

The Ministry of Education, Culture, Sports, Science and Technology (MEXT)
Grant-in-Aid for Scientific Research on Innovative Areas

Research Fee
科研費
KARENHI

重力波天体の多様な観測による宇宙物理学の新展開

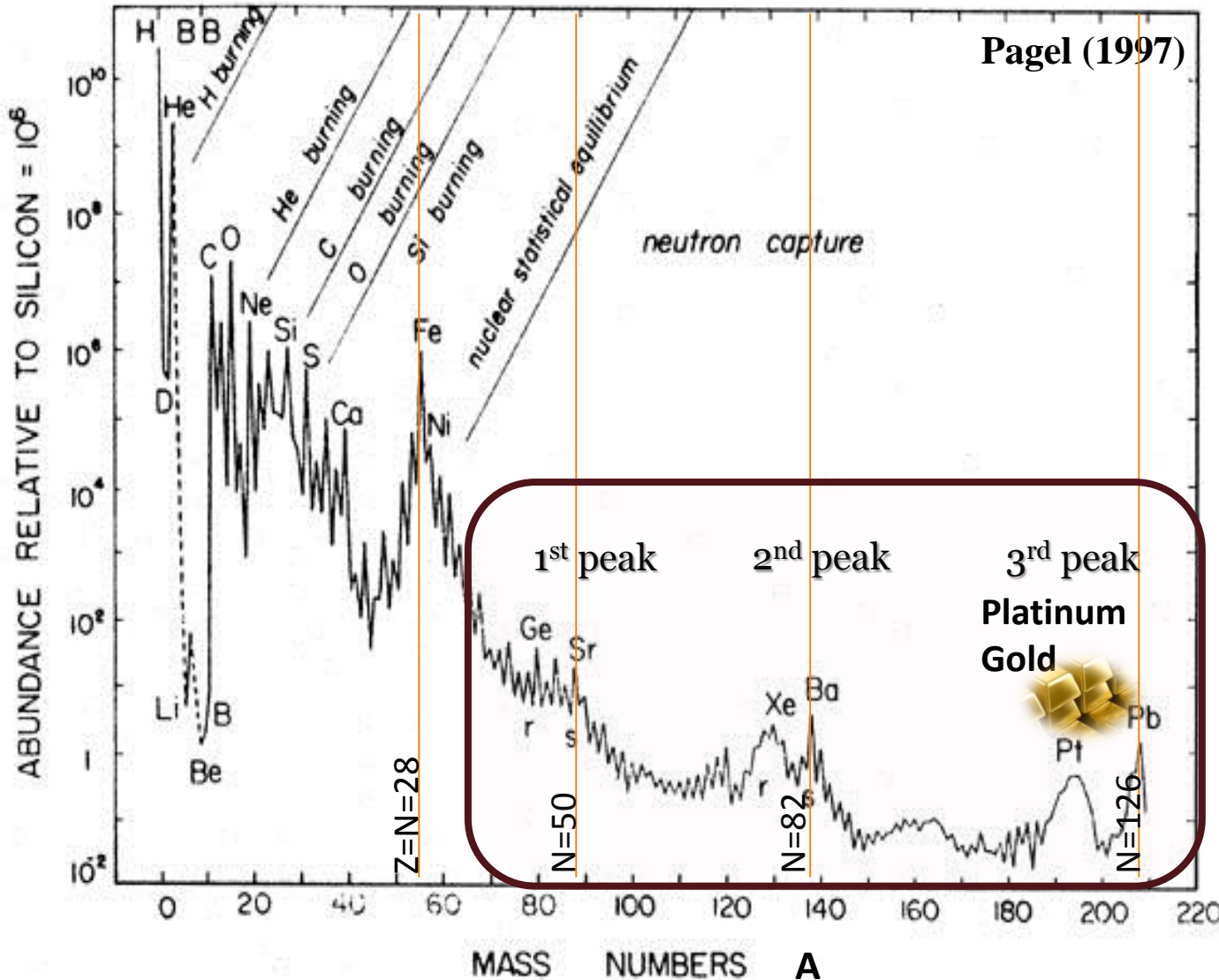
New development in astrophysics through multimessenger observations of gravitational wave sources



Binary neutron star merger and **r-process**

Yuichiro Sekiguchi (Toho Univ.)

Solar abundance of nuclei

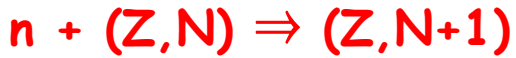


Pagel (1997)

- ▶ Basic feature : exponential decay with mass number + constant tail
- ▶ Characteristic features:
 - ▶ Peak in iron-group
 - ▶ Deficient of D, Li, Be, and B
 - ▶ Enhancement of α -nuclei (C, O, Ne, Si,..)
 - ▶ **Peaks in heavier region associated with n-magic numbers,**
 - ▶ **made by neutron capture processes**

Neutron capture processes:

free from Coulomb barrier



n-capture

versus



β -decay

$$\tau_n < \tau_\beta$$

rapid neutron-capture process
(r-process)

large neutron densities

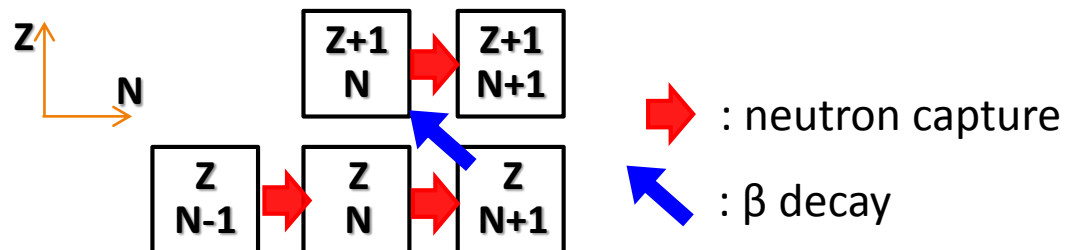
- Can synthesize all heavy nuclei

$$\tau_n > \tau_\beta$$

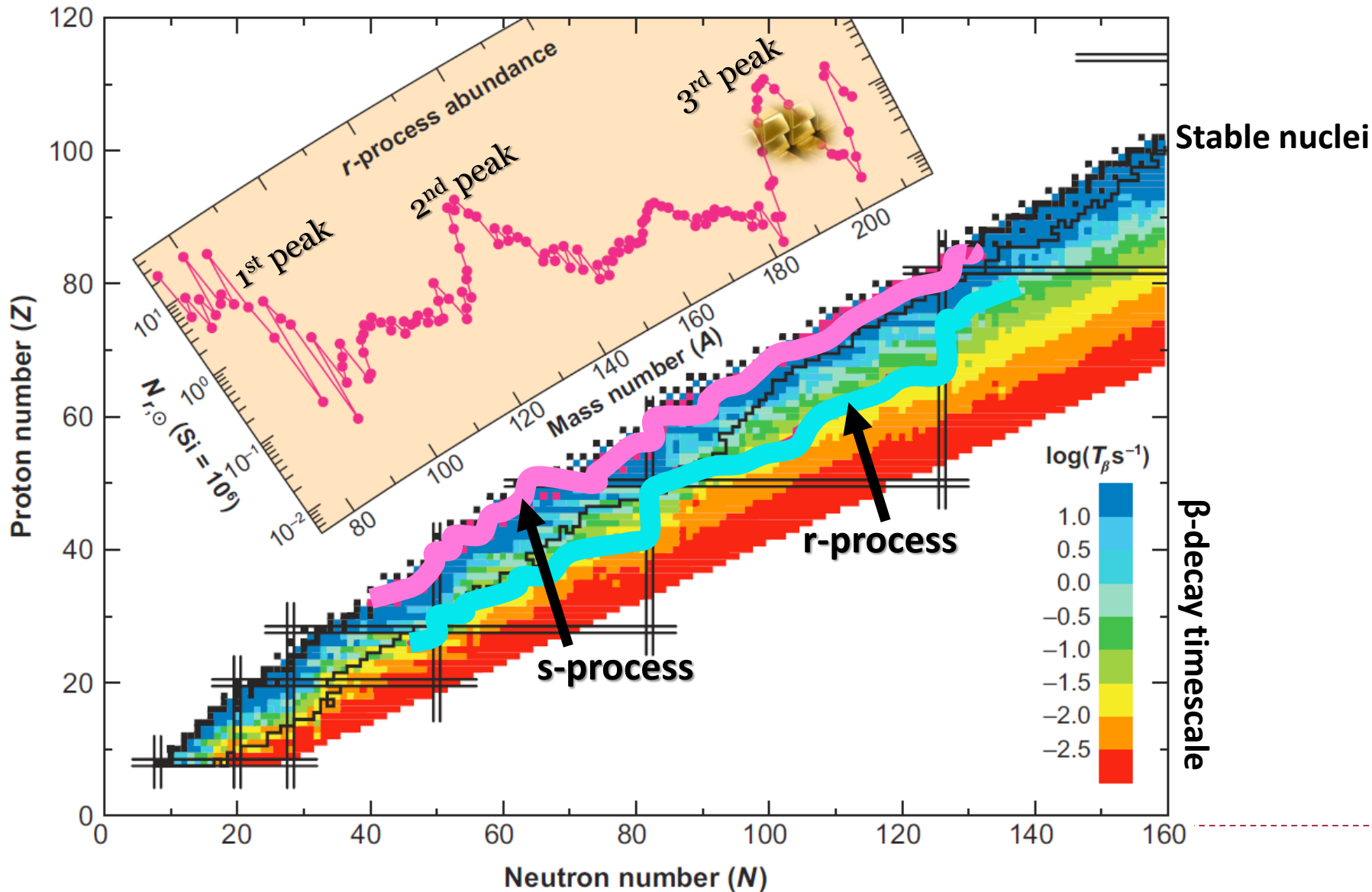
slow neutron-capture process
(s-process)

moderate neutron densities

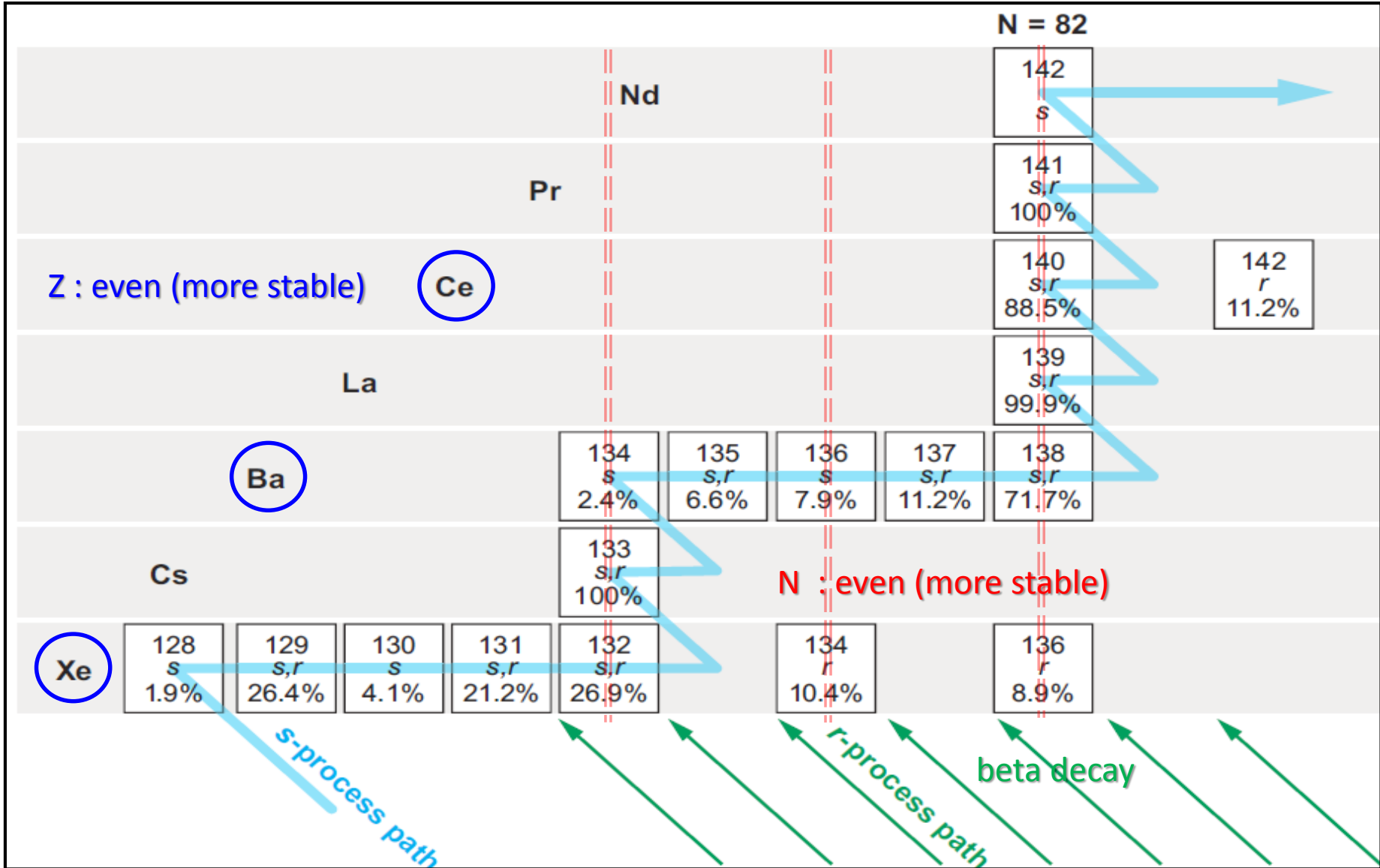
- does not synthesize all heavy nuclei
- terminates at Pb, Bi



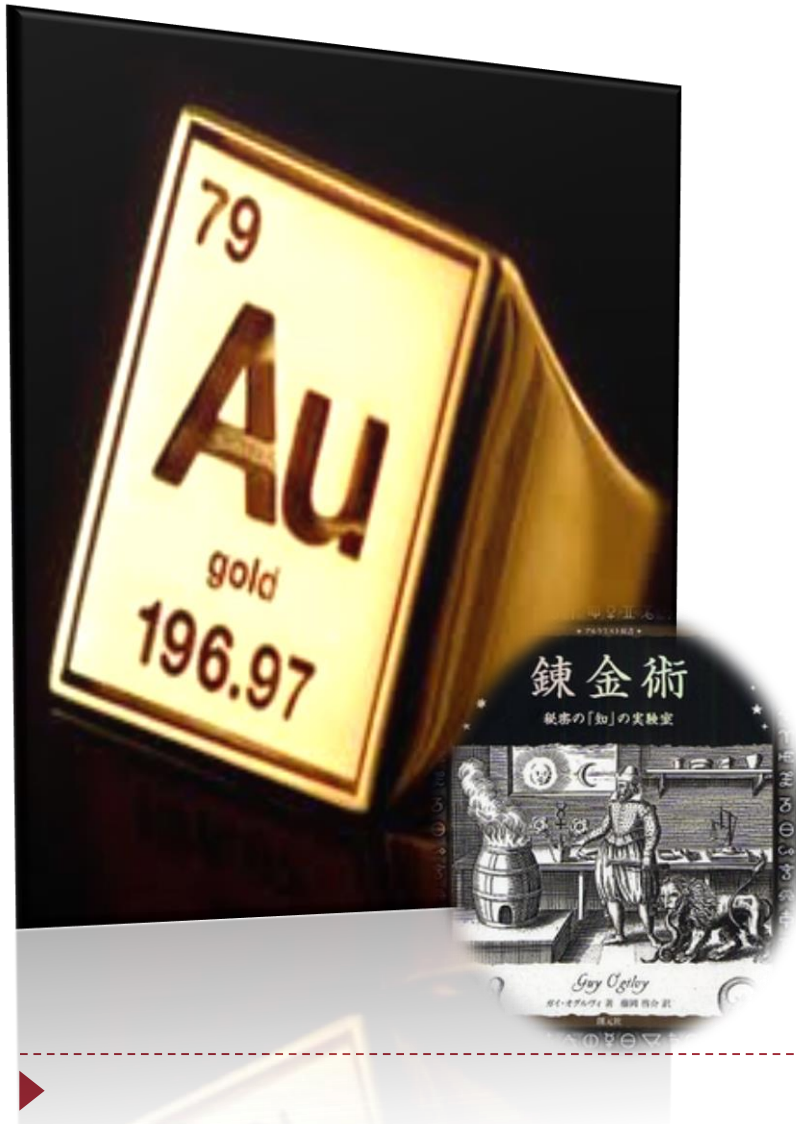
s-process / r-process path



s-process / r-process path



To be an alchemist : recipe to cook gold



- ▶ Neutron capture : packing neutrons into 'seed' nuclei $n + (Z,N) \Rightarrow (Z,N+1)$
 - ▶ Large #neutron/#seed ratio is required
 - ▶ $A(\text{gold}) - A(\text{seed}) \sim 100$
- ▶ **(1) Low electron fraction Y_e**
 - ▶ Y_e = number of electrons per baryon \sim # of proton $\sim 1 -$ # of neutron
 - ▶ To have a large number of free neutrons
- ▶ **(2) Higher entropy per baryon**
 - ▶ To slow the seed nuclei production
- ▶ **(3) Short expansion timescale**
 - ▶ To freeze seed production with rapid decrease of temperature

What is the melting pot for r-process ?

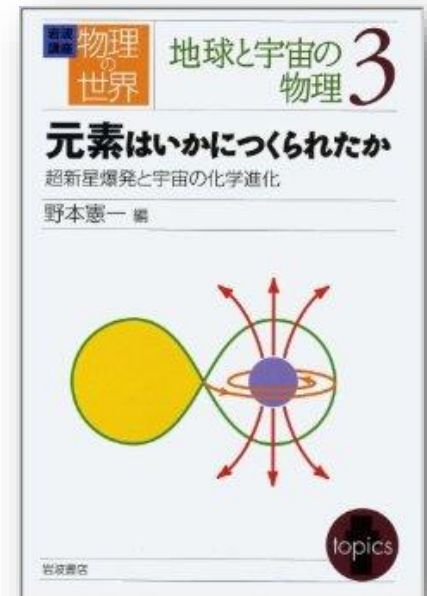


▶ Supernova (SN) explosion (+ PNS v-driven wind) : (*Burbidge et al. 1957*)

- ▶ Textbooks tell you that SNe are the origin of heavy elements, but
- ▶ theoretically disfavored (Roberts et al. 2010, 2012)

▶ NS-NS/BH binary merger: (*Lattimer & Schramm 1974*)

- ▶ Observationally disfavored ?? (Argust et al. 2004)
- ▶ Too neutron rich ???



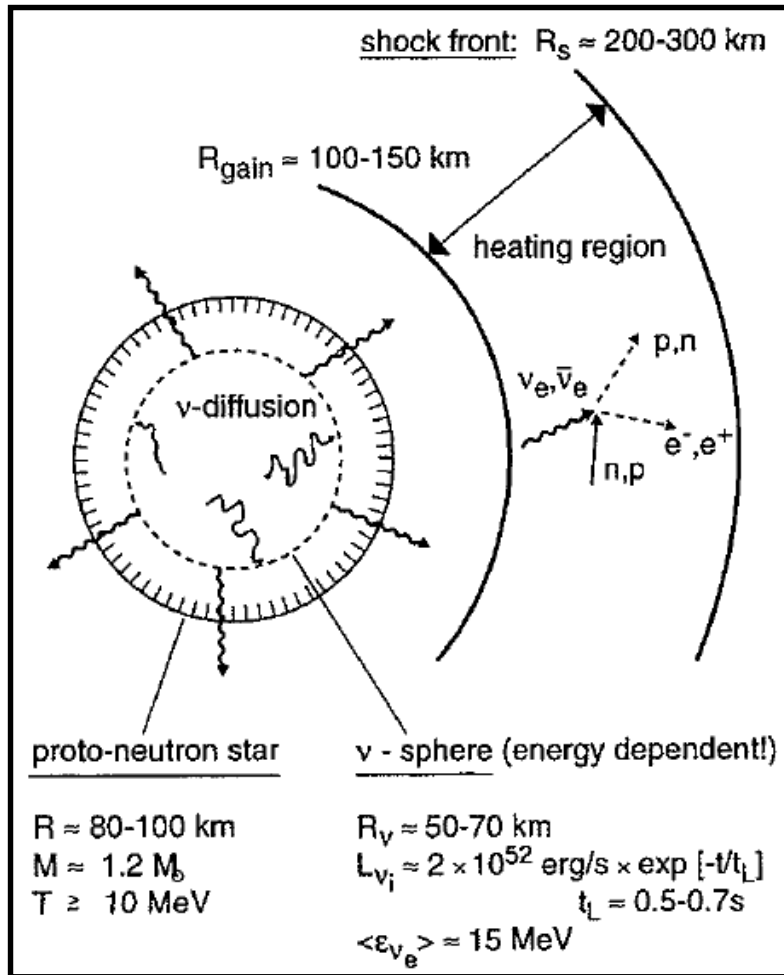
What is the melting pot for r-process ?



