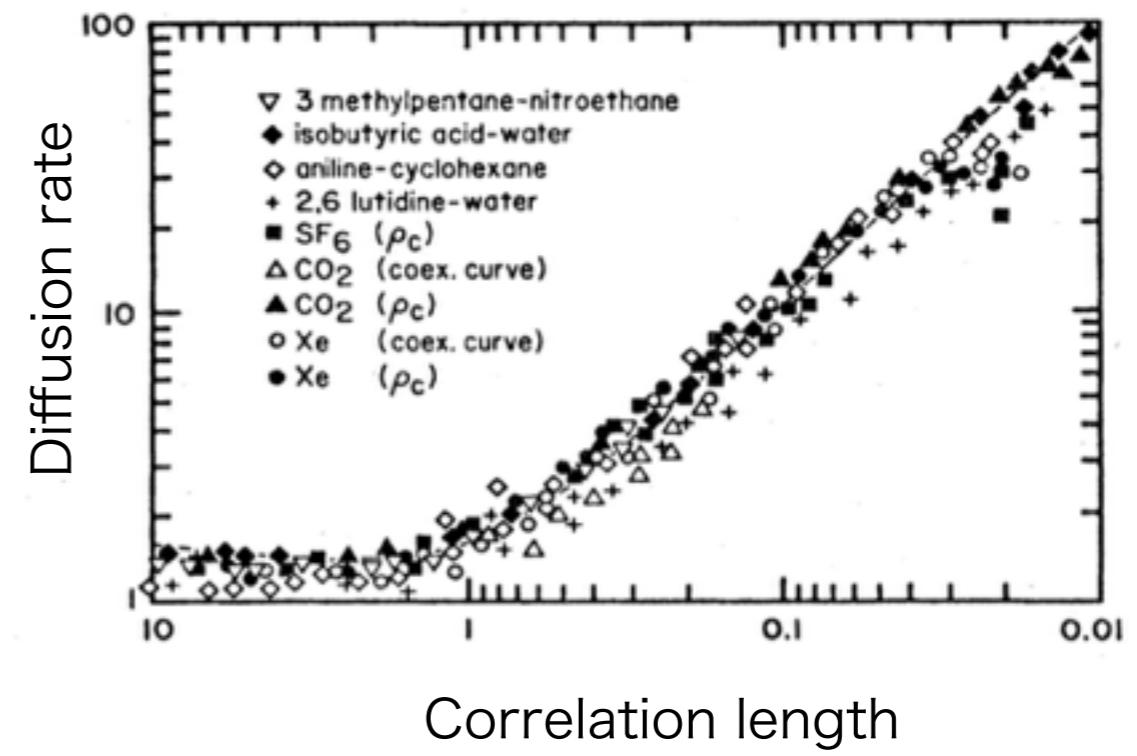


New dynamic critical phenomena of nonrelativistic photons in QCD

Noriyuki Sogabe (Keio University)

熱場の量子論 at YITP on September 4th, 2019
in a collaboration with Naoki Yamamoto (Keio University)

Critical phenomena



Correlation length

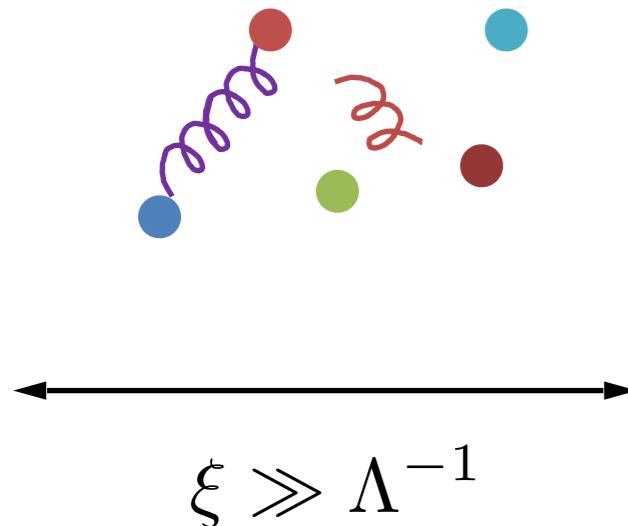
Swinney and Henry (1973)

