$$

► main system

$$H_{\text{sys}} = \sum_{\sigma=\uparrow,\downarrow} \int d\mathbf{r} \psi_{\sigma}^{\dagger}(\mathbf{r}) \frac{-\nabla^2}{2m} \psi_{\sigma}(\mathbf{r}) - \int d\mathbf{r} [\Delta(\mathbf{r}, t) \psi_{\uparrow}^{\dagger}(\mathbf{r}) \psi_{\downarrow}^{\dagger}(\mathbf{r}) + \text{H.c.}]$$

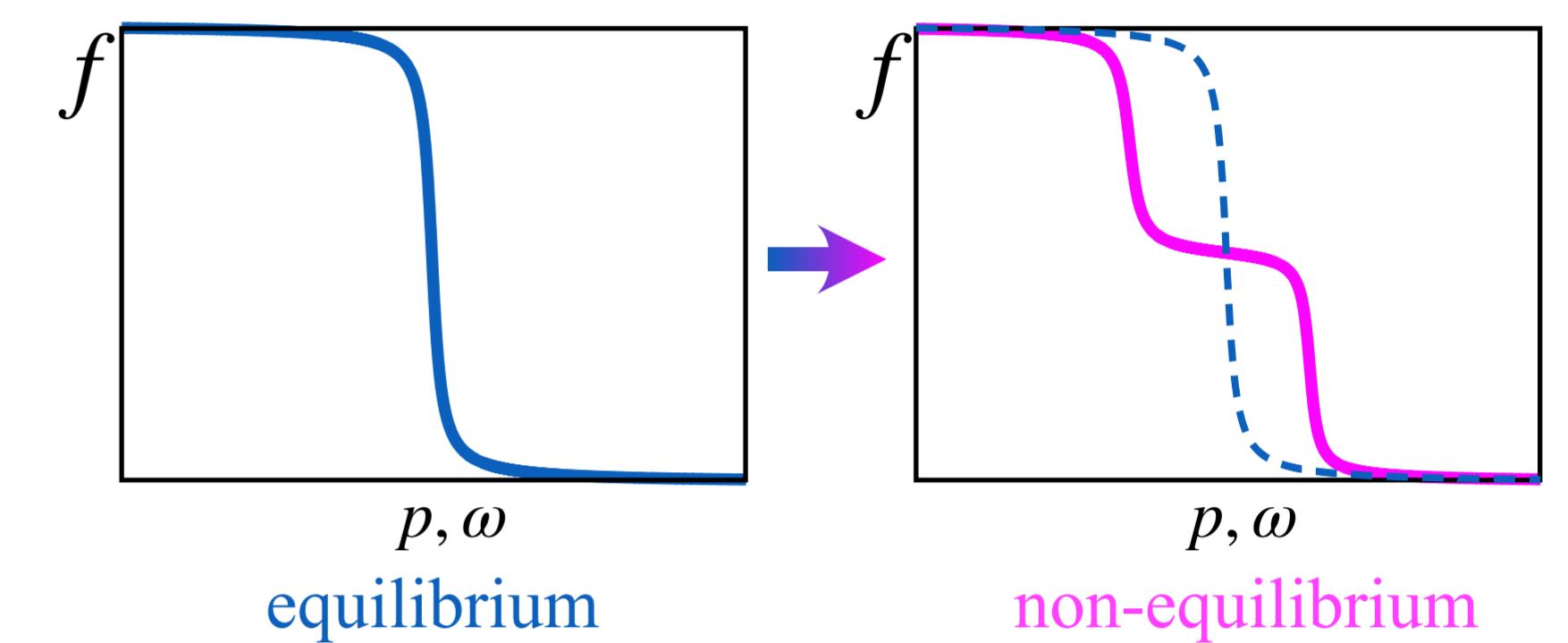
► environment system

$$H_{\text{env}} = \sum_{\alpha=L,R} \sum_{\sigma=\uparrow,\downarrow} \int d\mathbf{R} \phi_{\sigma}^{\alpha\dagger}(\mathbf{R}) \left[\frac{-\nabla^2}{2m} - \mu_{\alpha} \right] \phi_{\sigma}^{\alpha}(\mathbf{R})$$