- ▶ **Supernova (SN) explosion:** (*Burbidge et al. 1957*)
 - ▶ Smaller entropy/per baryon than previously expected (e.g., *Janka et al. 1997*)
 - ▶ Previous expectation ($s/kB > 200$) => recent update $s/kB \sim 100-150$
 - ▶ Neutrino heating mechanism of SNe explosion:
 - ▶ Neutrinos from PNS try to make the flow proton-rich via $n+\nu \rightarrow p+e$ and $p+\bar{\nu} \rightarrow n+e^+$



Overall picture of the neutrino heating mechanism



Cartoon by E. Muller (1998)

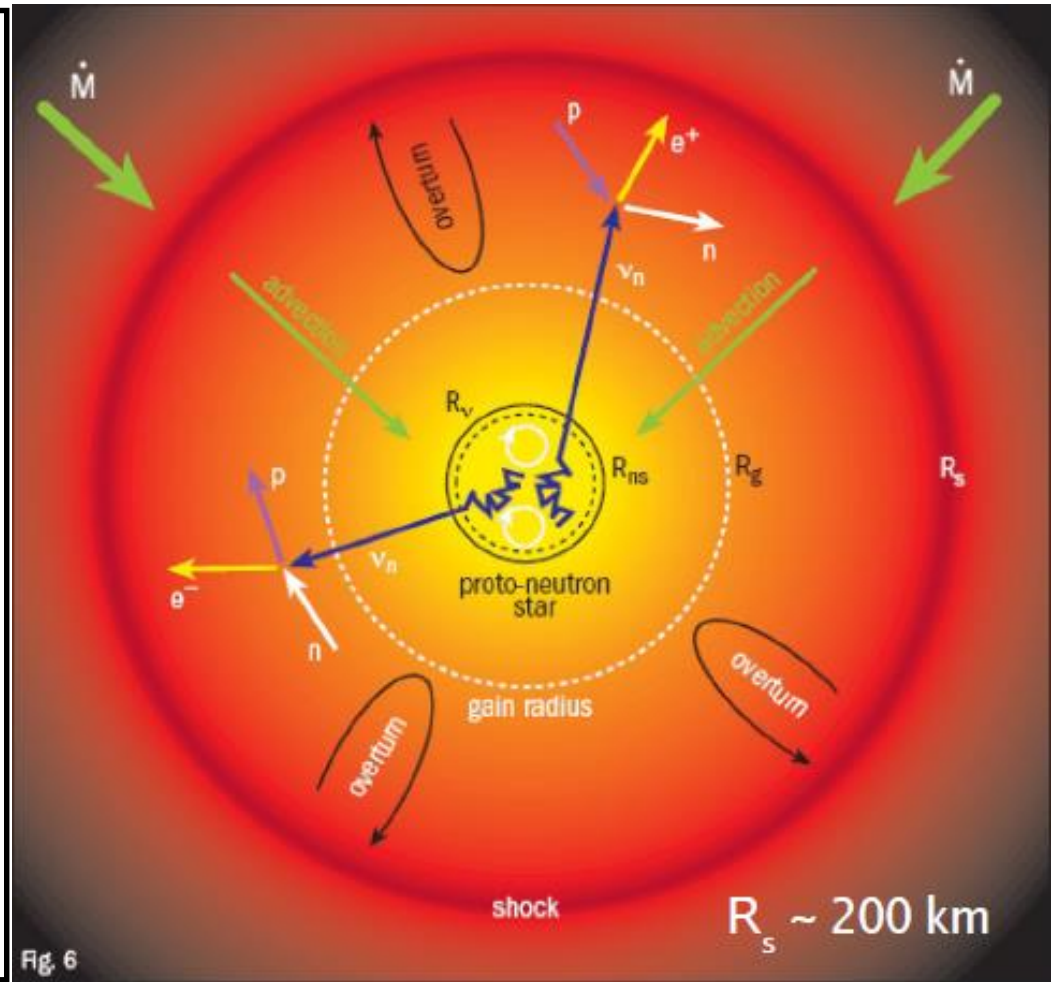


Fig. 6

Cartoon by T. Janka



What is the melting pot for r-process ?



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- ▶ Smaller entropy/per baryon than previously expected (e.g., *Janka et al. 1997*)
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- ▶ Neutrino heating mechanism of SNe explosion:
- ▶ Neutrinos from PNS may make the flow proton-rich via
 $n + \nu \rightarrow p + e$ and $p + \bar{\nu} \rightarrow n + e^+$
 - ▶ Note : neutrons are heavier than proton => tendency of being proton rich.
 - ▶ Whether the flow becomes proton rich or not depends on mean neutrino energy
 - ▶ Mass difference vs. neutrino energy difference (and luminosities)

$$\Delta\varepsilon = \bar{\varepsilon}_{\nu} - \varepsilon_{\nu} \quad \text{vs.} \quad \Delta m = m_n - m_p$$

$$\Delta\varepsilon > \sim 4\Delta m \quad (\text{neutron rich}) \quad \Delta\varepsilon < \sim 4\Delta m \quad (\text{proton rich})$$

- ▶ Higher electron anti-neutrino energy => effectively larger proton mass

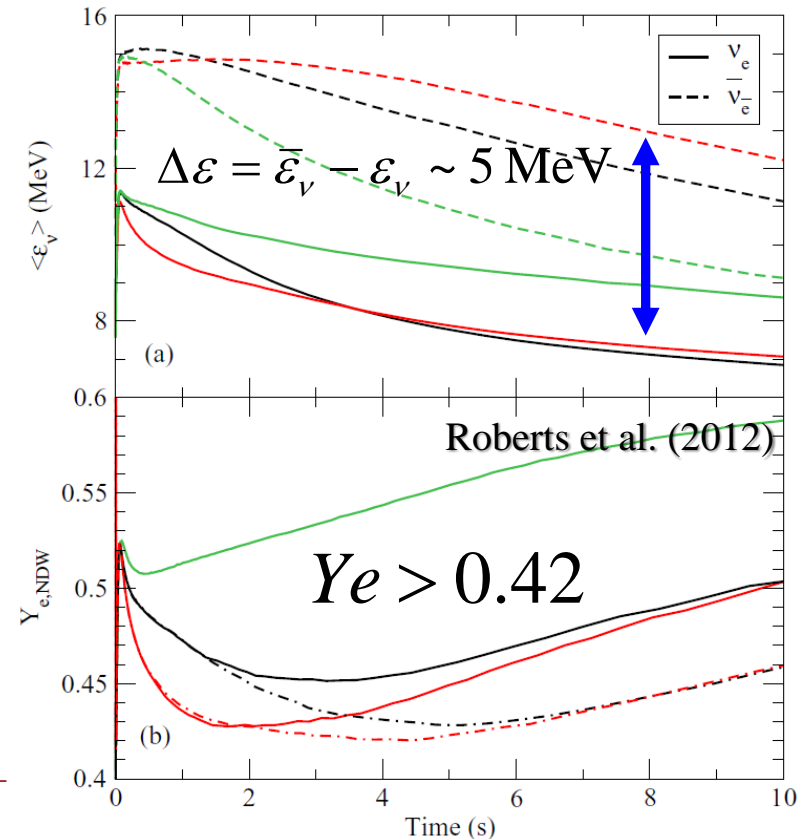
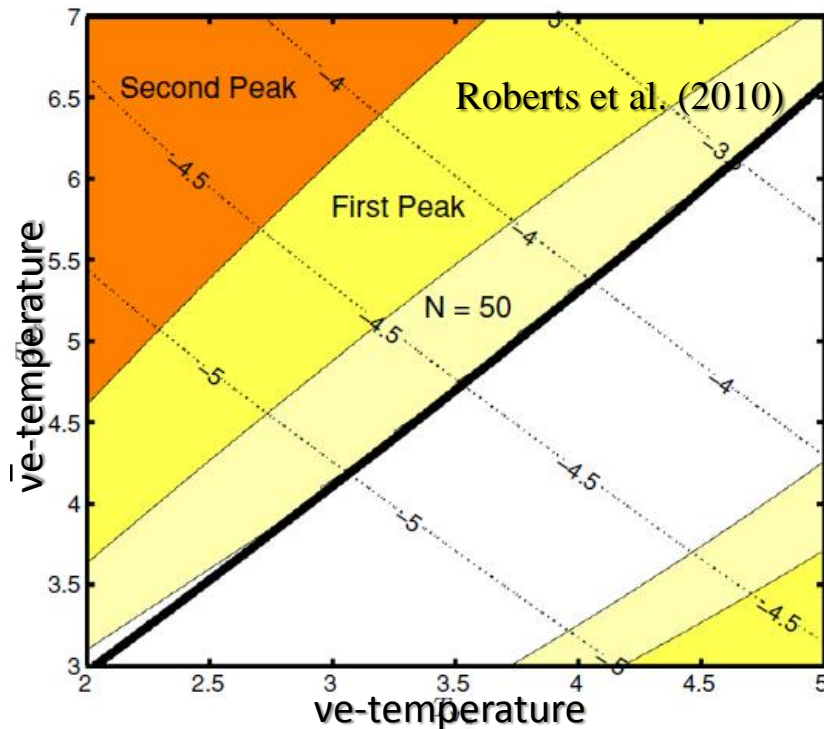


What is the melting pot for r-process ?



▶ **Supernova (SN) explosion:** (*Burbidge et al. 1957*)

- ▶ Smaller entropy/per baryon than previously expected (e.g., *Janka et al. 1997*)
- ▶ Neutrinos from PNS make the flow proton-rich via weak interactions
- ▶ \Rightarrow only weak r-process (up to 2nd peak, no gold (3rd peak)!) (*Roberts et al. 2010, 2012;*
Wajajo et al, 2013 etc.)



What is the melting pot for r-process ?



▶ **Supernova (SN) explosion:** (*Burbidge et al. 1957*)

- ▶ Smaller entropy/per baryon than previously expected (e.g., *Janka et al. 1997*)
 - ▶ Previous expectation ($s/kB > 200$) => recent update $s/kB \sim 100-150$
- ▶ Neutrinos from PNS try to make the flow proton-rich via $n + \nu \rightarrow p + e$ and $p + \bar{\nu} \rightarrow n + e^+$
 - ▶ Note : neutrons are heavier than proton
 - ▶ Whether the flow becomes proton rich or not depends on neutrino energy
- ▶ According to the recent studies, only weak r-process occurs (up to 2nd peak, no gold (3rd peak)!) (*Roberts et al. 2010, 2012*)
 - ▶ Electron capture SN : *Hoffman et al. 2008; Wanajo et al. 2009*
 - ▶ (Iron) core collapse SN : *Fisher et al. 2010; Hudepohl et al. 2010; Wanajo et al. 2011; Roberts et al. 2012*



What is the melting pot for r-process ?

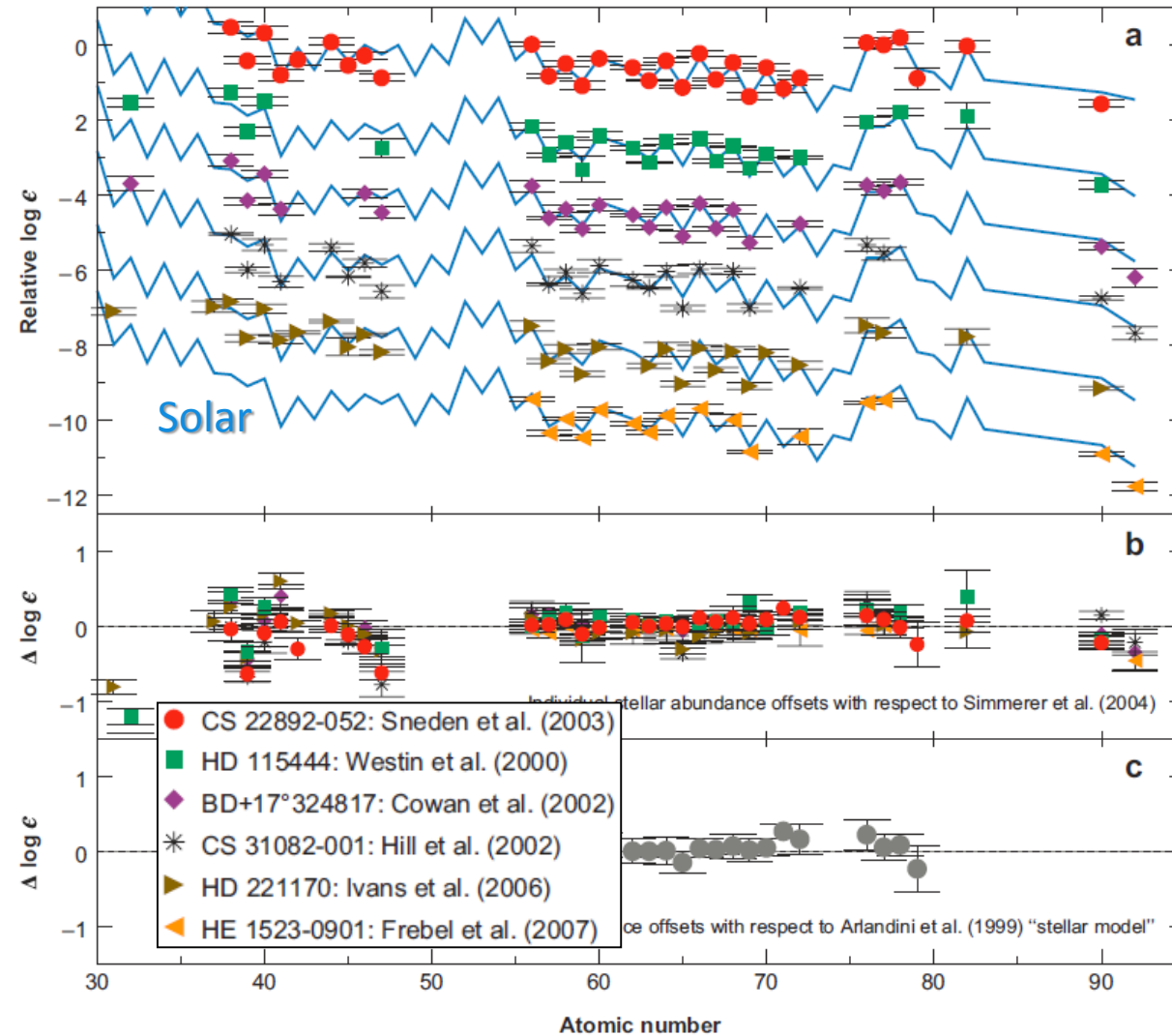


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- ▶ Supernova can be the origin of r-process nuclei only if
 - ▶ The explosion mechanism is not due to the popular neutrino heating (e.g., magneto-rotational; *Winteler et al. 2012*)
 - ▶ or
 - ▶ Our knowledge of neutrino and nuclear physics is insufficient

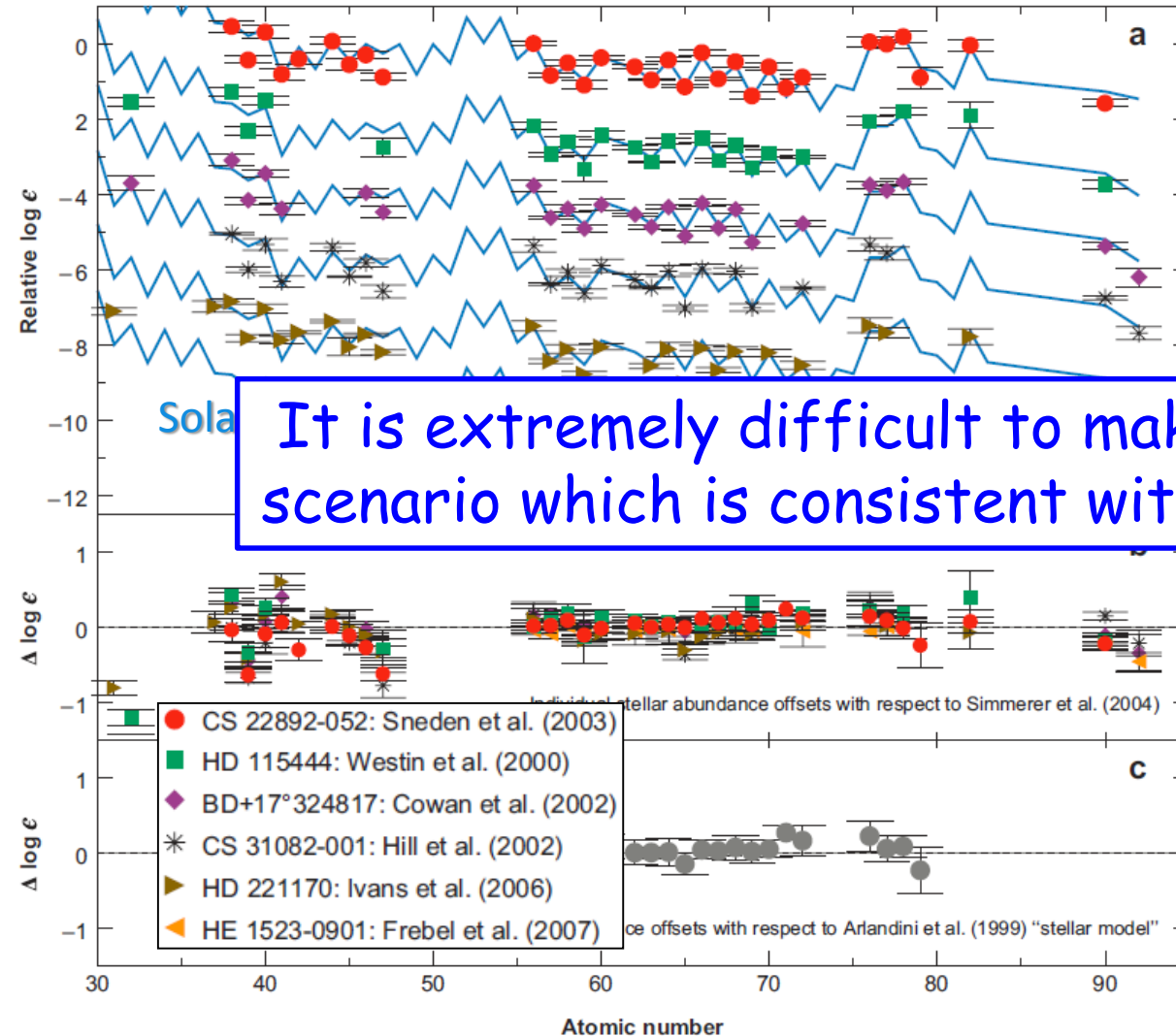


A key observation to resolve the problem: Universality of the r-process cite



- ▶ Abundance pattern comparison :
 - ▶ r-rich low metallicity stars
 - ▶ Solar neighborhood
- ▶ Low metallicity means
 - ▶ Such stars experience only one/two r-process events
 - ▶ Such stars preserve the original pattern of the r-process events (chemical fossil)

A key observation to resolve the problem: Universality of the r-process cite



- ▶ Abundance pattern comparison :
- ▶ The solar and chemical fossil pattern agree well
- ▶ for $Z > 35-40$ ($A > 85-90$)

It is extremely difficult to make a model in SNe scenario which is consistent with the universality.

the original pattern

- ▶ => these observations strongly suggest that the (main) r-process event synthesize the elements with a pattern similar to solar (Universality)

What is the melting pot for r-process ?