from Brookhaven National Laboratory

Dynamic universality class

P. C. Hohenberg and B. I. Halperin (1977)

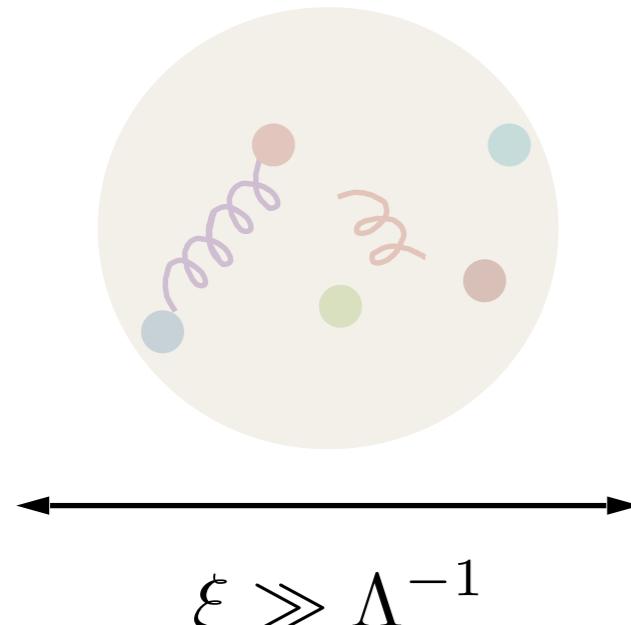
Microscopic theory



Dynamic universality class

P. C. Hohenberg and B. I. Halperin (1977)

Microscopic theory



Integrating out

Effective theory

Hydrodynamic variables:

- Order parameters
- Conserved charge densities
- Nambu-Goldstone modes

Same Symmetries

Dynamic universality class

P. C. Hohenberg and B. I. Halperin (1977)

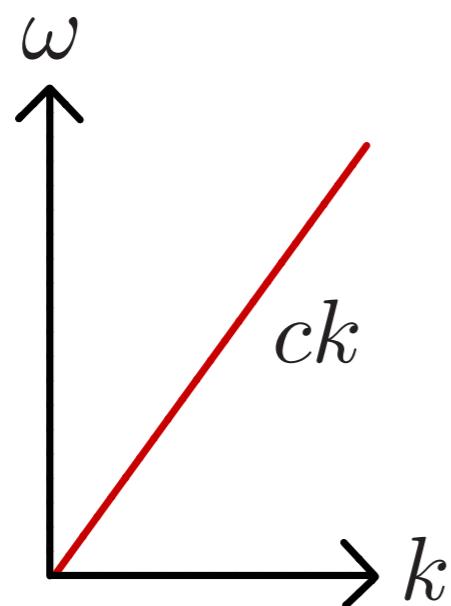
Hohenberg and Halperin: Theory of dynamic critical phenomena

TABLE I. Some dynamical models treated by renormalization-group methods.

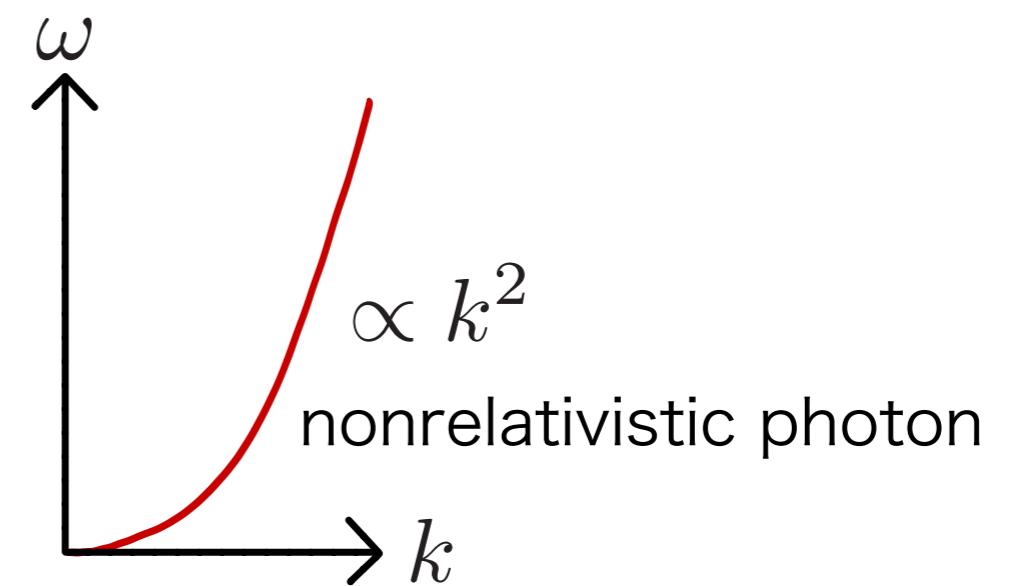
Model	Designation	System	Dimension order of parameter	Non-conserved fields	Conserved fields	Non-vanishing Poisson bracket
Relaxational	A	Kinetic Ising anisotropic magnets	n	ψ	None	None
	B	Kinetic Ising uniaxial ferromagnet	n	None	ψ	None
	C	Anisotropic magnets structural transition	n	ψ	m	None
Fluid	H	Gas-liquid binary fluid	1	None	ψ, \mathbf{j}	$\{\psi, \mathbf{j}\}$
Symmetric planar magnet	E	Easy-plane magnet, $h_z = 0$	2	ψ	m	$\{\psi, m\}$
Asymmetric planar magnet	F	Easy-plane magnet, $h_z \neq 0$ superfluid helium	2	ψ	m	$\{\psi, m\}$
Isotropic antiferromagnet	G	Heisenberg antiferromagnet	3	ψ	m	$\{\psi, m\}$
Isotropic ferromagnet	J	Heisenberg ferromagnet	3	None	ψ	$\{\psi, \psi\}$