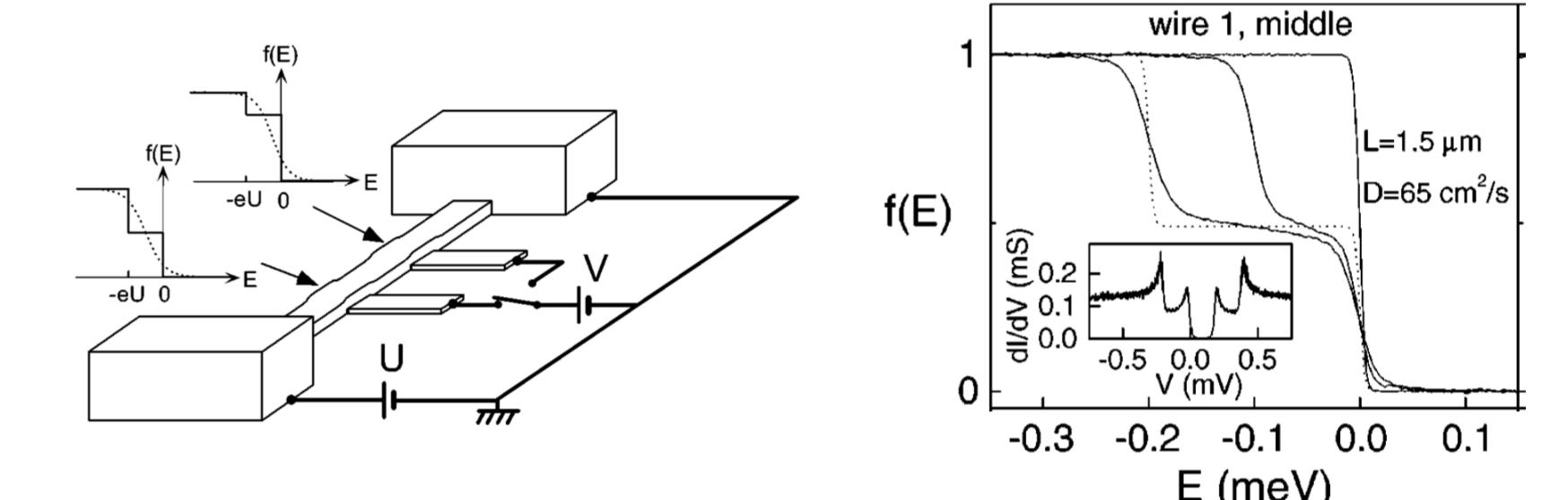
► mixing term

$$H_{\text{mix}} = \sum_{\alpha=L,R} \sum_{\sigma=\uparrow,\downarrow} \sum_{i=1}^{N_t} \int d\mathbf{r}_i \int d\mathbf{R}_i [e^{i\mu_{\alpha}t} \Lambda_{\alpha} \phi_{\sigma}^{\alpha\dagger}(\mathbf{R}_i) \psi_{\sigma}(\mathbf{r}_i) + \text{H.c.}]$$

Atomic momentum/energy distribution



(c.f.) electron energy distribution in a nanowire



H. Pothier, S. Guéron, O. Birge, D. Esteve, and M.H. Devoret, PRL **79**, 18 (1997)

✓ kinetic equation for the driven-dissipative Fermi gas

► Nambu lesser Green's function

$$g^{<}(\mathbf{r}_1 t_1, \mathbf{r}_2 t_2) = i \begin{pmatrix} \langle \psi_{\uparrow}^{\dagger}(\mathbf{r}_2, t_2) \psi_{\uparrow}(\mathbf{r}_1, t_1) \rangle & \langle \psi_{\downarrow}(\mathbf{r}_2, t_2) \psi_{\uparrow}(\mathbf{r}_1, t_1) \rangle \\ \langle \psi_{\uparrow}^{\dagger}(\mathbf{r}_2, t_2) \psi_{\downarrow}^{\dagger}(\mathbf{r}_1, t_1) \rangle & \langle \psi_{\downarrow}(\mathbf{r}_2, t_2) \psi_{\downarrow}^{\dagger}(\mathbf{r}_1, t_1) \rangle \end{pmatrix}$$