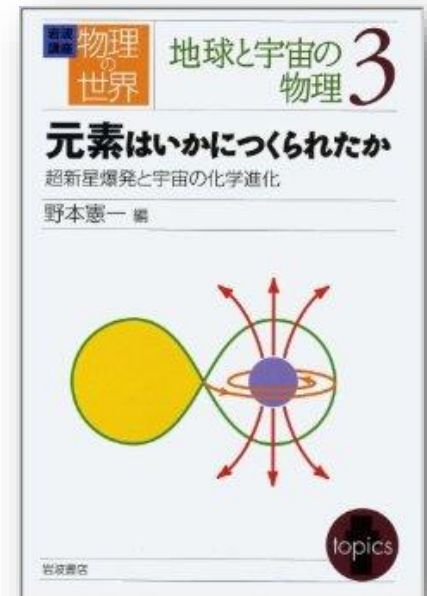


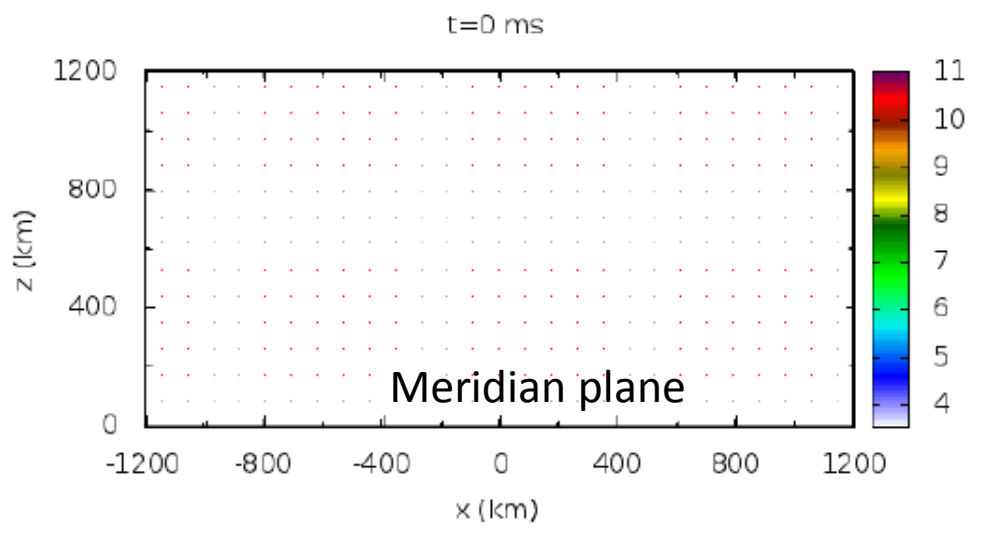
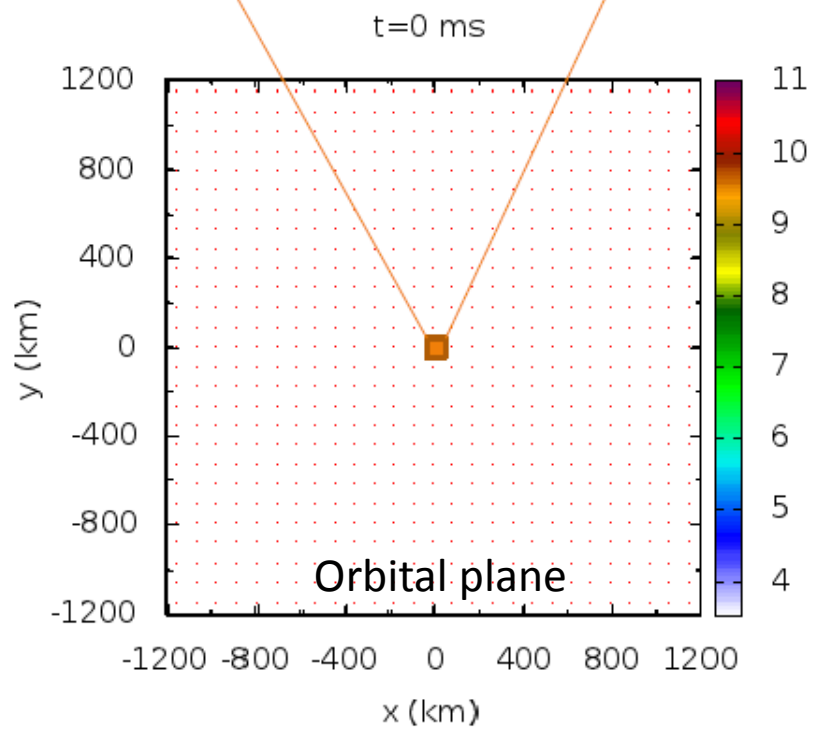
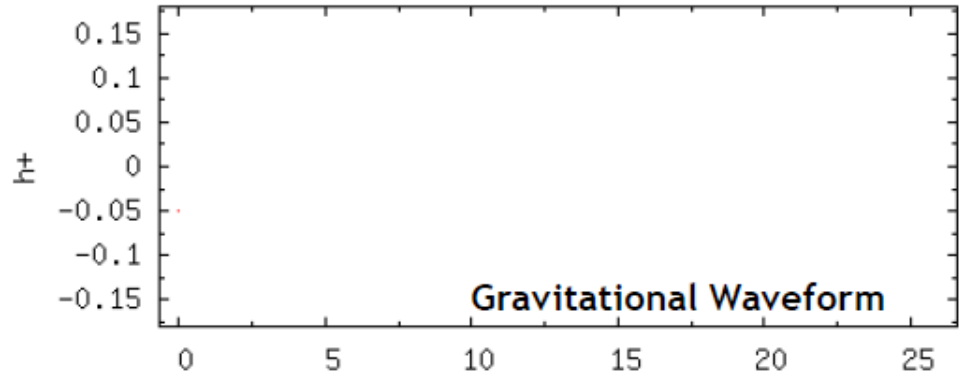
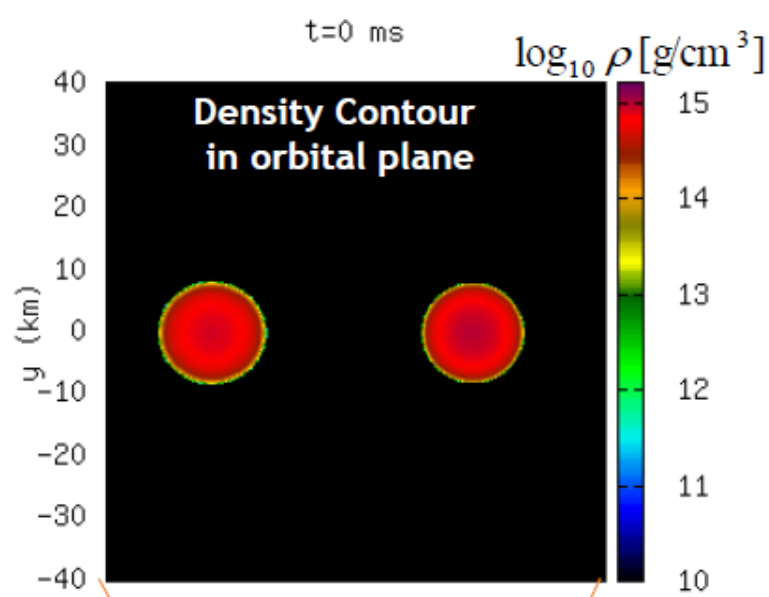
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- ▶ Observationally disfavored ?? (Argust et al. 2004)
- ▶ Too neutron rich ???





Animation by Hotokezaka

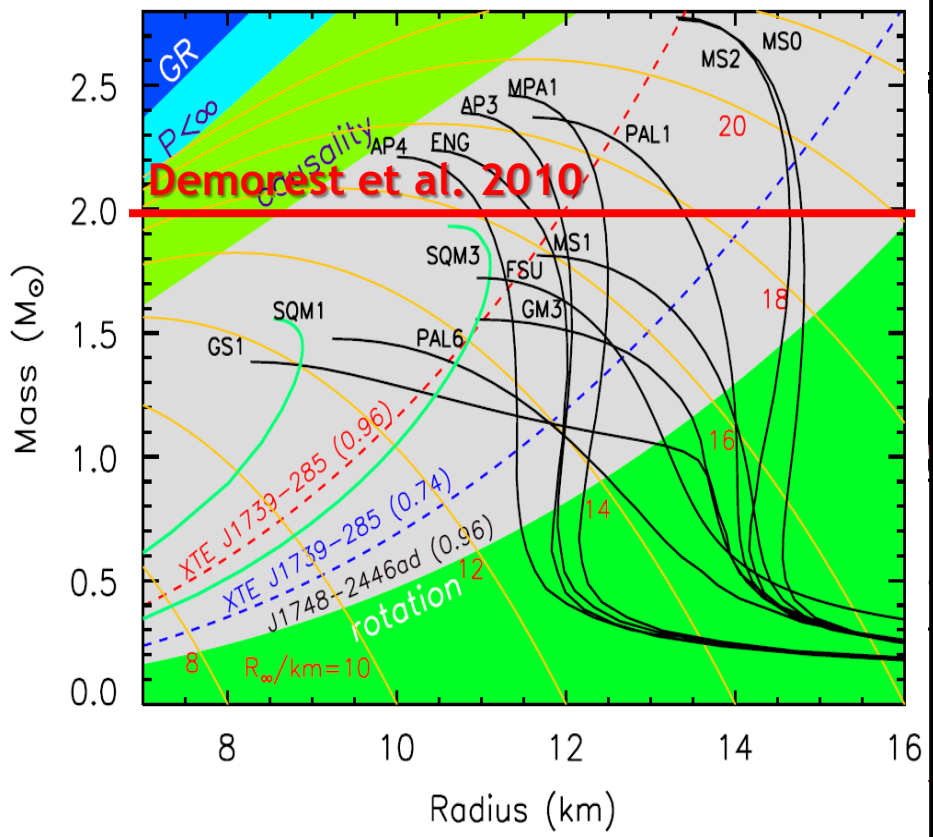
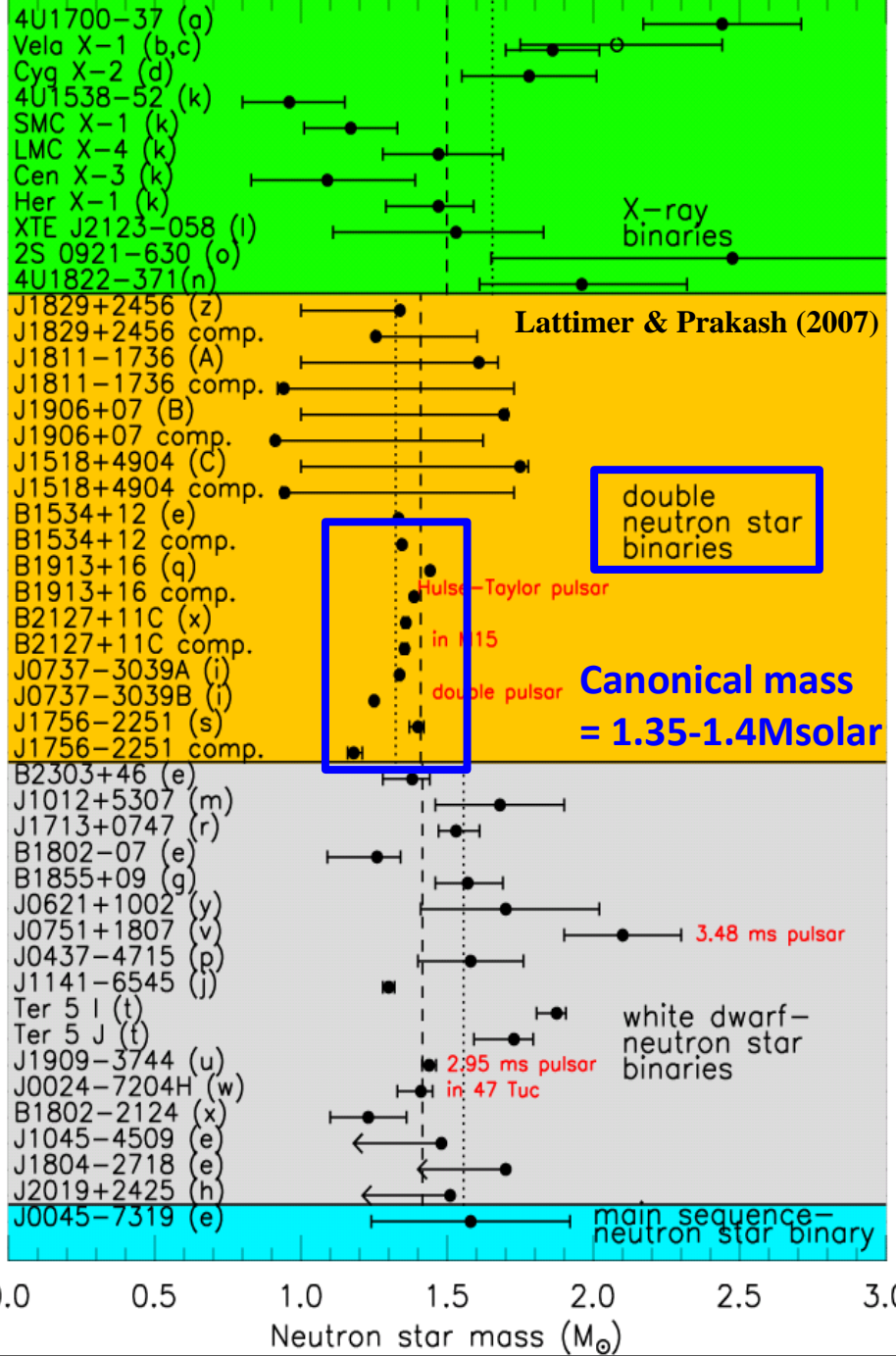
Sekiguchi et al. PRL (2011a, 2011b)
 Kiuchi et al. PRL (2010); Hotokezaka et al. (2013)

Bill Saxton,
NRAO/AUI/NSF

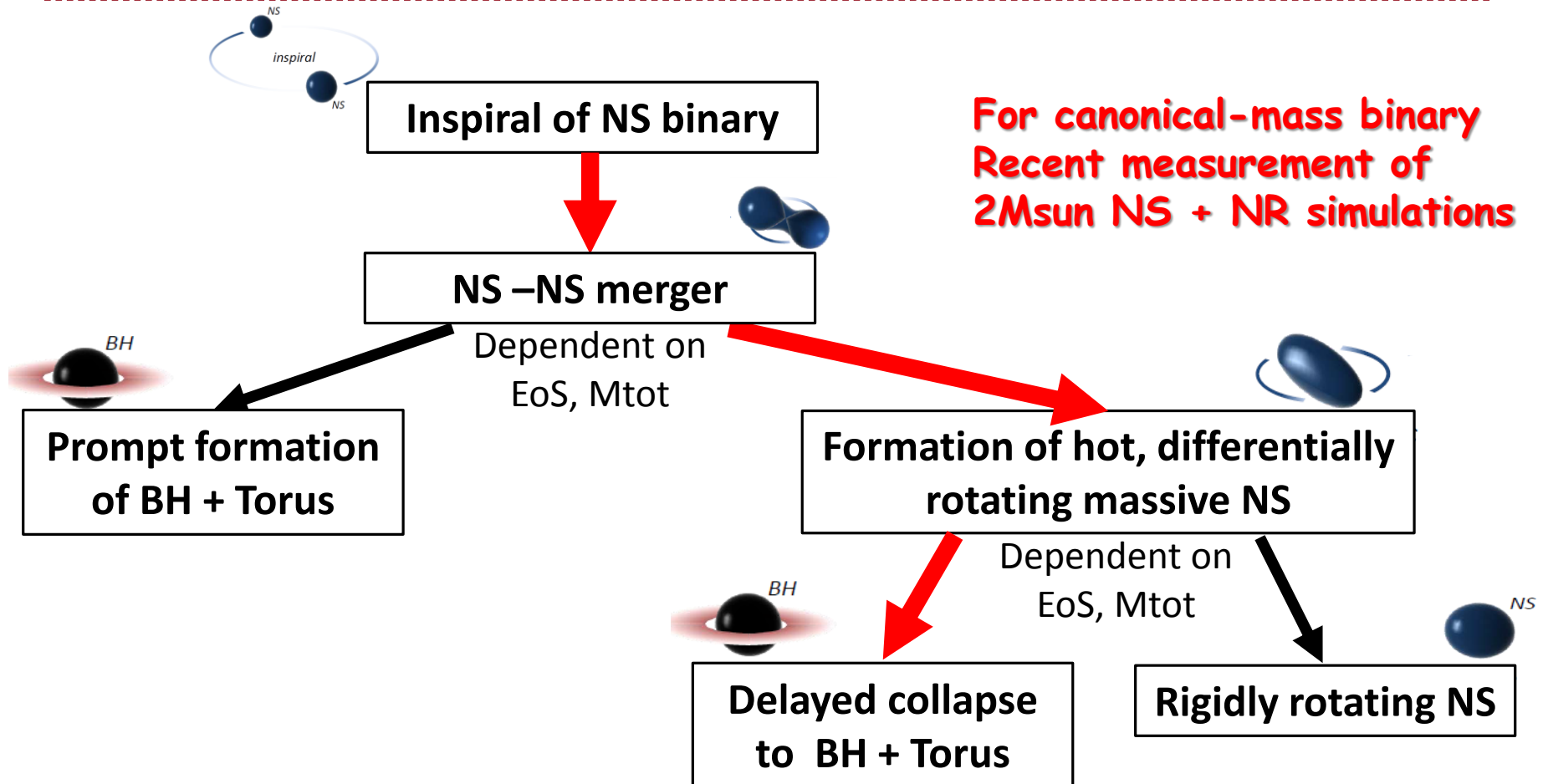
NS-NS 1

of NS binary

updated 8 December 2006



Evolution of NS-NS mergers



Causality limit : NS radius estimation

- ▶ The measurement of flux and temperature yields an apparent angular size (pseudo-BB)

$$\frac{R_\infty}{D} = \frac{R}{D} \frac{1}{\sqrt{1 - GM/Rc^2}} \quad F \propto T_{\text{eff}}^4 \frac{R_\infty^2}{D^2}$$

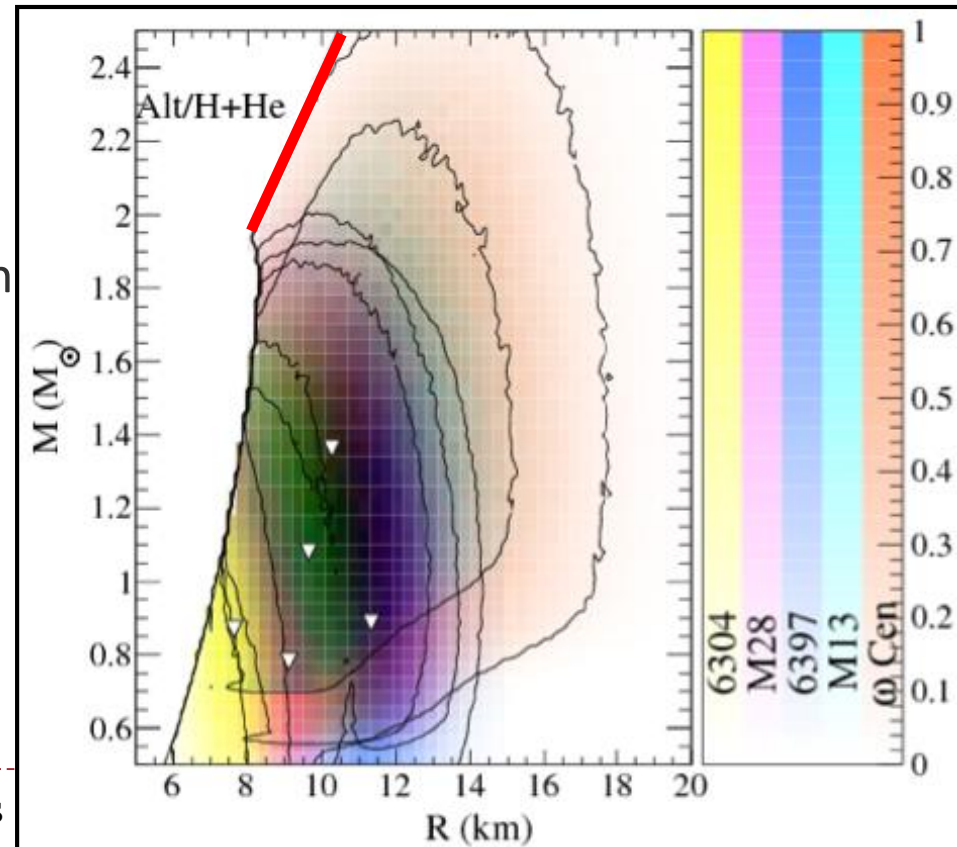
- ▶ Many uncertainties : **redshift**, distance, interstellar absorption, atmospheric composition