Gapless modes are important.

Electromagnetic waves



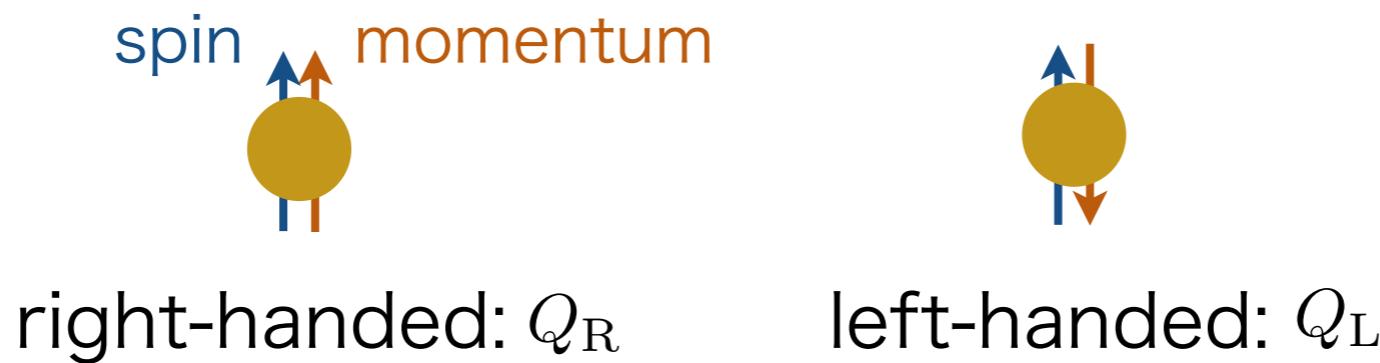
Maxwell equations in vacuum



Today's talk

Chirality and Quantum anomaly

- Chirality



$$Q_5 \equiv Q_R - Q_L \quad \text{chirality charge}$$

$$Q \equiv Q_R + Q_L \quad \text{total charge}$$

- Quantum anomaly

$$\partial_t n_5 + \nabla \cdot j_5 = CE \cdot B$$

Outline

- Setup
- Static critical phenomena
- Dynamic critical phenomena
- Conclusion

Setup

- Massless two-flavor QCD at finite T and B_{ex}
- Chiral symmetry:

$$\begin{array}{ccc} B_{\text{ex}} = 0 & B_{\text{ex}} \neq 0 & \text{ground state} \\ \text{SU}(2)_L \times \text{SU}(2)_R & \longrightarrow & \text{U}(1)^{\tau^3} \times \text{U}(1)_5^{\tau^3} \\ q = \begin{pmatrix} u \\ d \end{pmatrix} & q \rightarrow e^{i\alpha\tau^3} e^{i\alpha_5\gamma_5\tau^3} q & \cancel{\text{U}(1)^{\tau^3} \times \text{U}(1)_5^{\tau^3}} \\ \text{"isospin"} & & \langle \bar{q}q \rangle \neq 0 \end{array}$$

- Second-order chiral phase transition

Hydrodynamic variables

- (Complex) order parameter: Φ
- Chirality charge density: n_5
- Electric field: E
- Magnetic field: $B = B_{\text{ex}} + \delta B$

Ginzburg-Landau Theory

$$F = \int_V \left(\frac{r}{2} |\Phi|^2 + \frac{u}{4} |\Phi|^4 + \frac{a}{2} |\nabla \Phi|^2 + \frac{1}{2} n_5^2 + \frac{1}{2} \mathbf{E}^2 + \frac{1}{2} \mathbf{B}^2 \right)$$