diagonal : particle density

off-diagonal : pair amplitude

► Kadanoff-Baym equation

$$[g^{-1} g^{<} - g^{<} g^{-1}] (\mathbf{p}, \mathbf{r}, t) = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} [\Sigma^{\text{R}} \otimes g^{<} - g^{<} \otimes \Sigma^{\text{A}} - g^{\text{R}} \otimes \Sigma^{<} + \Sigma^{<} \otimes g^{\text{A}}] (\mathbf{p}, \omega, \mathbf{r}, t)$$

driving term

collision term

self-energy

$$\Sigma = \begin{array}{c} \text{---} \\ \text{---} \end{array} \hat{\mathcal{G}} \begin{array}{c} \text{---} \\ \text{---} \end{array} -U$$

Hartree-Fock-Bogoliubov

interaction effects

$$+ \begin{array}{c} \text{---} \\ \text{---} \end{array} \Lambda \begin{array}{c} \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \end{array} \hat{D} \begin{array}{c} \text{---} \\ \text{---} \end{array} \Lambda$$

2nd Born

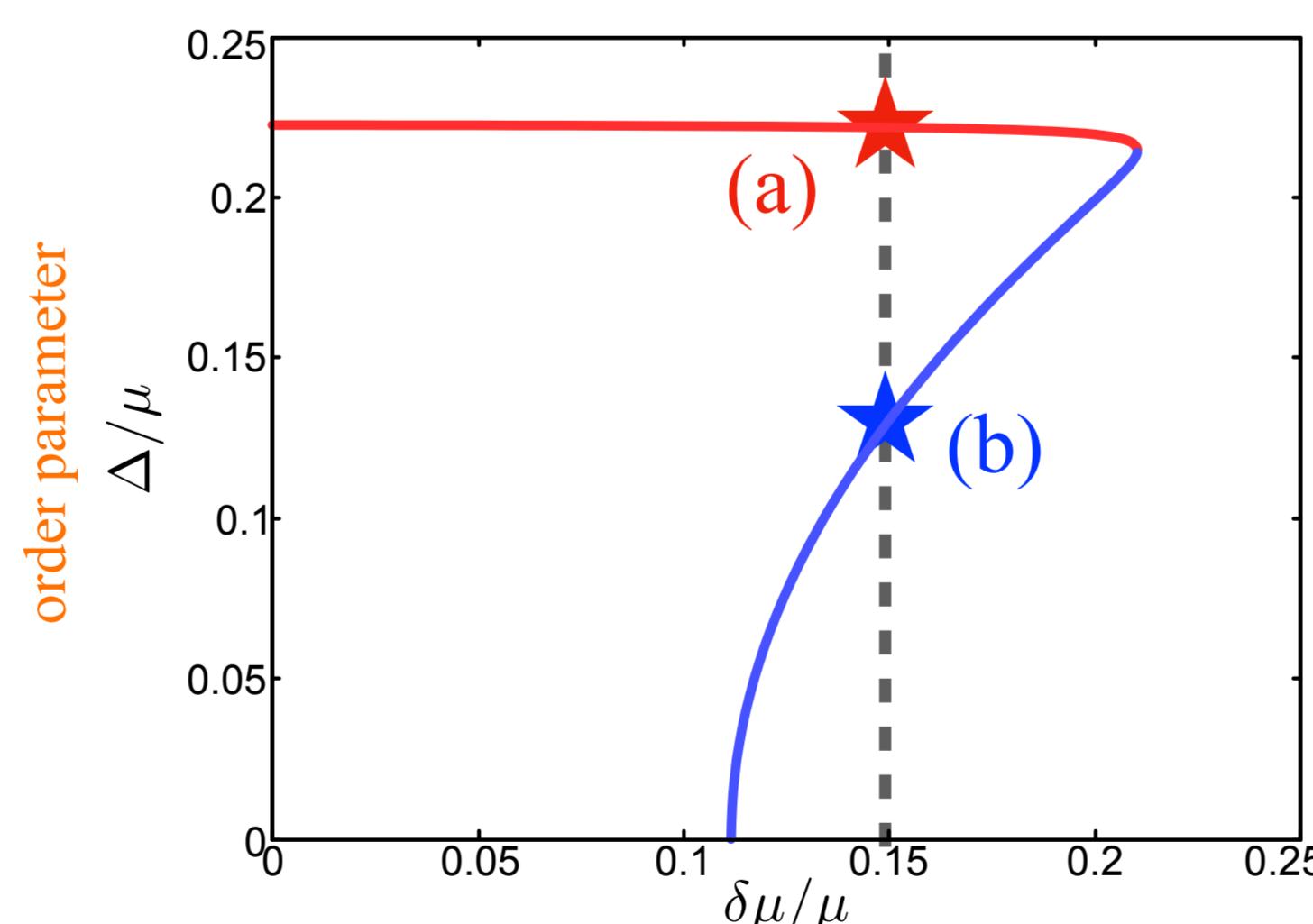
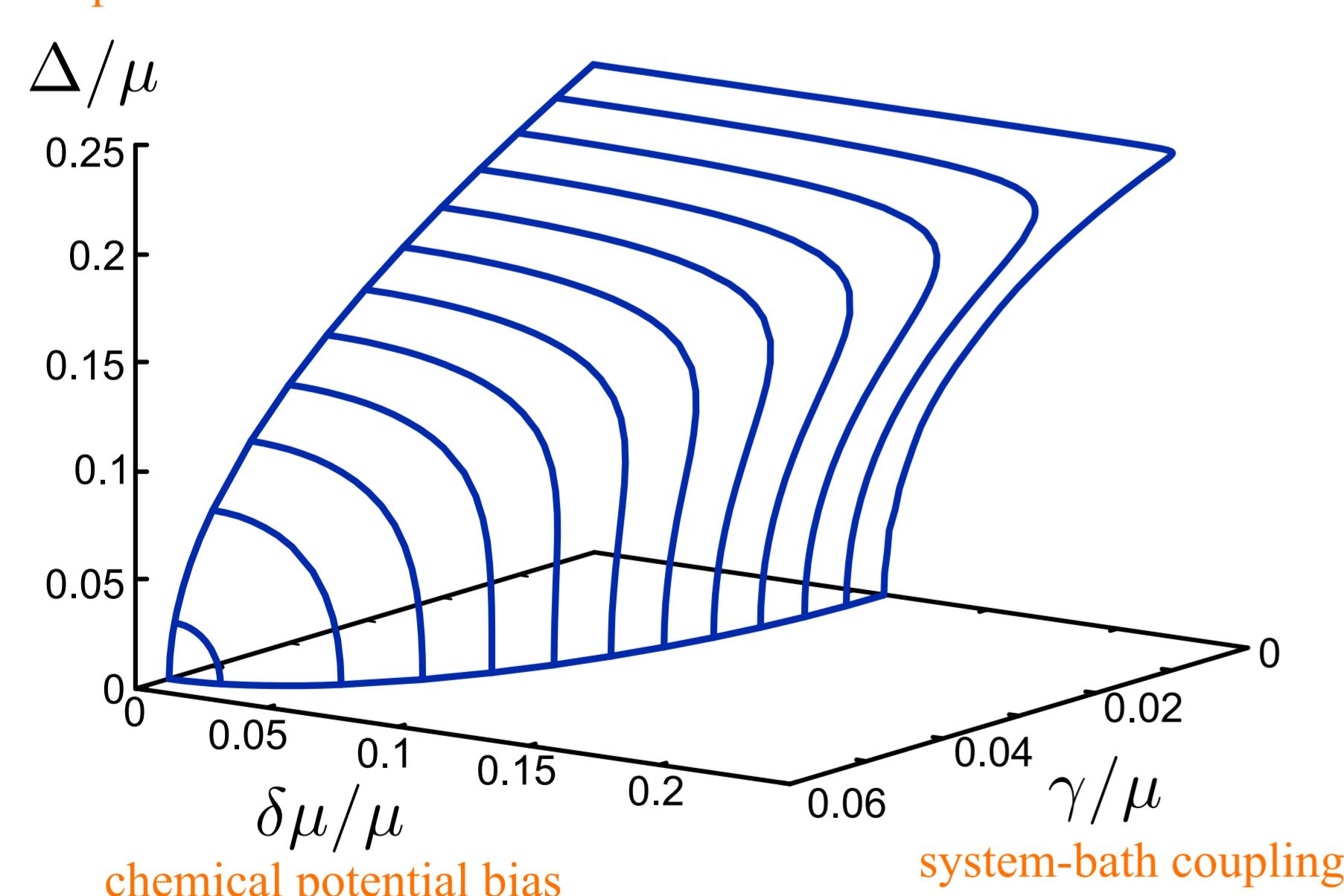
non-equilibrium steady state $\partial_t g^{<} = 0$