▶ Good Targets:

- ▶ Quiescent X-ray binaries in globular clusters (D, composition)
- ▶ Bursting sources with peak flux close to Eddington limit (M)

▶ Imply rather small radius

- ▶ **If true, maximum mass may not be much greater than 2Msun**

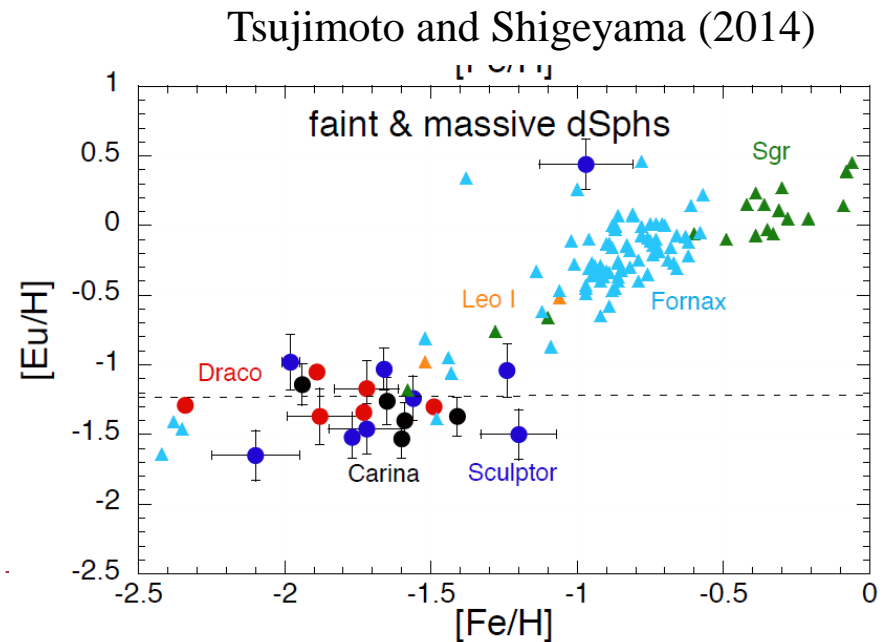
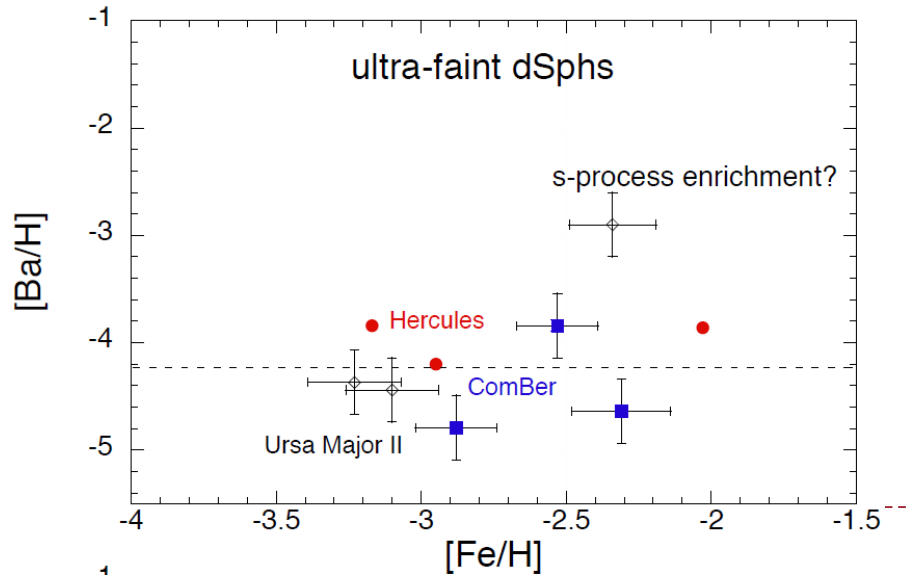


- ▶ Lattimer & Steiner 2014 for quiescent LMXBs

What is the melting pot for r-process ?



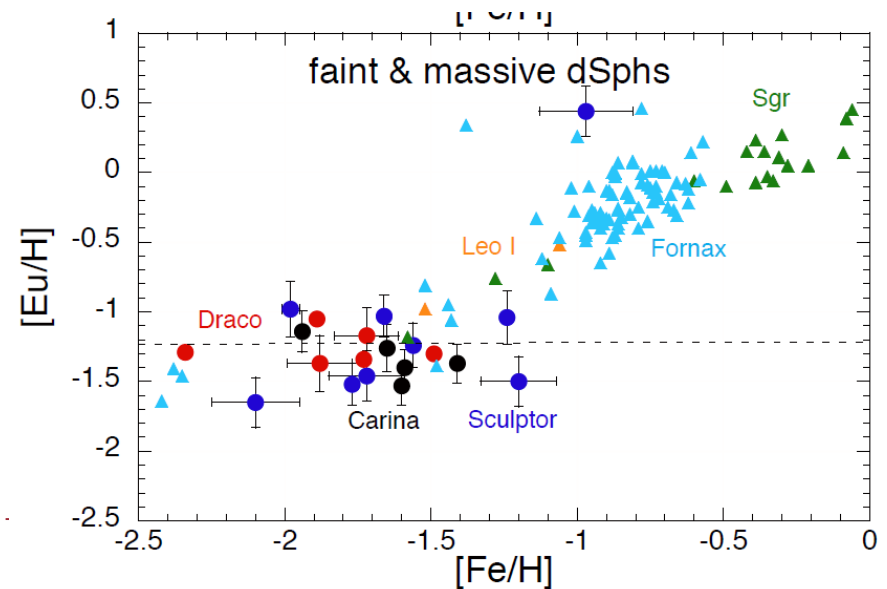
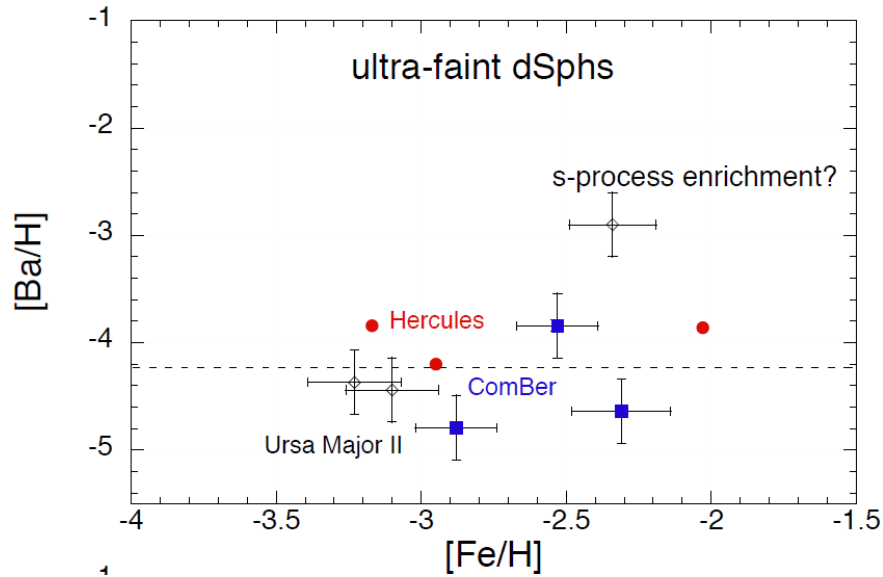
- ▶ **Observationally NOT disfavored ??** (*Tsujimoto and Shigeyama. 2014*)
 - ▶ No enrichment of Eu in ultra dwarf galaxies but Fe increases
 - ▶ No r-process events (No Eu) but a number of SNe (Fe \uparrow)
 - ▶ If SNe are the r-process cite, both Eu and Fe should increase
 - ▶ Suggest different origin for Fe and Eu



What is the melting pot for r-process ?



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 - ▶ No r-process events (No Eu) but a number of SNe (Fe \uparrow)
 - ▶ If SNe are the r-process cite, both Eu and Fe should increase
 - ▶ Suggest different origin for Fe and Eu
 - ▶ **Enrichment of Eu in massive dwarfs**
 - ▶ **event rate is estimate as 1/1000 of SNe : consistent with BNS merger**



Further observational evidence ?

Kilo-nova / Macro-nova / r-process-nova

- ▶ **EM transients possibly powered by radioactivity of the r-process elements were expected (Li & Paczynski 1998) and found ([important GW counterpart](#))**

LETTER

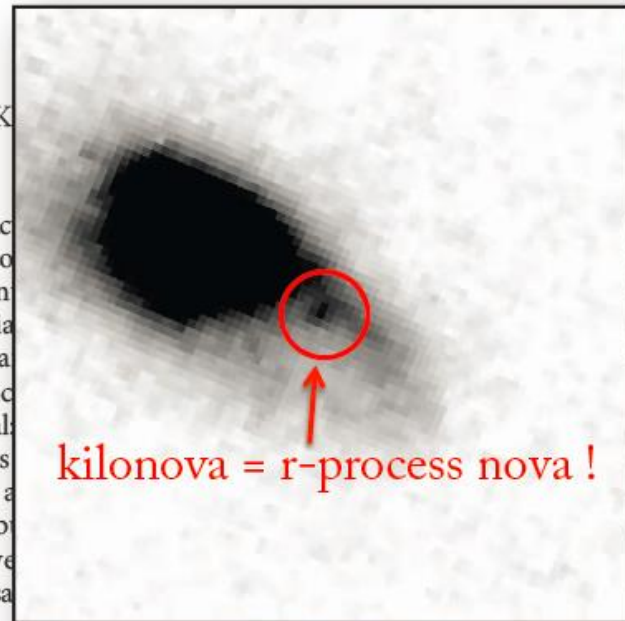
doi:10.1038/nature12505

A 'kilonova' associated with the short-duration γ -ray burst GRB 130603B

N. R. Tanvir¹, A. J. Levan², A. S. Fruchter³, J. Hjorth⁴, R. A. Hounsell³, K.

Short-duration γ -ray bursts are intense flashes of cosmic γ -rays, lasting less than about two seconds, whose origin is unclear^{1,2}. The favoured hypothesis is that they are produced by a relativistic jet created by the merger of two compact stellar objects (specifically two neutron stars or a neutron star and a black hole). This is supported by indirect evidence such as the properties of their host galaxies³, but unambiguous confirmation of the model is still lacking. Mergers of this kind are also expected to create significant quantities of neutron-rich radioactive species^{4,5}, whose decay should result in a faint transient, known as a 'kilonova', in the days following the burst⁶⁻⁸. Indeed, it is speculated that this mechanism may be the predominant source of stable r-process elements in the Universe^{5,9}.

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Short summary

- ▶ The origin of r-process nuclei : SNe vs. BNS merger
 - ▶ Key words
 - ▶ low Y_e required, universality of the pattern
 - ▶ Nice lecture by Evan for nucleosynthesis
 - ▶ SNe
 - ▶ Difficult to preserve n-rich condition necessary for the r-process
 - ▶ Extremely difficult to satisfy the universality
 - ▶ BNS
 - ▶ Recent theoretical and observational studies indicate BNS mergers are a promising candidate
 - ▶ **Kilonova-like signal : important as EM counterpart to GW**
 - ▶ How about from the universality point of view
 - Ye profile of merger ejecta



From the ‘Universality’ point of view : NS-NS merger ejecta: too neutron-rich ?

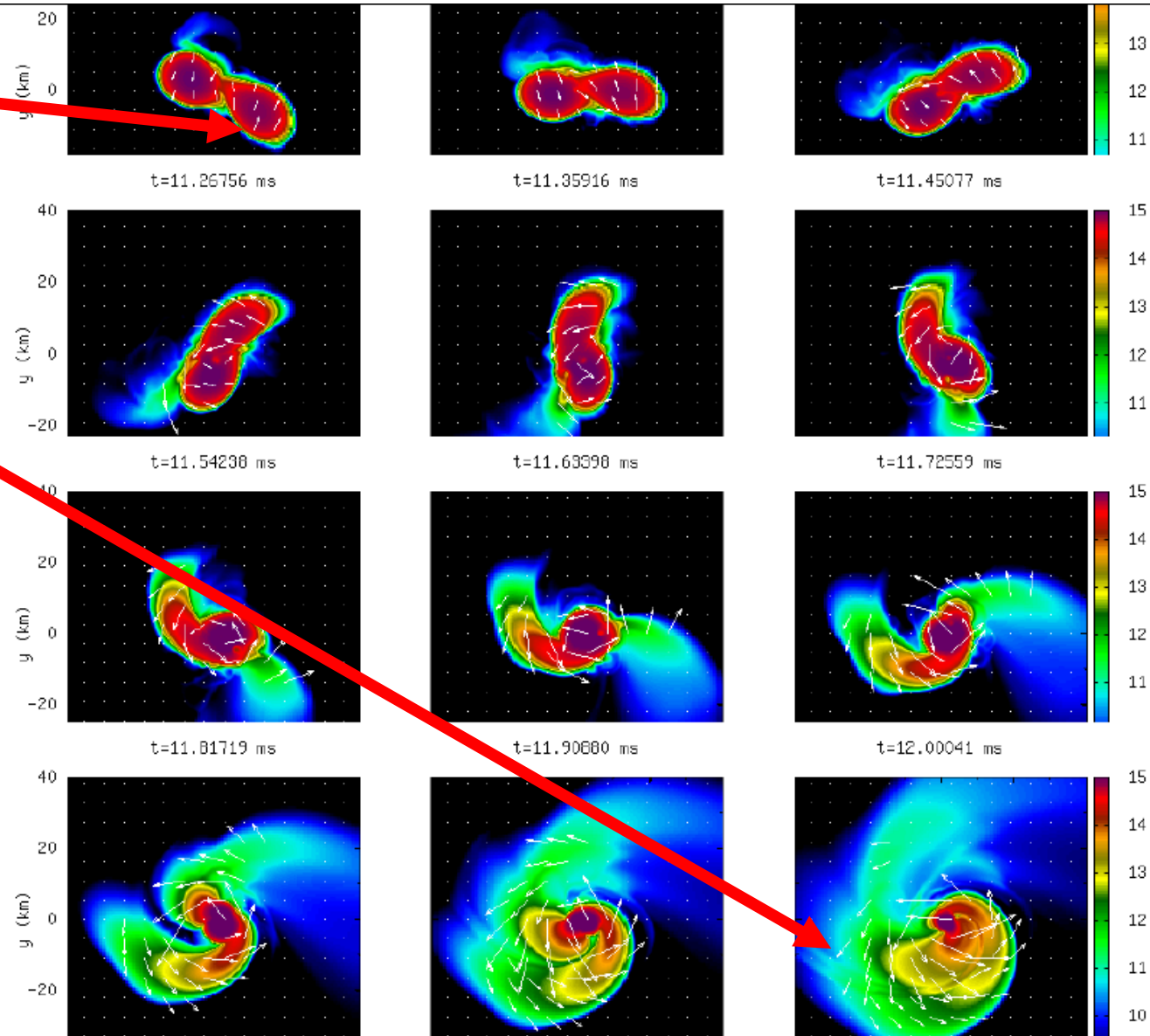
- ▶ Goriely et al. 2011; Bauswein et al. 2013
 - ▶ Approx. GR SPH sim. **without** weak interactions
 - ▶ No way to change $Y_e \Rightarrow$ ejecta remains n-rich (initial low Y_e)
 - ▶ See also post-process calculation of weak interactions
- ▶ Korobkin et al. 2012; Rosswog et al. 2013
 - ▶ Newtonian SPH sim. with neutrino
 - ▶ **tidal mass ejection (explained in the next slide)** of ‘pure’ neutron star matter
- ▶ Ejecta is very n-rich with $Y_e < 0.1$



Mass ejection from BNS merger (1) : Tidal torque + centrifugal force

Hotokezaka et al. (2013)

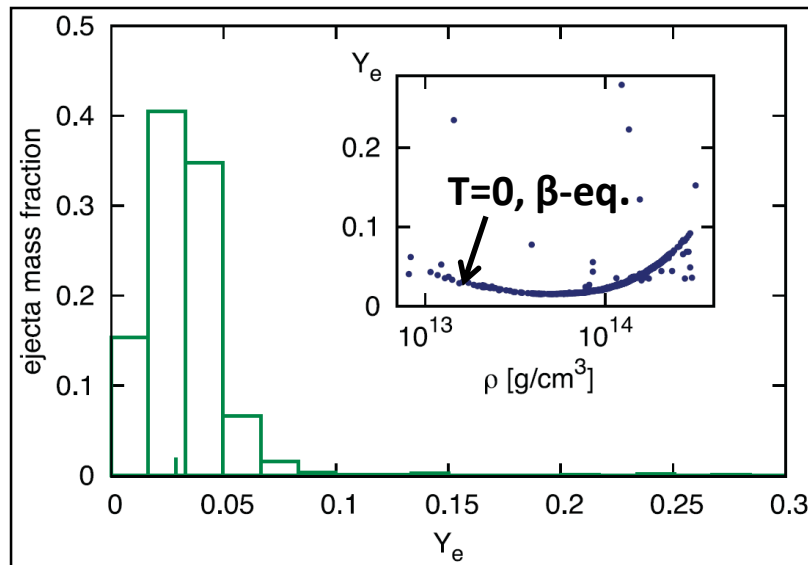
- ▶ Less massive NS is tidally deformed
- ▶ Angular momentum transfer by spiral arm and swing-by
- ▶ A part of matter is ejected along the orbital plane
- ▶ reflects low Y_e of cold NS (β -eq. at $T \sim 0$), no shock heating, rapid expansion (fast T drop), no time to change Y_e by weak interactions



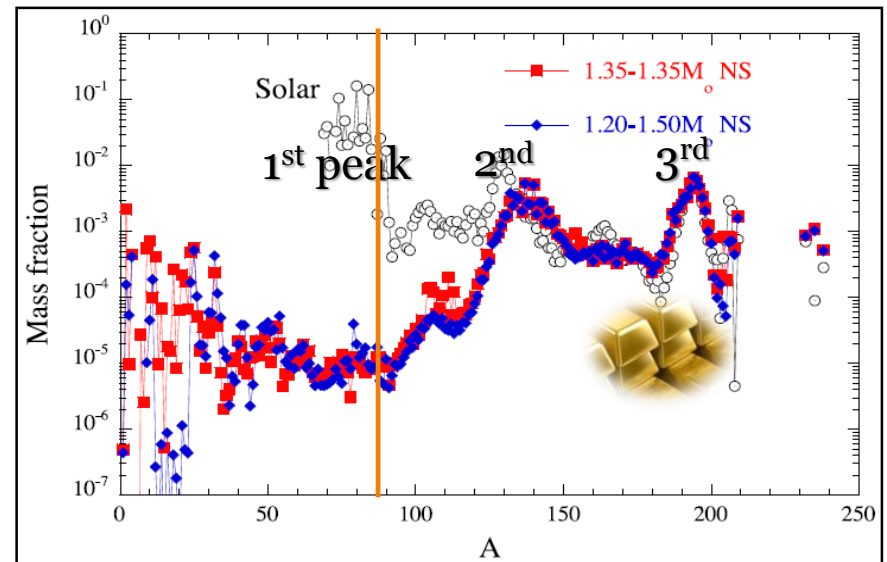
▶ Density contour
[$\log(\text{g/cm}^3)$]

From the ‘Universality’ point of view : NS-NS merger ejecta: too neutron-rich ?