- Near second-order phase transition \rightarrow small order parameters
- Long-range behavior \rightarrow derivative expansion
- QCD symmetries \rightarrow constraints on the expansion
 - chiral symmetry and CPT symmetries

Ginzburg-Landau Theory

$$F = \int_V \left(\frac{r}{2} |\Phi|^2 + \frac{u}{4} |\Phi|^4 + \frac{a}{2} |\nabla \Phi|^2 + \frac{1}{2} n_5^2 + \frac{1}{2} \mathbf{E}^2 + \frac{1}{2} \mathbf{B}^2 \right)$$

$$h_0 \equiv \langle |\Phi| \rangle \sim \tau^\beta, \quad a \sim \xi^\eta, \quad \xi \sim \tau^{-\nu} \quad \left(\tau \equiv \frac{T - T_c}{T_c} \right)$$

Wilson-Fisher fixed point (2-components)

Dynamics

- Preliminary

$$\Phi = (h + h_0)e^{i\pi^0/h_0}$$

↑
amplitude mode
Nambu-Goldstone mode
(linearly decoupled)

Langevin theory 1/2

- π^0 and n_5

$$\partial_t \pi^0 = -\kappa \frac{\delta F}{\delta \pi^0} + \int_V [\pi^0, n_5] \frac{\delta F}{\delta n_5} + \xi_{\pi^0}$$

$$\partial_t n_5 = \lambda \nabla^2 \frac{\delta F}{\delta n_5} + \int_V [n_5, \pi^0] \frac{\delta F}{\delta \pi^0} + \int_V [n_5, E^i] \frac{\delta F}{\delta E^i} + \xi_{n_5}$$

$$[\pi^0(t, \mathbf{r}), n_5(t, \mathbf{r}')] = h_0 \delta(\mathbf{r} - \mathbf{r}') \delta(t - t')$$

$$[n_5(t, \mathbf{r}), E^i(t, \mathbf{r}')] = C B_i \delta(\mathbf{r} - \mathbf{r}') \delta(t - t')$$

Anomalous commutation relation

S. Adler and D. Boulware (1969)

Langevin theory 2/2

- Maxwell equations:

$$\partial_t \mathbf{E} = \nabla \times \mathbf{B} - \mathbf{j}, \quad \partial_t \mathbf{B} = -\nabla \times \mathbf{E}$$

- Current term:

$$\mathbf{j} = \sigma \mathbf{E} + C n_5 \mathbf{B} + \frac{C}{h_0} \nabla \pi^0 \times \mathbf{E} + \boldsymbol{\xi}$$

↑
Chiral magnetic effect

$$\langle \xi_i(t, \mathbf{r}) \xi_j(t, \mathbf{r}') \rangle = 2\sigma_{ij} \delta(\mathbf{r} - \mathbf{r}') \delta(t - t')$$

Linear analysis

- Linearize: $B = B_{\text{ex}} + \delta B$ etc.
- Hydrodynamic (gapless) modes:
 - 1 diffusive mode
 - nonrelativistic photon:

$$\omega = \pm \sqrt{\frac{v^2}{C^2 B_{\text{ex}}^2} - D^2 k_{\perp}^2 - i D k_{\perp}^2} \quad (k_{\perp} \perp B_{\text{ex}})$$

$$v^2 \equiv a h_0^2, \quad D \equiv \frac{\kappa}{2} + \frac{\sigma v^2}{2 C^2 B_{\text{ex}}^2}$$

Dynamic scaling hypotheses

- $\omega \sim \xi^{-z}$ for all terms

$$\omega = \pm \sqrt{\frac{v^2}{C^2 B_{\text{ex}}^2} - D^2 \mathbf{k}_\perp^2} - iD \mathbf{k}_\perp^2$$

$$D \sim \frac{v}{CB_{\text{ex}}} \sim \xi^{-z+2}$$

- From statics: $v^2 \equiv ah_0^2 \sim \xi^{d-2}$
- Dynamic critical exponent:

$$z = \frac{d}{2} + 1$$

New dynamic universality class
beyond Hohenberg and Halperin's classification

Conclusion

- Dynamic critical phenomena of the second-order chiral phase transition under external magnetic field with dynamical electromagnetic fields
- Hydrodynamic modes: Nonrelativistic photon
- New dynamic universality class beyond conventional classification