$$\Rightarrow \Delta_0 = -iU \sum_{\mathbf{p}} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} g^{<}(\mathbf{p}, \omega)_{12} \quad \text{non-equilibrium gap equation}$$

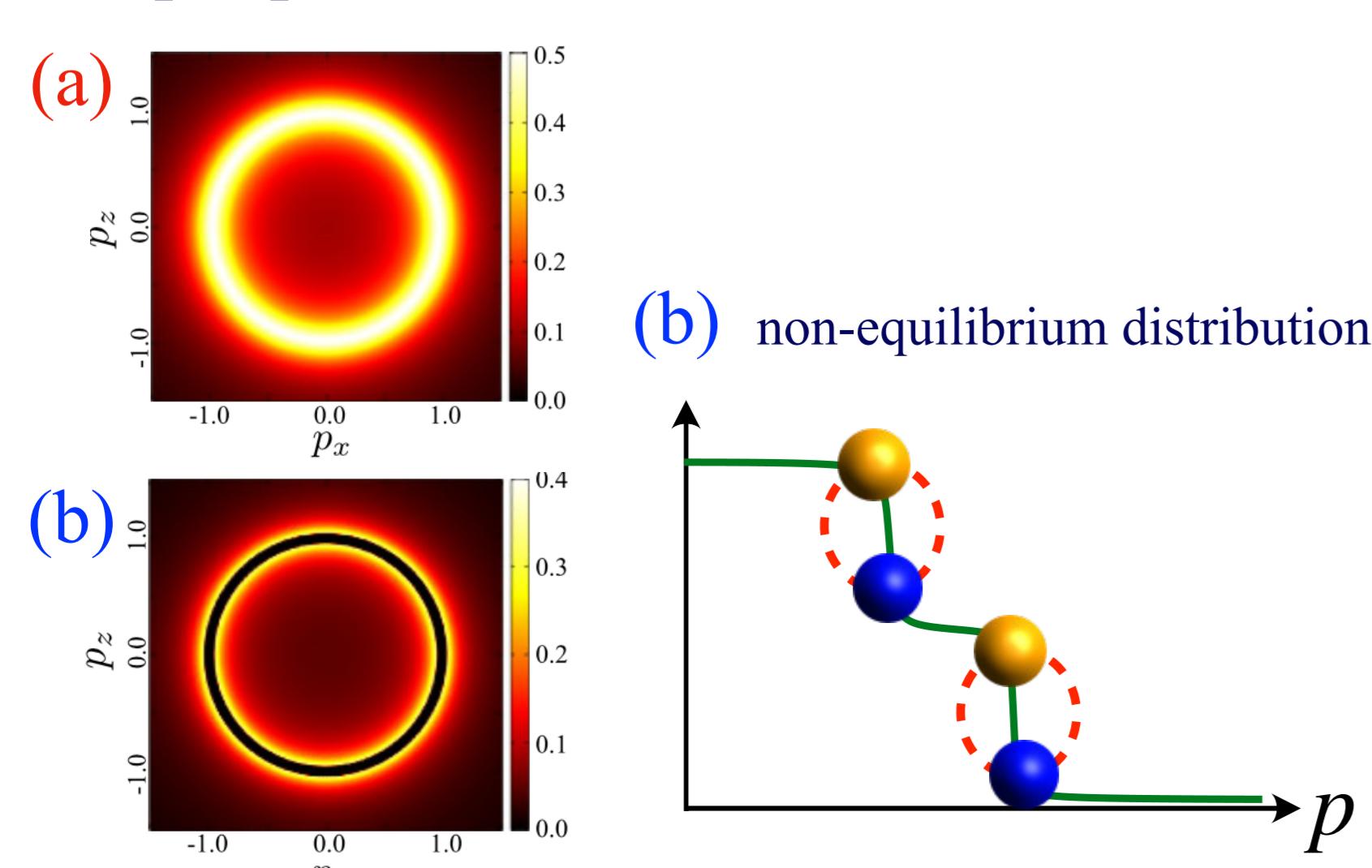
Result

✓ non-equilibrium steady states

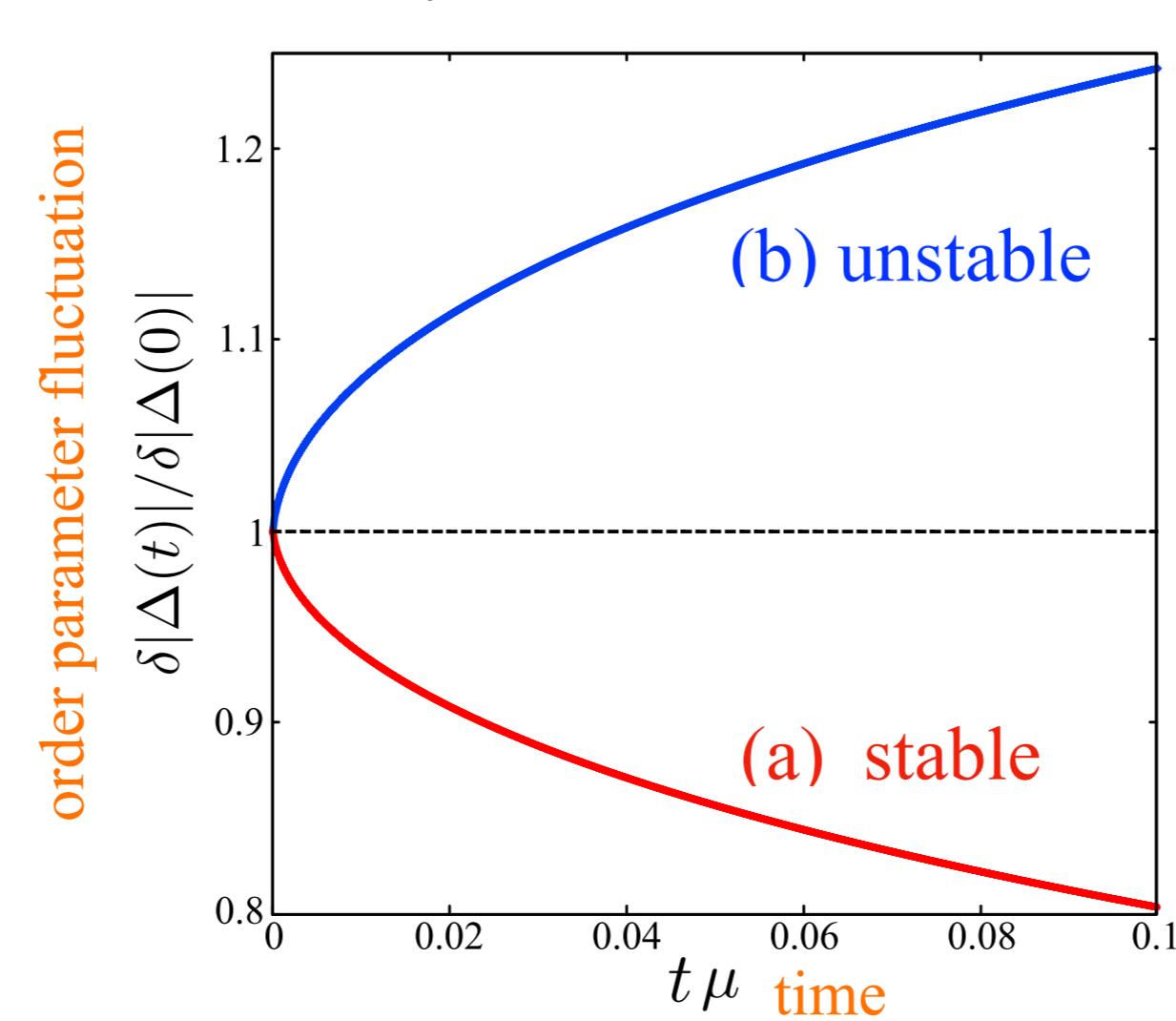
order parameter