- ▶ **Korobkin et al. 2012; Rosswog et al. 2013; see also Goriely et al. 2011**
 - ▶ tidal mass ejection of ‘pure’ neutron star matter (very n-rich) with $Y_e < 0.1$
 - ▶ Y_e is that of $T=0$, β -equilibrium
 - ▶ strong r-process with fission recycling only 2nd ($A \sim 130$; $N=82$) and 3rd ($A \sim 195$; $N=126$) peaks are produced (few nuclei in $A=90-120$)
 - ▶ the resulting abundance pattern does not satisfy universality in $A=90-120$



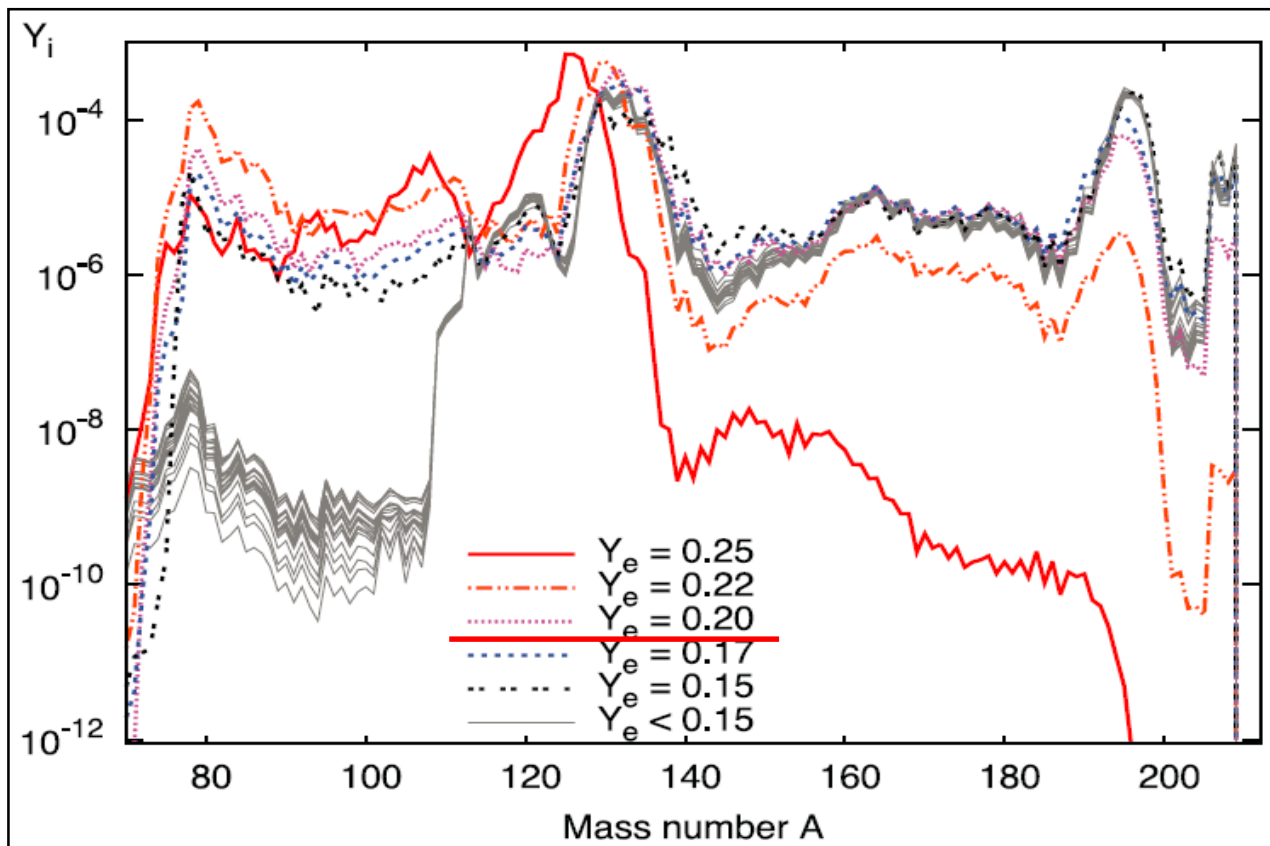
Korobkin et al. (2012) MNRAS 426 1940



Goriely et al. (2011) ApJL 738 32

How to satisfy the universality

- ▶ **Electron fraction (Y_e) is a key parameter : $Y_e \sim 0.2$ is critical threshold**
 - ▶ $Y_e < 0.2$: strong r-process \Rightarrow nuclei with $A > 130$ (the pattern is robust)
 - ▶ $Y_e > 0.2$: weak r-process \Rightarrow nuclei with $A < 130$ (for larger Y_e , nuclei with smaller A)



We need ejecta with higher Y_e

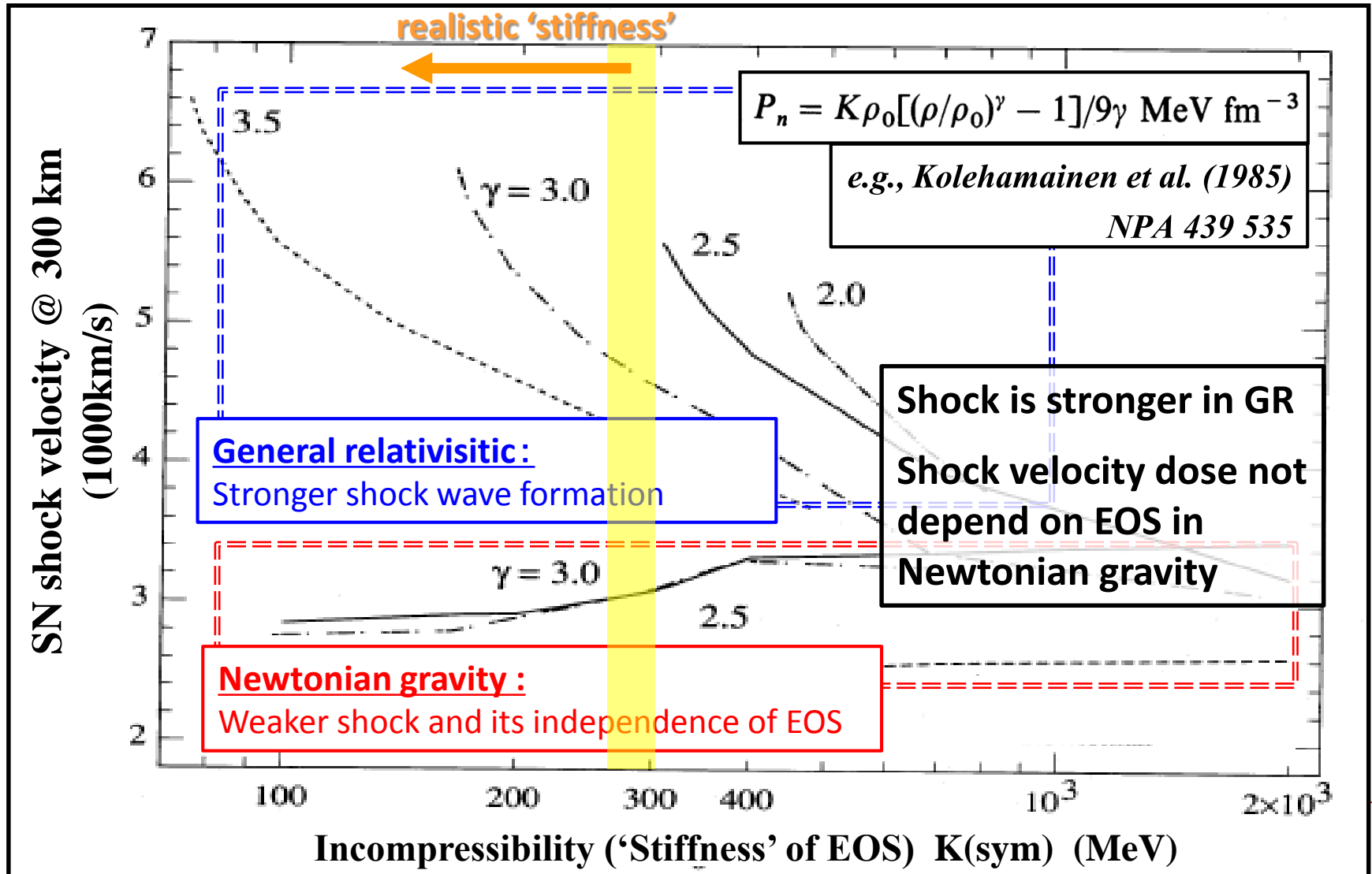
How to satisfy the universality

- ▶ Introduce new ejecta components
 - ▶ Neutrino driven winds from the remnant system
 - ▶ Dessart et al. (2009); Grossman et al. (2014); Perego et al. (2014); Just et al. (2015)
 - ▶ late time disk/torus disintegration
 - ▶ Fernandez & Metzger (2013)
- ▶ Take into account effects of both GR and weak interaction in the dynamical ejecta (this talk)



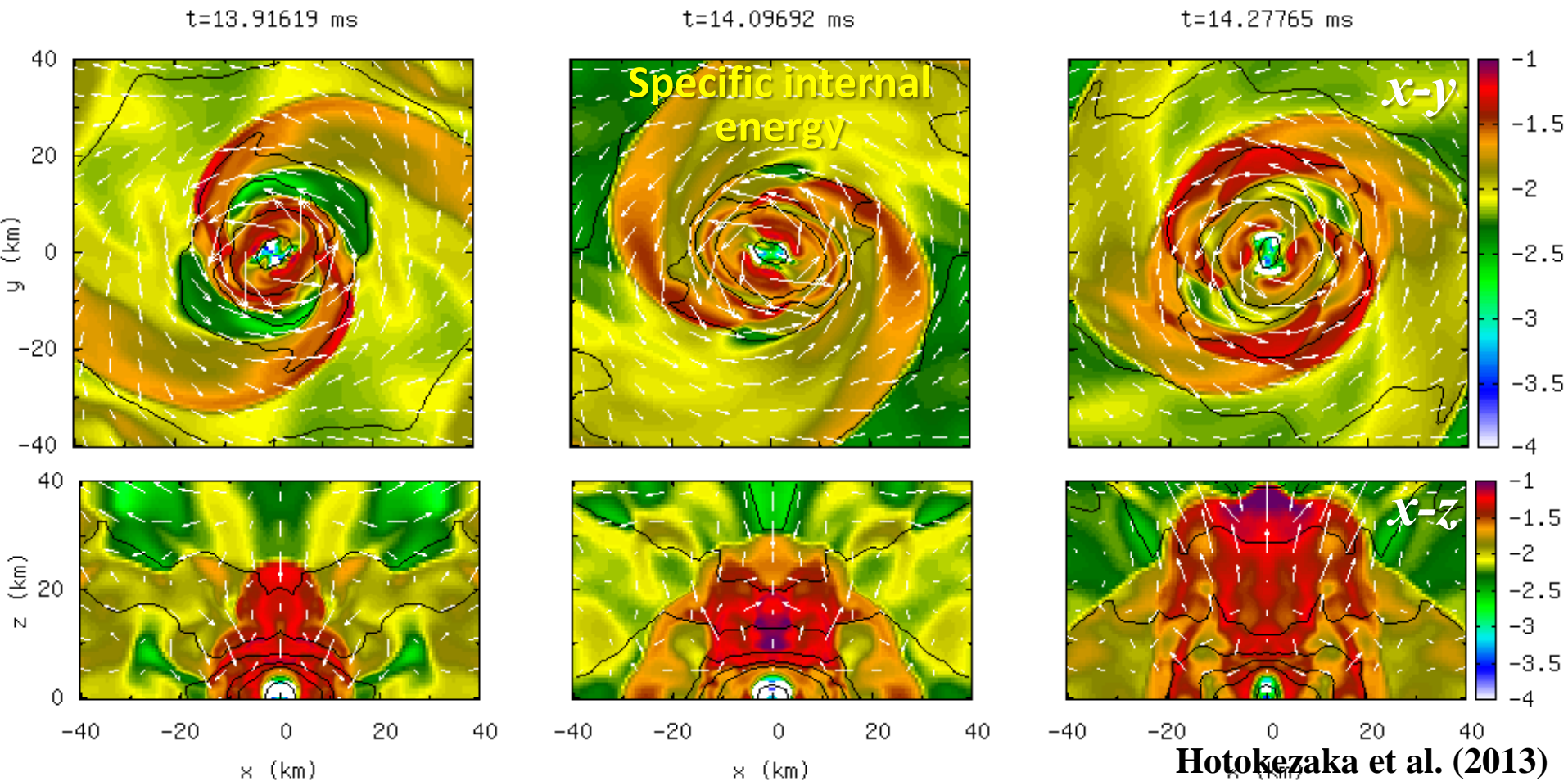
What will change if you include GR and microphysics (1) : Stronger shock in GR

van Riper (1988) *ApJ* 326 235



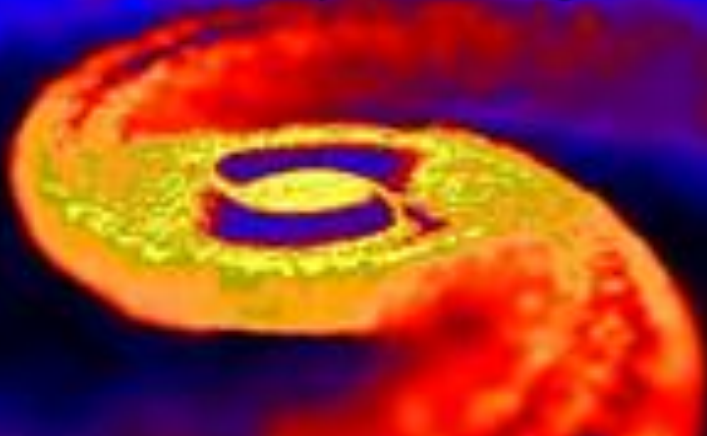
Mass ejection from BNS merger (2): Shock driven components

- ▶ Shocks occur due to oscillations of massive NS and collisions of spiral arms
- ▶ Isotropic mass ejection, higher temperature (weak interactions set in)



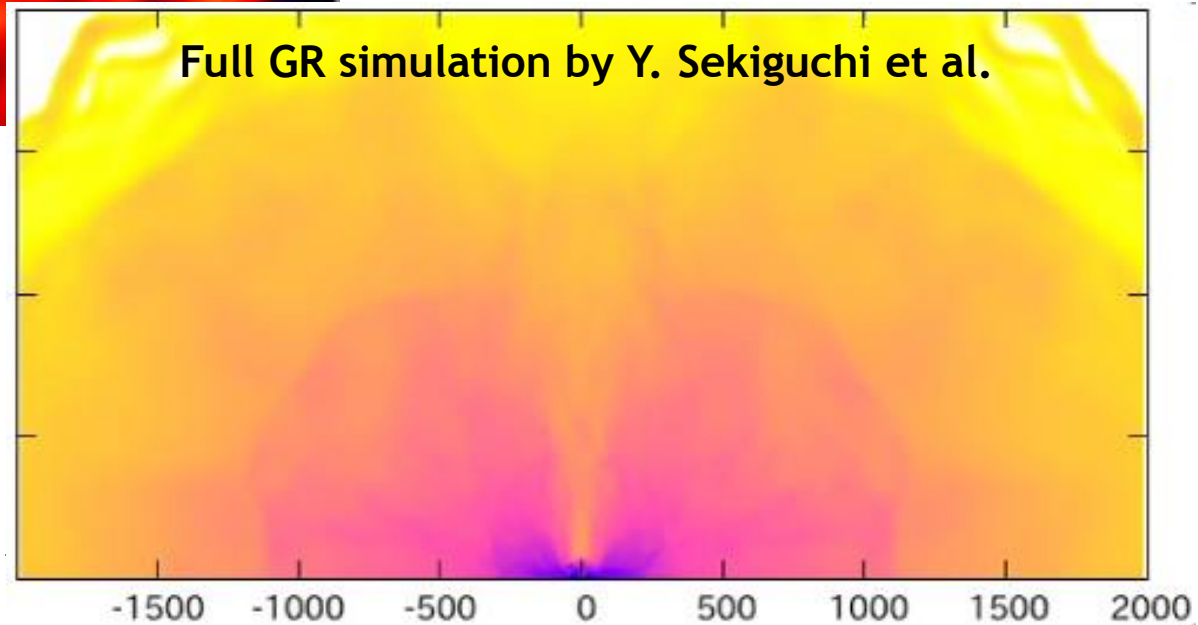
What will change if you include GR and microphysics (1) : Stronger shock in GR

Newtonian simulation by S. Rosswog et al.



Isotropic component less dominant (shock-driven) in Newtonian simulation
Only the tidal component

Full GR simulation by Y. Sekiguchi et al.



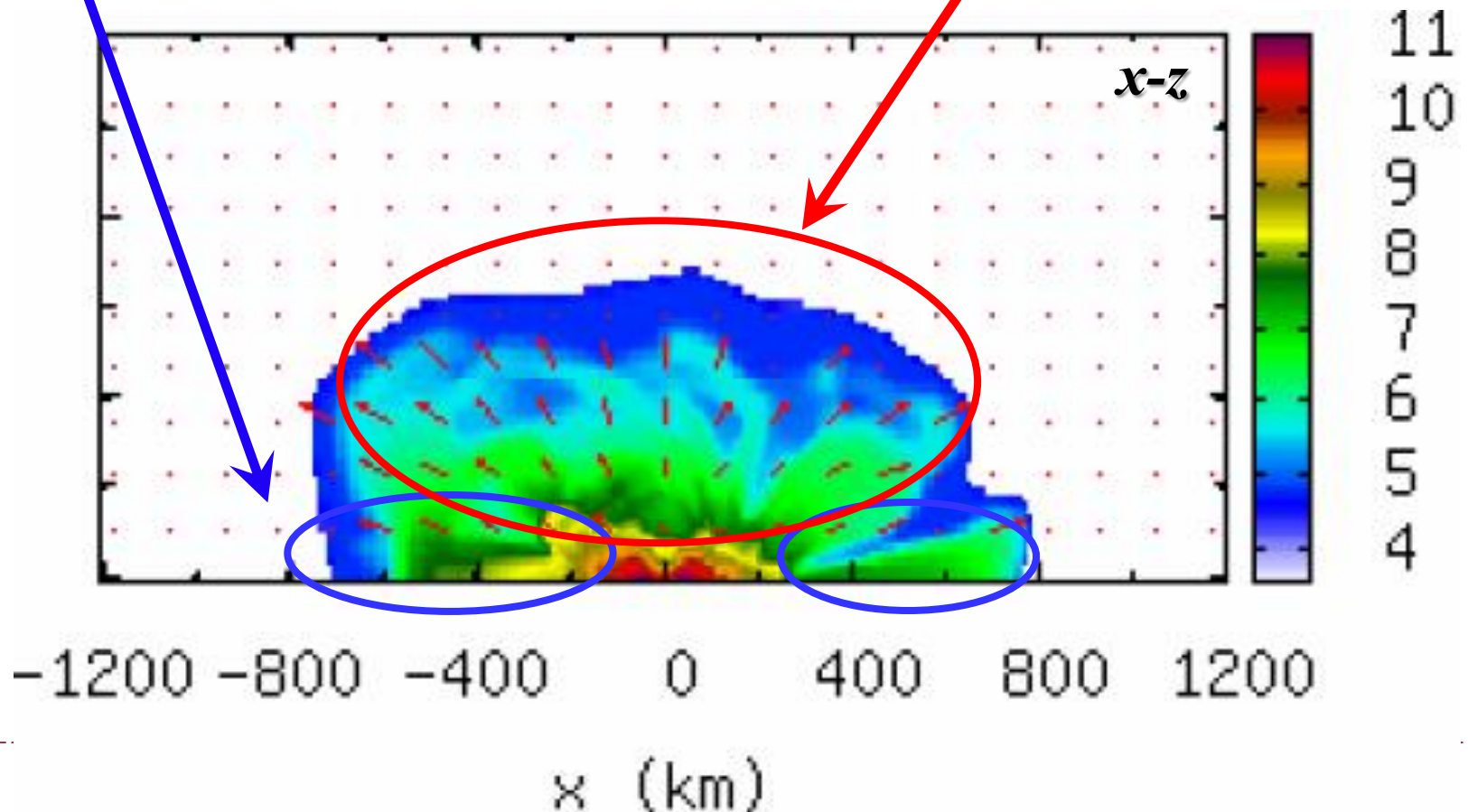
What will change if you include GR and microphysics (2) : Ye can change via weak interaction

▶ Driven by tidal interactions

Consists of cold NS matter in β -equilibrium \Rightarrow **low Ye and T**

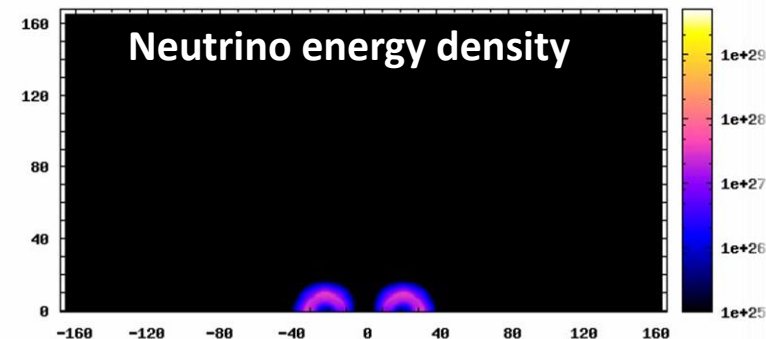
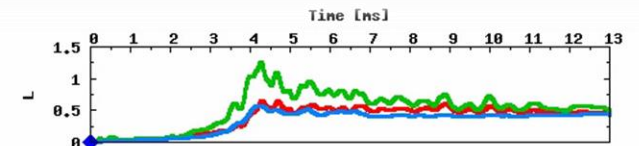
▶ Driven by shocks

Consists of shock heated matter
higher temperature \Rightarrow
Weak interaction can change Ye



Previous studies and our study

- ▶ **Korobkin et al. 2012** : Newtonian SPH simulations with neutrinos
- ▶ **Bauswein et al. 2013**: Relativistic SPH simulations with many EOS but without neutrons
- ▶ **This Study** : Full GR, approximate gray radiation hydrodynamics simulation with [multiple EOS and neutrinos \(brief summary of code is in appendix of lecture note\)](#)
 - ▶ Einstein's equations: Puncture-BSSN/Z4c formalism
 - ▶ GR radiation-hydrodynamics (*neutrino heating can be approximately treated*)
 - ▶ Advection terms : Truncated **Moment scheme** (*Shibata et al. 2011*)
 - ▶ EOS : any tabulated EOS with 3D smooth connection to Timmes EOS
 - ▶ gray or multi-energy but advection in energy is not included
 - ▶ Fully covariant and relativistic M-1 closure
 - ▶ Source terms : two options
 - ▶ **Implicit treatment : Bruenn's prescription**
 - ▶ **Explicit treatment : trapped/streaming ν 's**
 - e-captures: thermal unblocking/weak magnetism; NSE rate
 - Iso-energy scattering : recoil, Coulomb, finite size
 - e^\pm annihilation, plasmon decay, bremsstrahlung
 - diffusion rate (Rosswog & Liebendoerfer 2004)
 - two (beta- and non-beta) EOS method
 - ▶ Lepton conservation equations



Adopted finite-temperature EOS

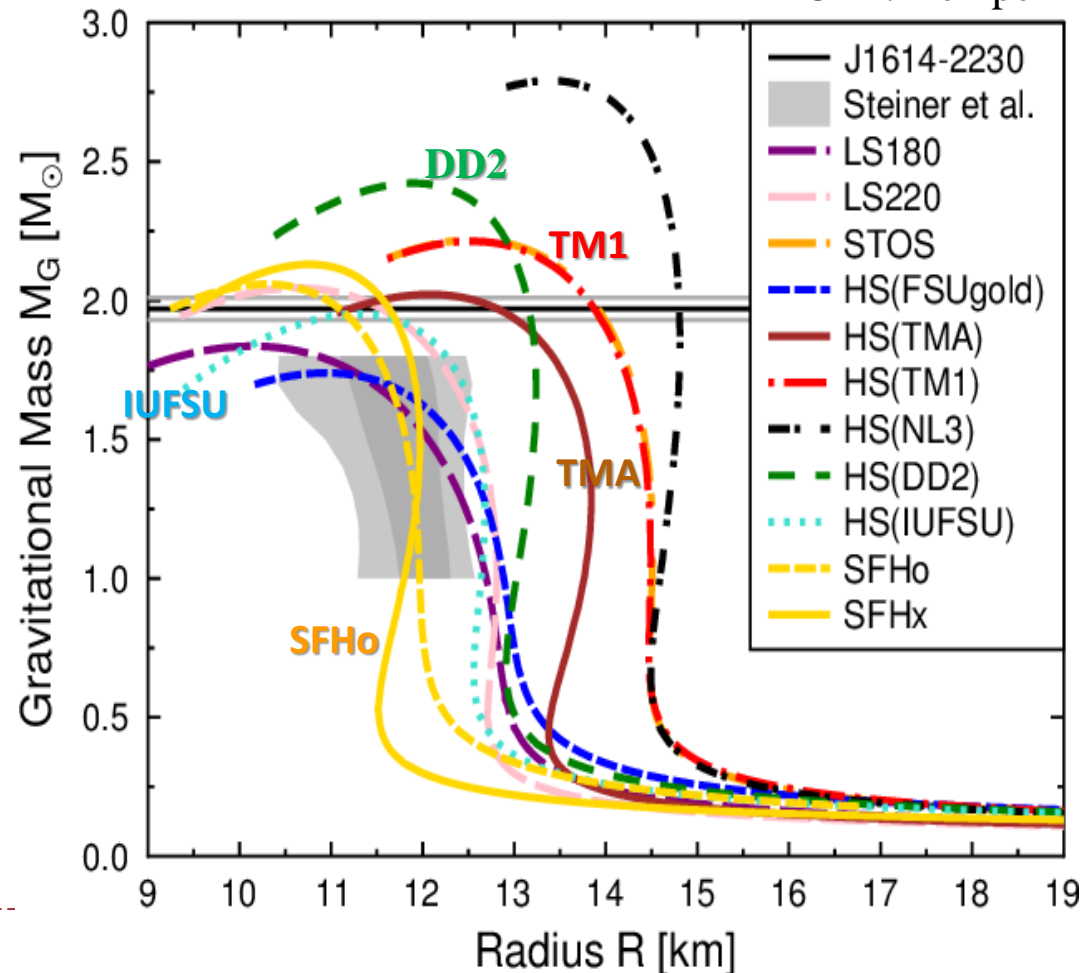
▶ Multi-EOS study (Thanks to M. Hempel)

See also, Bauswein et al. (2013);
Just et al. (2014)

▶ Adopted EOS

- ▶ 14.5km ▶ **TM1 (Shen EOS)**
 - ▶ TMA
 - ▶ 13.2km ▶ **DD2**
 - ▶ IUFSU
 - ▶ 11.8km ▶ **SFHo**
- Consistent with
- ▶ NS radius estimation
 - ▶ Chiral effective theory

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(Expected) Mass ejection mechanism & EOS

▶ 'Stiffer EOS'

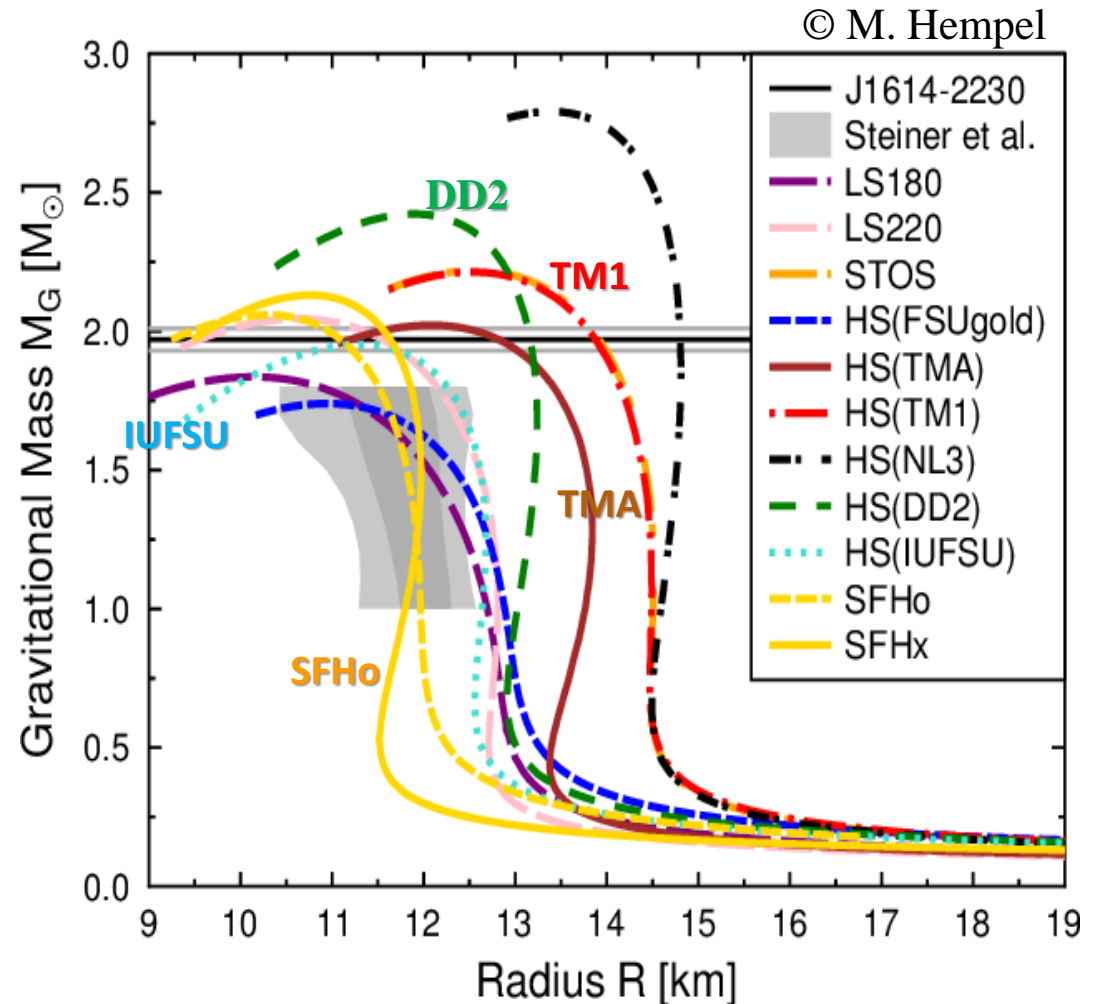
- ▶ $\Leftrightarrow R_{\text{NS}}$: larger
- ▶ **TM1, TMA**
- ▶ Tidal-driven dominant
- ▶ **Ejecta consist of low T & Y_e NS matter**

▶ 'Intermediate EOS'

- ▶ **DD2**

▶ 'Softer EOS'

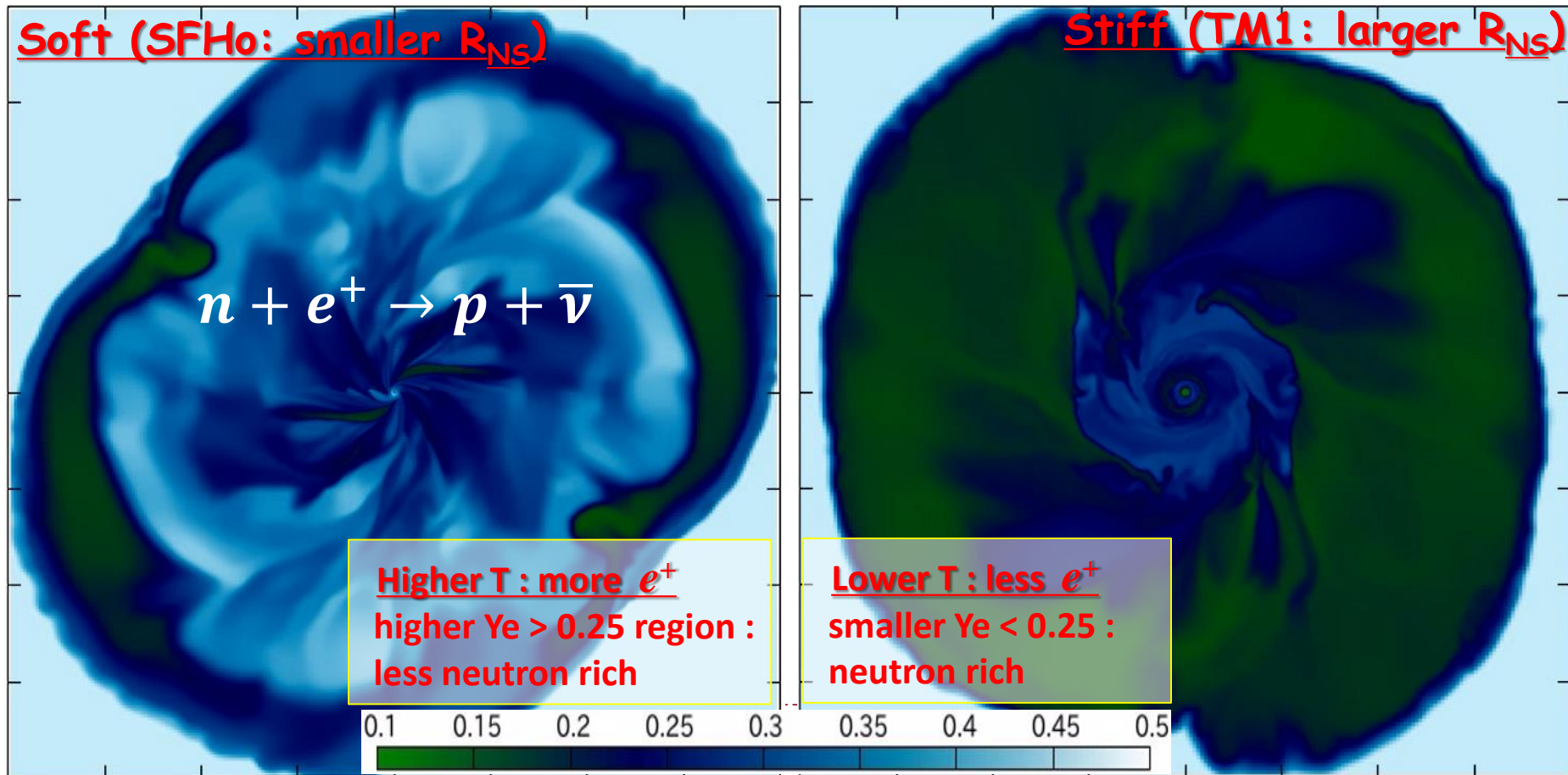
- ▶ $\Leftrightarrow R_{\text{NS}}$: smaller
- ▶ **SFHo, IUFSU**
- ▶ Tidal-driven less dominant
- ▶ Shock-driven dominant
- ▶ **Y_e can change via weak processes**



▶ See also, Bauswein et al. (2013); Just et al. (2014)

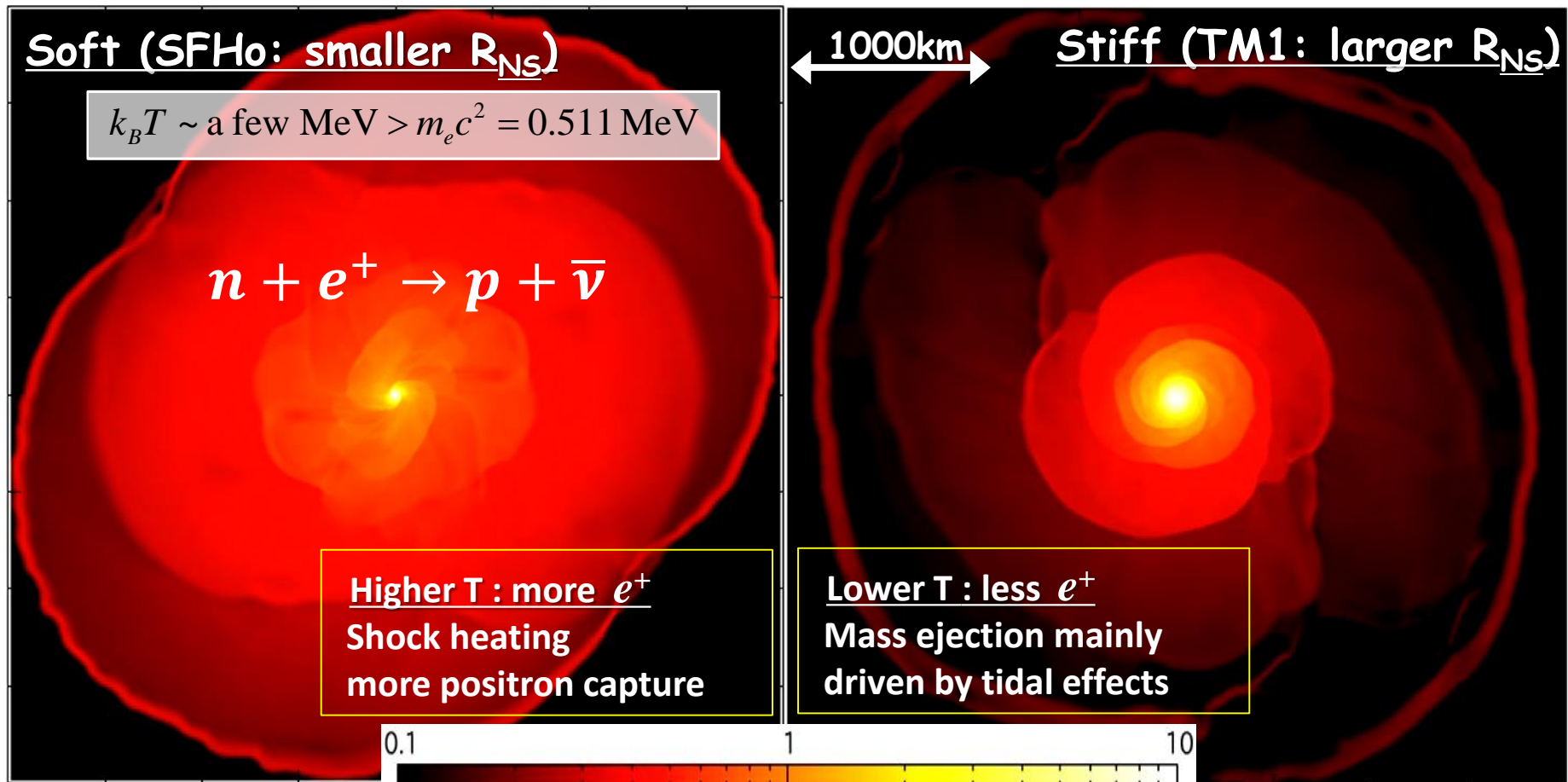
Soft(SFHo) vs. Stiff(TM1): Ejecta $Y_e = 1 - Y_n$

- ▶ Soft (SFHo): In the shocked regions, $Y_e \gg 0.2$ by weak processes
- ▶ Stiff (TM1): Y_e is low as < 0.2 (only strong r-process expected)



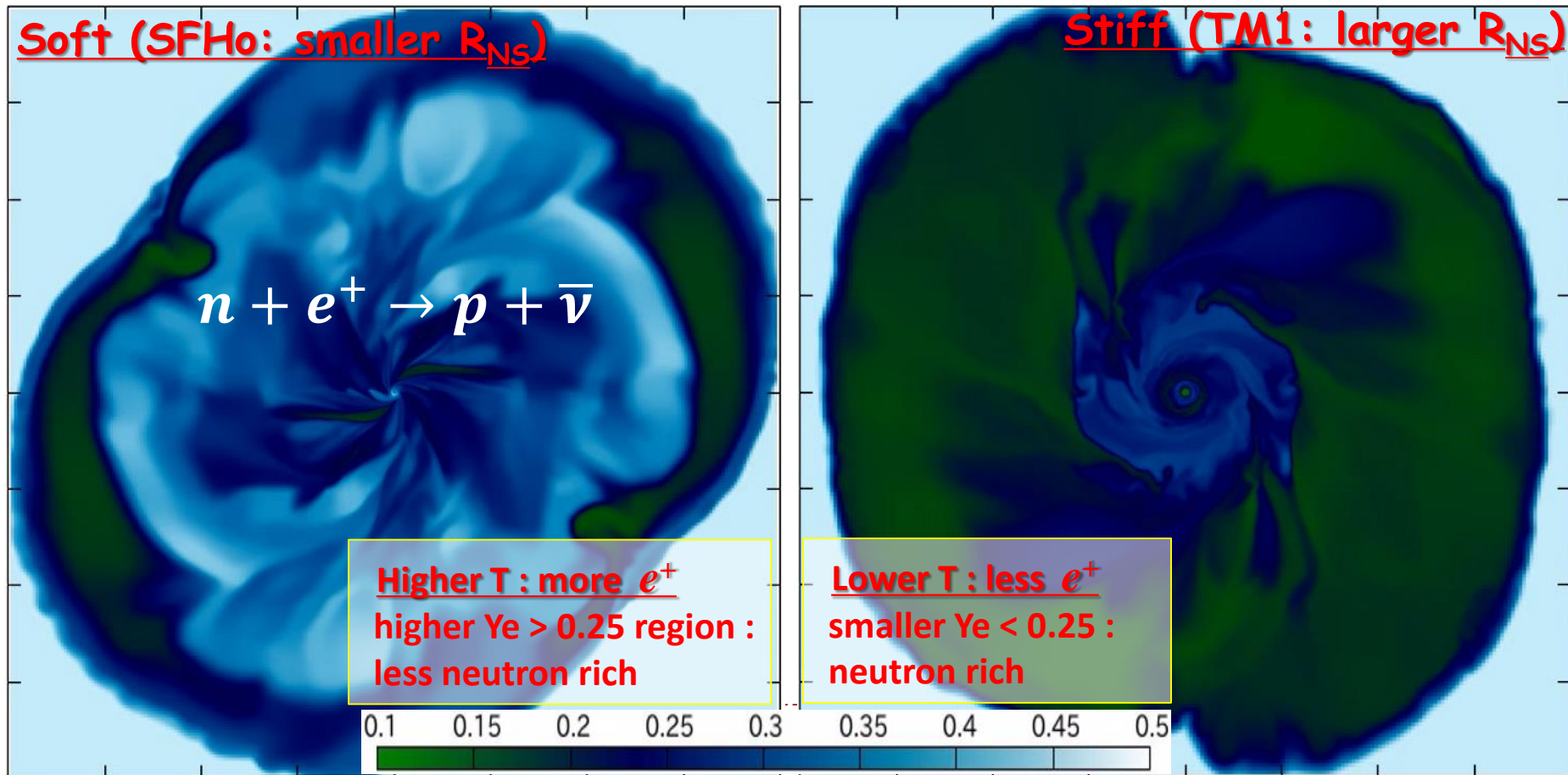
Soft(SFHo) vs. Stiff(TM1): Ejecta temperature