Cooper pair wave function $\langle a_{-\mathbf{p}\downarrow} a_{\mathbf{p},\uparrow} \rangle$



linear stability $\delta\Delta(t) = \Delta(t) - \Delta_0$



✓ exotic pairing state

non-equilibrium T-matrix

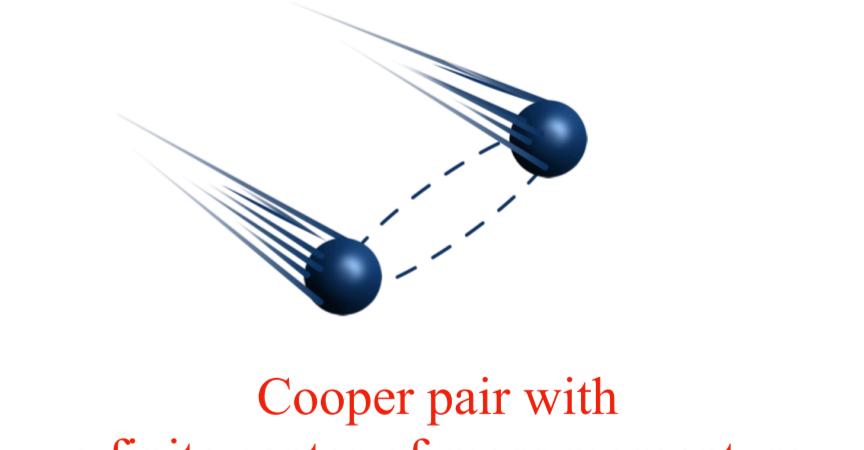
$$\Gamma = \begin{array}{c} \text{---} \\ \text{---} \end{array} -U \begin{array}{c} \text{---} \\ \text{---} \end{array} + \begin{array}{c} \text{---} \\ \text{---} \end{array} + \dots$$

non-equilibrium propagator including environment effects

Thouless criterion

$$[\Gamma^{\text{R}}(\mathbf{q} = 0, \nu = 0; T = T_c)]^{-1} = 0$$

$$[\Gamma^{\text{R}}(\mathbf{q} = \mathbf{Q}, \nu = 0; T = T_c)]^{-1} = 0$$



superfluid transition temperature