- ▶ Soft (SFHo): temperature of unbound ejecta is higher (as 1MeV) due to the shock heating, and produce copious positrons
- ▶ Stiff (TM1): temperature is much lower

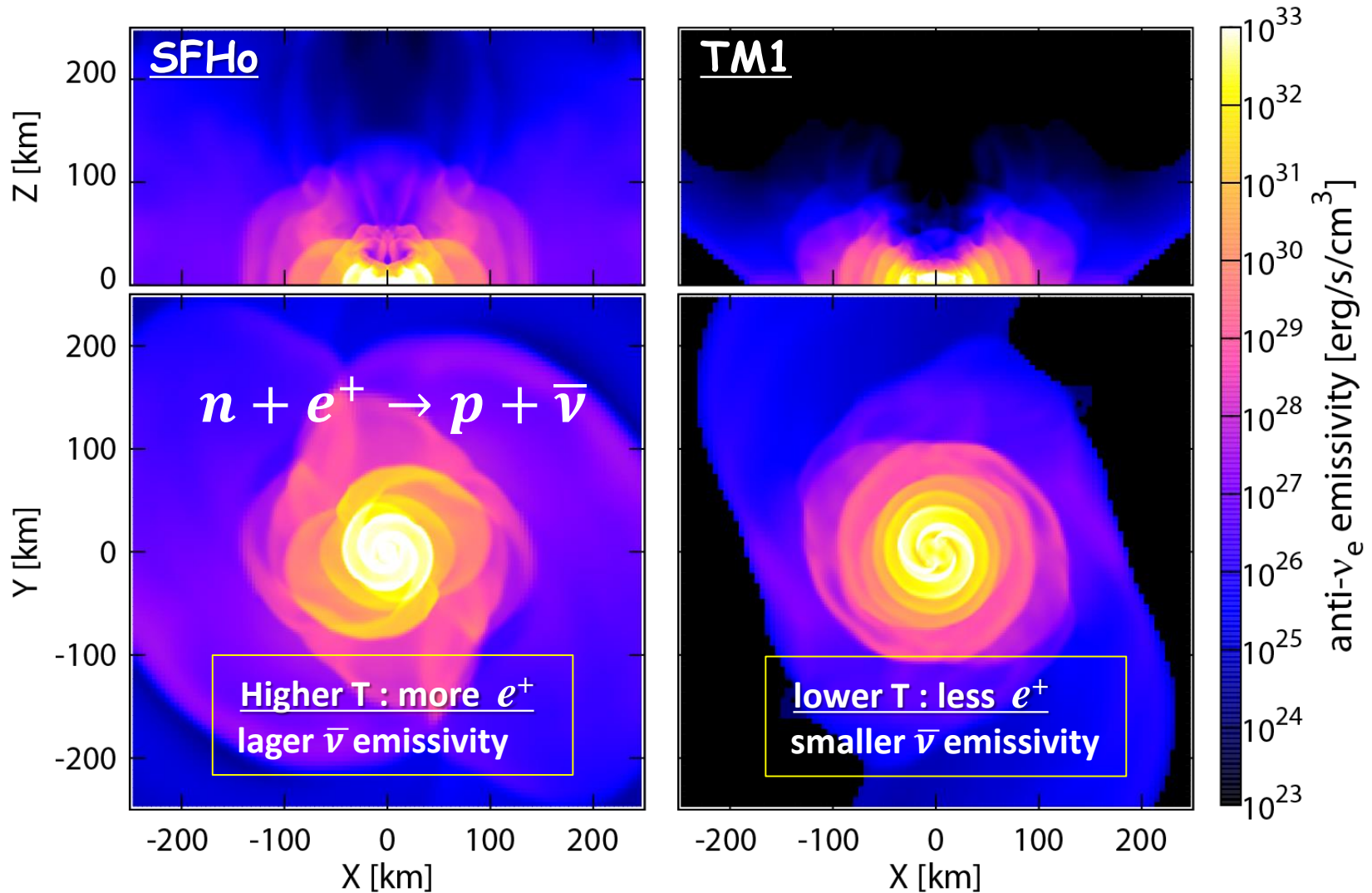


Soft(SFHo) vs. Stiff(TM1): Ejecta $Y_e = 1 - Y_n$

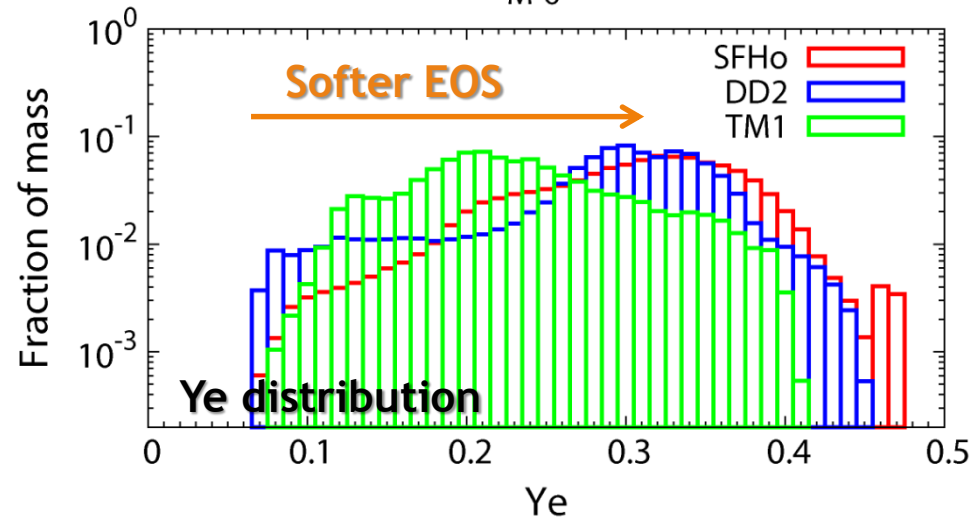
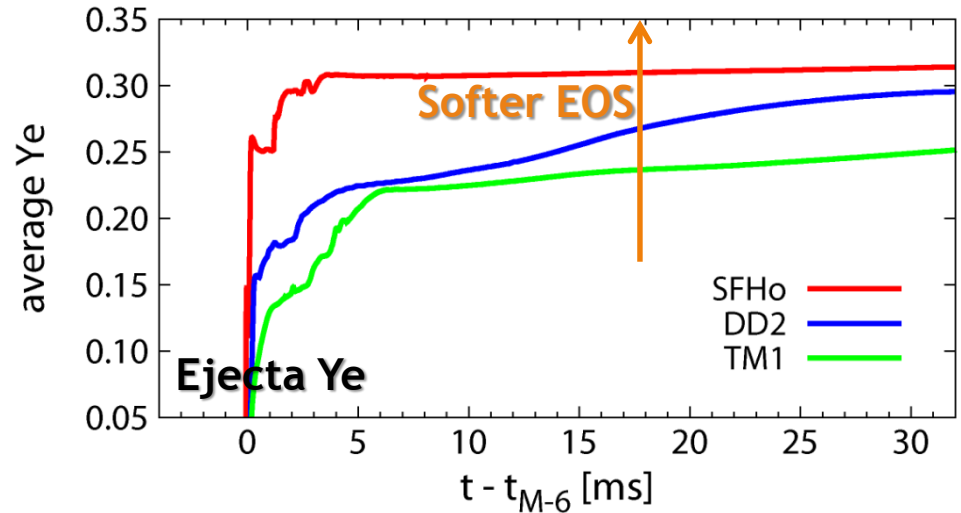
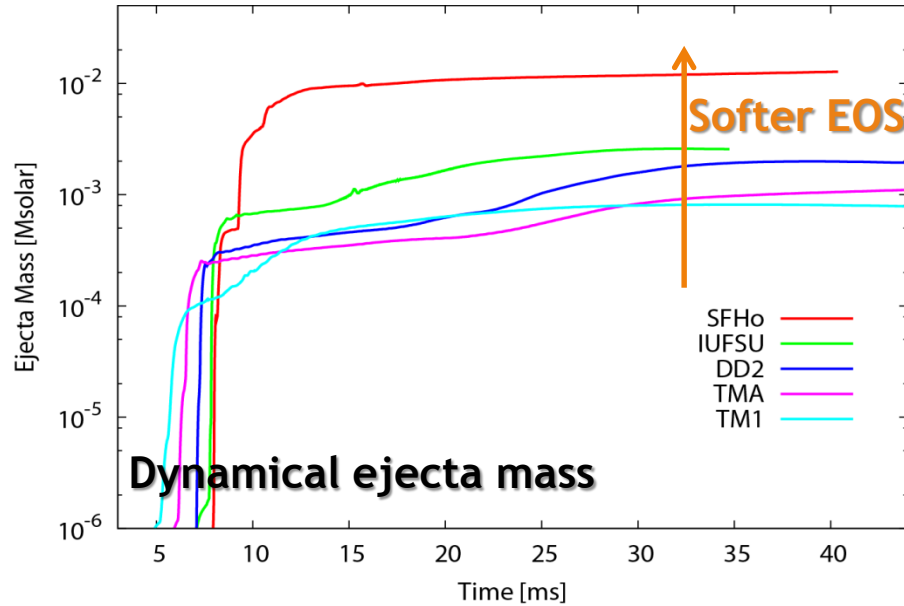
- ▶ Soft (SFHo): In the shocked regions, $Y_e \gg 0.2$ by weak processes
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SFHo vs. TM1: $\bar{\nu}_e$ emissivity



EOS dependence : 1.35-1.35 NS-NS

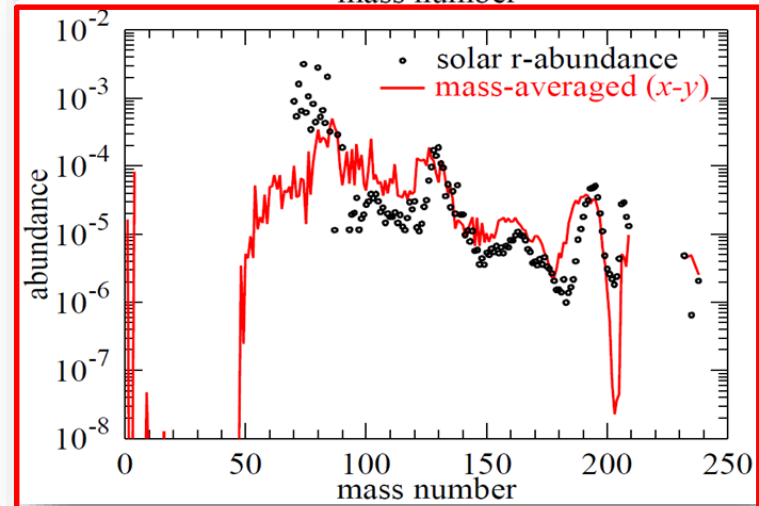
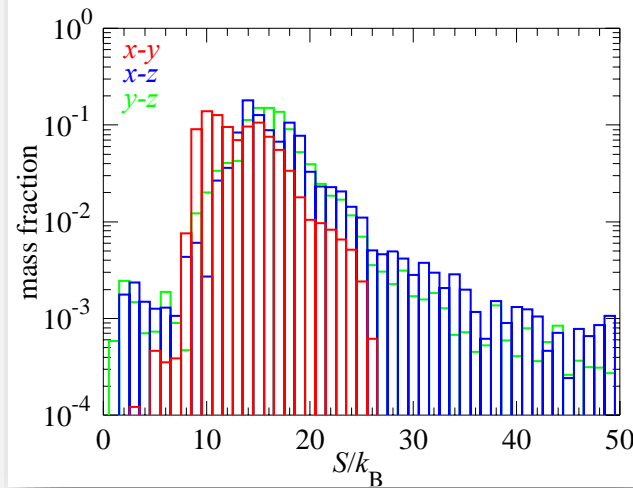
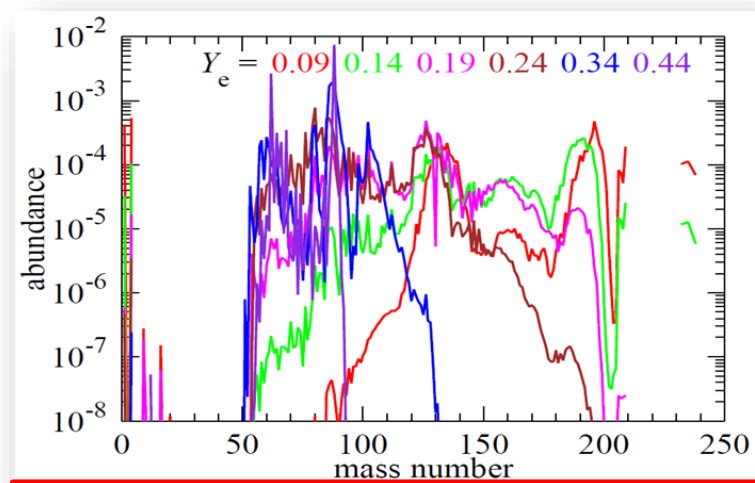
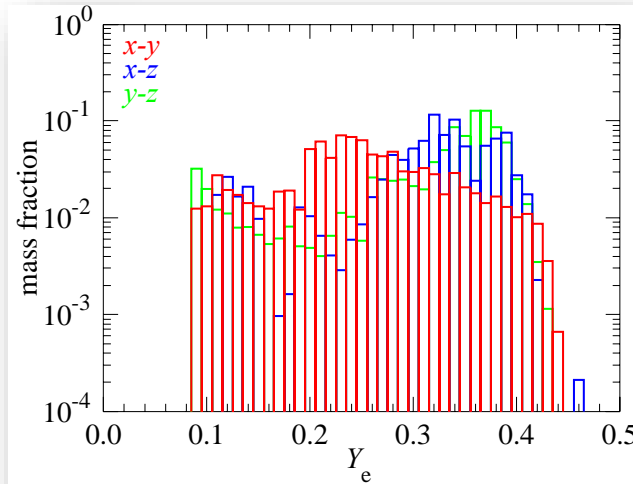


- ▶ **Mej is larger for softer EOS**
Consistent with piecewise-polytrope studies
- ▶ **Only SFHo will give Mej ~ 0.01 Msun**
 - ▶ a value required by the total amount of r-process elements and flux of the 'kilonova' event (GRB 130603B)
- ▶ If BNS is the origin, EOS should be soft



Achievement of the universality

(soft EOS (SFHo), equal mass (1.35-1.35))

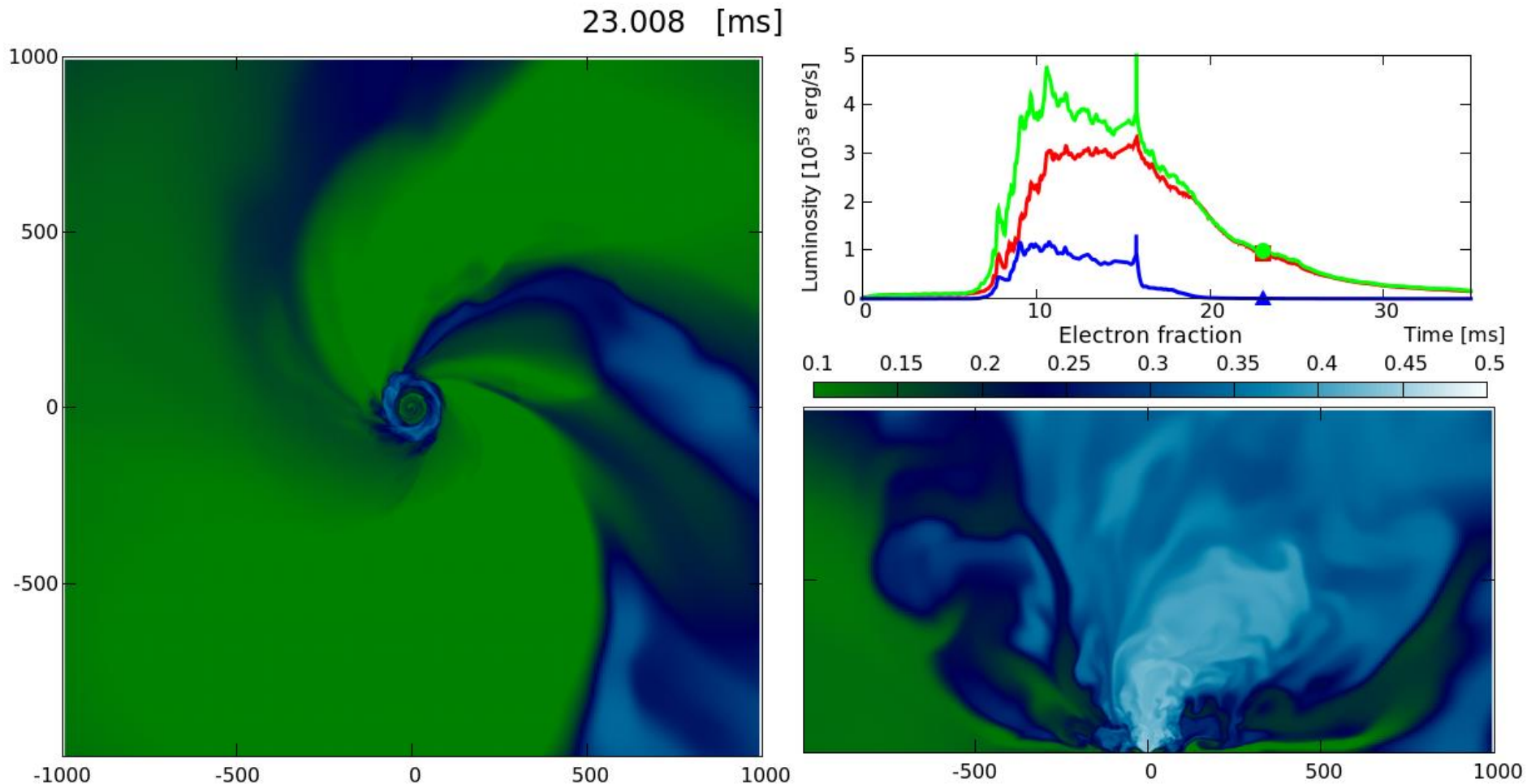


► The Y_e -distribution histogram has a broad, flat structure (*Wanajo, Sekiguchi, et al. (2014).*)

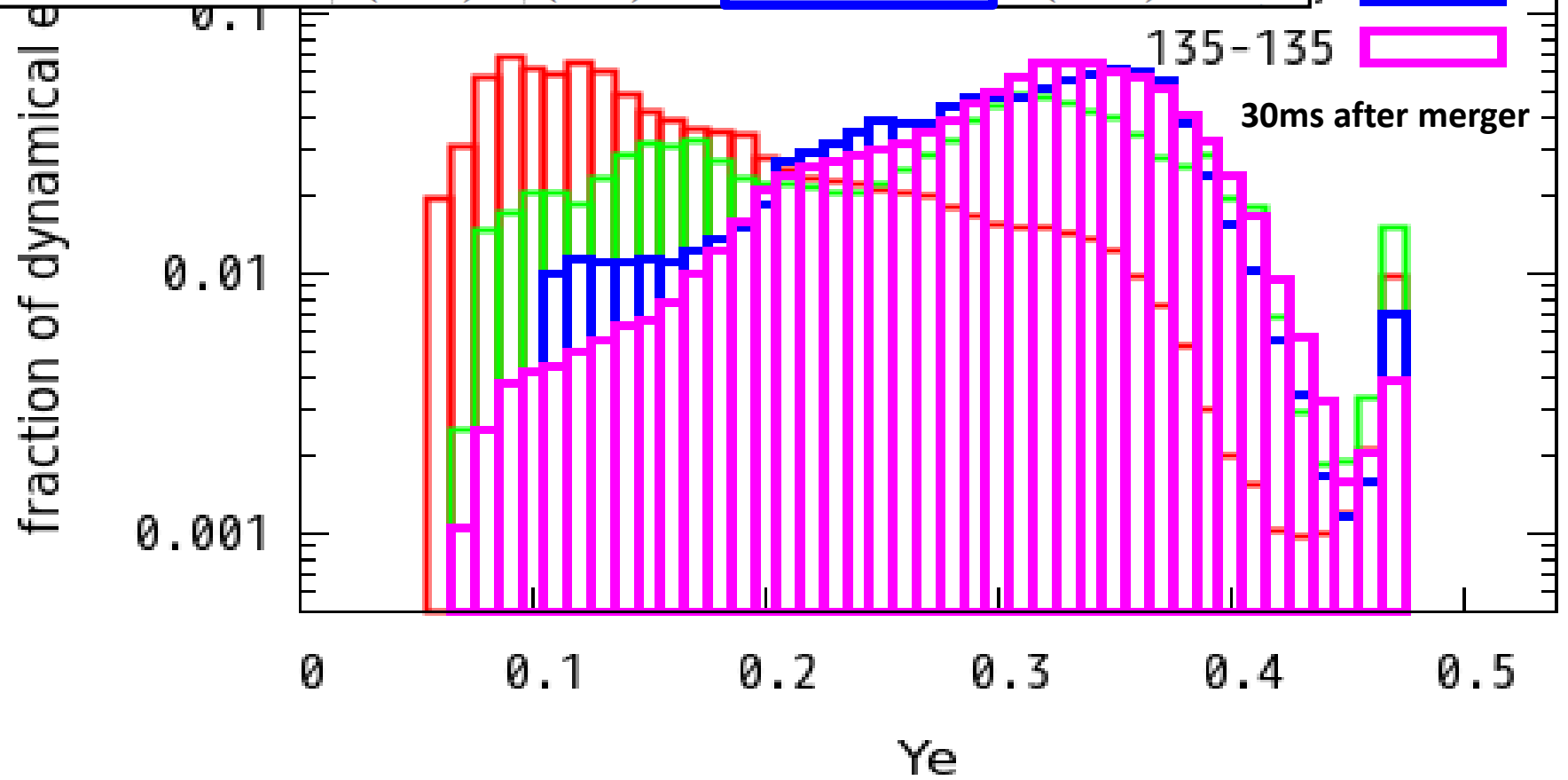
- Mixture of all Y_e gives a good agreement with the solar abundance !
- Robustness of Universality (dependence on binary parameters)

Unequal mass NS-NS system: SFHo1.25-1.45

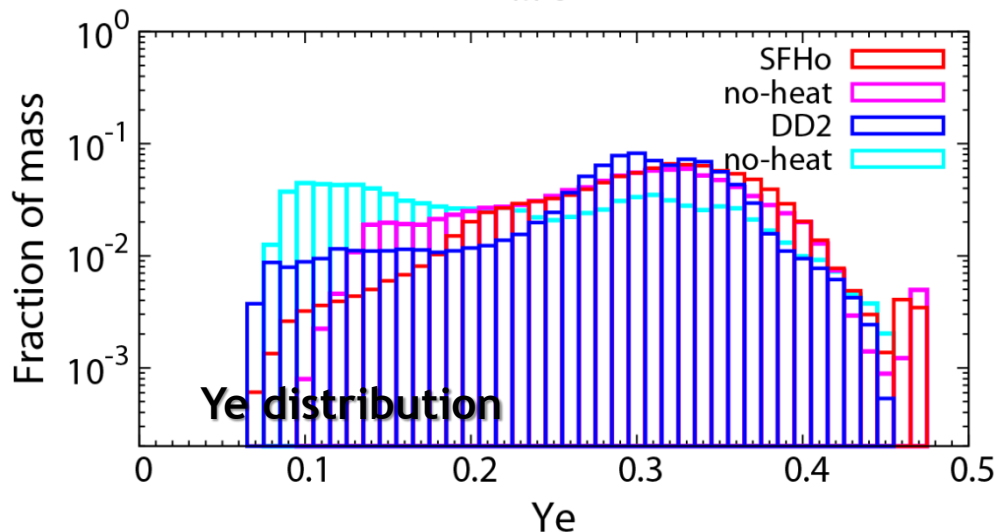
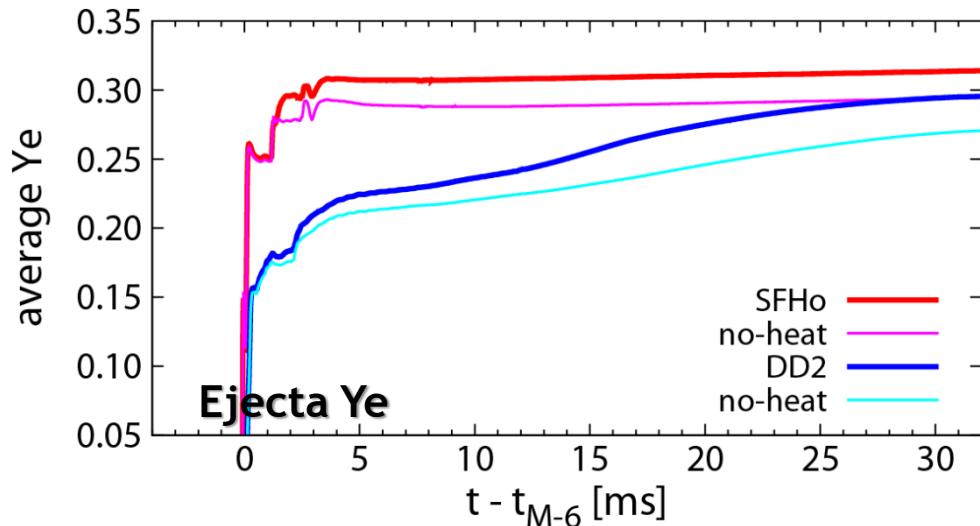
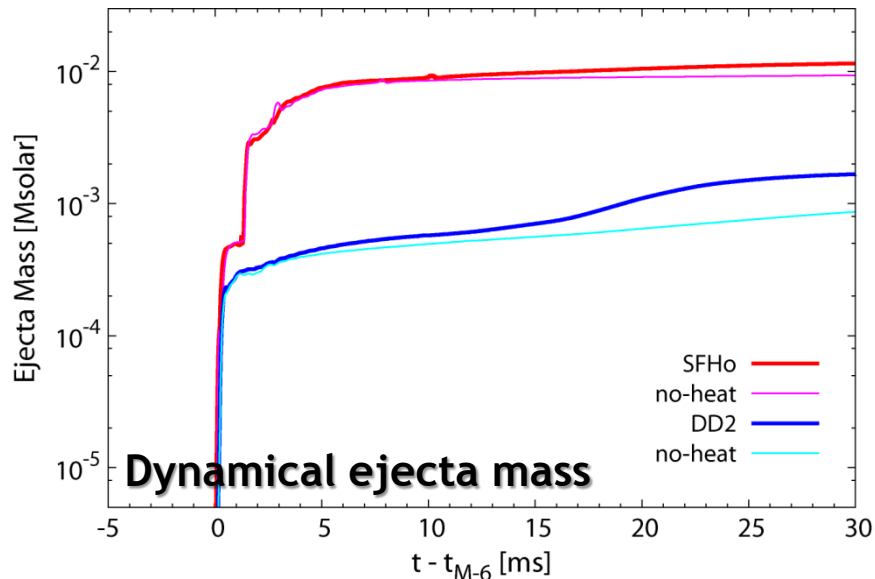
- ▶ Orbital plane : Tidal effects play a role, ejecta is neutron rich
- ▶ Meridian plane : shock + neutrinos play roles, ejecta less neutron rich



	PSR	$\log B(\text{G})$	P_{rot} (ms)	$M(M_{\text{sun}})$	T_{Mag}	T_{GW}
1.	B1913+16	10.4	59.0	1.441/1.387	1.0	3.0
2.	B1534+12	10.0	37.9	1.333/1.345	2.5	27
3.	B2127+11C	10.7	30.5	1.36/1.35	1.0	2.2
4.	J0737-3039	9.8/12.2	22.7/2770	1.34/1.25	2.0/0.5	0.86
5.	J1756-2251	9.7	28.5	1.34/1.23	4.0	17
6.	J1906+746	(12.2)	(144)	1.29/1.32	(<0.1)	3.1



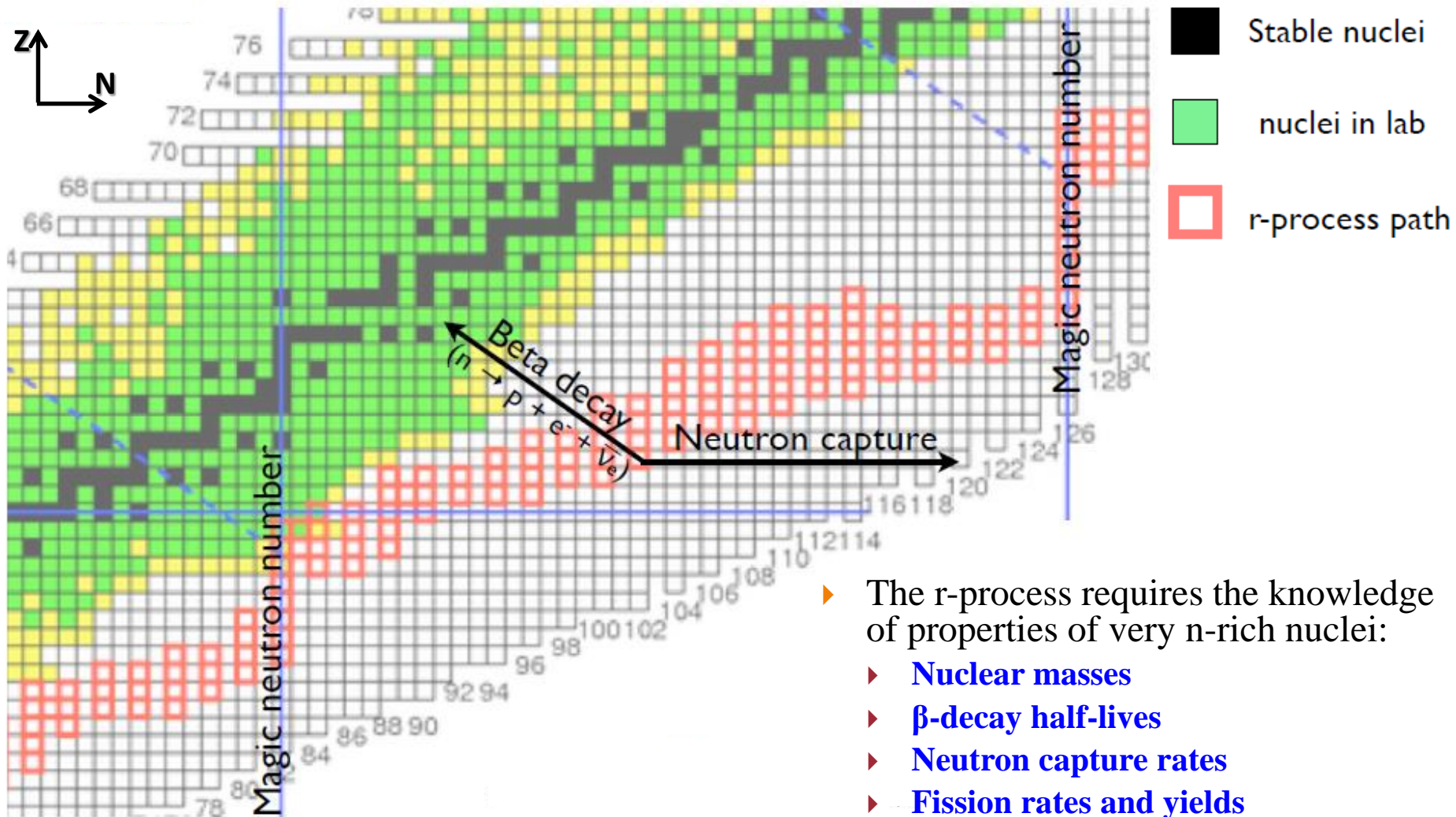
Importance of neutrino heating (absorption)



- ▶ Amount of ejecta mass can be increased order of 10^{-3} Msun
- ▶ Average Ye can change 0.02~0.03 depending on EOS : effect is stronger for stiffer EOS where HMNS survive in a longer time



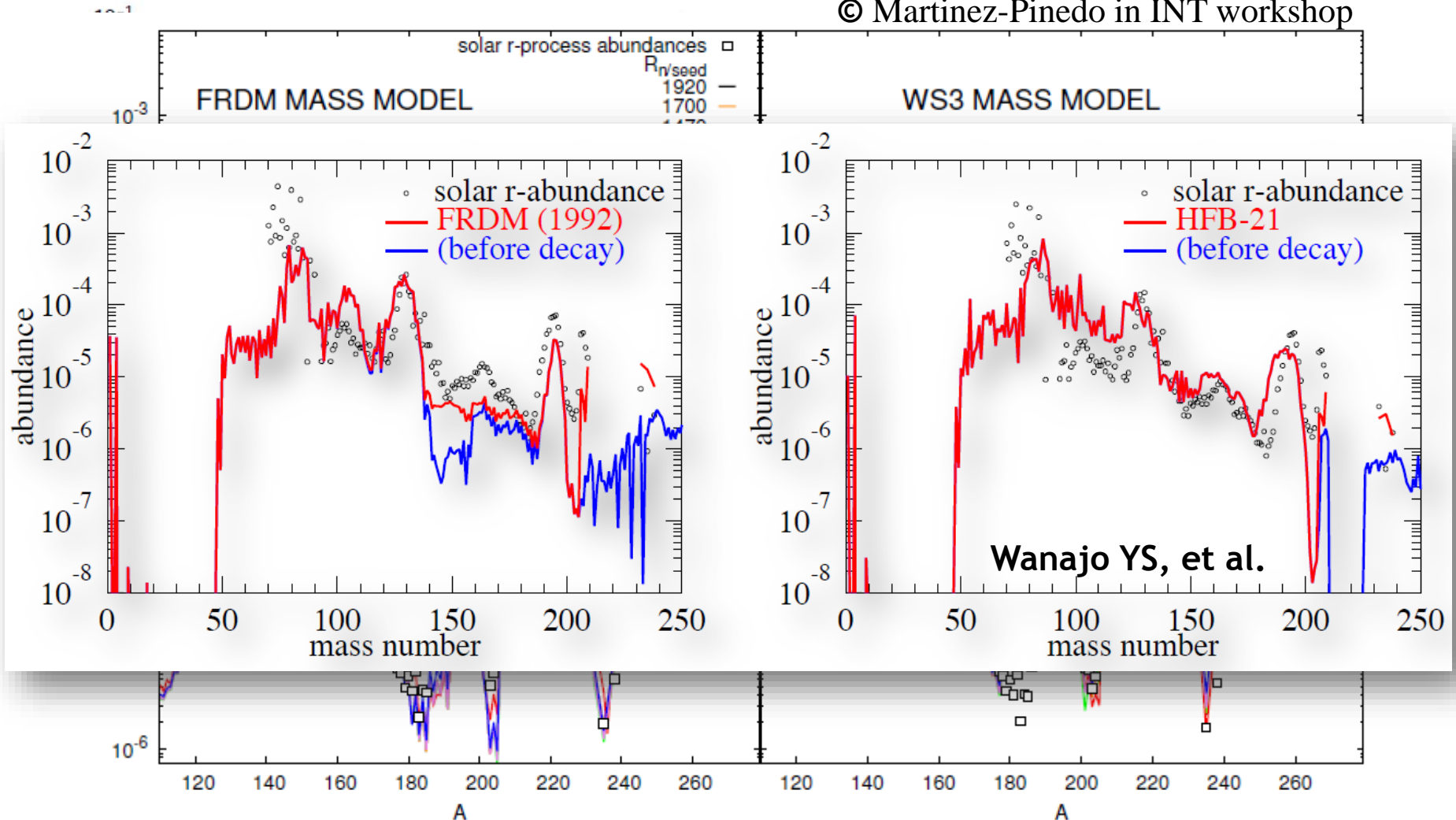
r-process nucleosynthesis: nuclear physics inputs



- ▶ The r-process requires the knowledge of properties of very n-rich nuclei:
 - ▶ **Nuclear masses**
 - ▶ **β -decay half-lives**
 - ▶ **Neutron capture rates**
 - ▶ **Fission rates and yields**

Dependence on mass model

© Martinez-Pinedo in INT workshop



Summary

- ▶ A long history about the origin of heavy elements
- ▶ Supernovae vs. Neutron star mergers
 - ▶ Supernova scenario is now theoretically and observationally disfavored
- ▶ Key observation: Universality
 - ▶ r-process cite synthesize elements with pattern which is similar to the observed solar pattern
 - ▶ Less diverse 'Initial condition' or some physical attracter (fission yields)
- ▶ General relativity and weak/strong interactions are key to resolve the problem
 - ▶ If EOS of NS is soft (like APR or SFHo), then it is strongly suggested that the origin of heavy elements are BNS mergers.



Final words:

Towards the first direct detection of GW

▶ The first detection of neutrinos

- ▶ Simultaneous observation of SN 1987A (EM counterpart to neutrino) was very important
- ▶ Similarly, EM counterpart to GW could play a role

▶ Possible EM counterpart to GW

- ▶ Short GRBs : likely to be collimated => Most of them are off axis and faint
 - ▶ SGRB111020A : $\theta_j \sim 3-8^\circ$ (Fong et al. 2012)
 - ▶ SGRB051121A : $\theta_j \sim 7^\circ$ (Burrows et al. 2006)
- ▶ We need 4π counterparts
- ▶ 'kilonova' like event from radioactive decay of r-process nuclei is one of promising candidates
- ▶ Many studies : Rosswog, Goriely, Bauswein, Metzger, Kasen, YS, Wanajo, Hotokezaka, Tanaka
- ▶ Studies based on full GR simulations on going in collaboration with Wanajo, Tanaka (Subaru group), ...



Kilonova / Macronova / r-process nova

- ▶ Merger ejecta will be neutron rich: rapid neutron capture (r-process) proceeds (Lattimer & Schramm 1974) : $n + (Z, N) \Rightarrow (Z, N+1)$
- ▶ Competition with the β -decay : $(Z, N+1) \Rightarrow (Z+1, N) + e + \bar{\nu}_e$
 - ▶ The r-process is very sensitive to how much neutrons are there, that is, to the electron fraction Y_e ($= Y_p = 1 - Y_n$) : we need microphysics !
- ▶ Then, **EM transients powered by radioactivity of the r-process elements** are expected (Li & Paczynski 1998; Kulkarni 2005; Metzger et al. 2010)

$$t_{\text{peak}} \sim 1 \text{ days} \left(\frac{v}{0.3c} \right)^{-1/2} \left(\frac{M}{0.01 M_{\text{solar}}} \right)^{1/2} \left(\frac{\kappa}{0.1 \text{ cm}^2 / \text{g}} \right)^{1/2}$$

$$L_{\text{peak}} \sim 10^{42} \text{ erg/s} \left(\frac{f}{10^{-6}} \right) \left(\frac{v}{0.3c} \right)^{1/2} \left(\frac{M}{0.01 M_{\text{solar}}} \right)^{1/2} \left(\frac{\kappa}{0.1 \text{ cm}^2 / \text{g}} \right)^{-1/2}$$

$$T_{\text{peak}}^{\text{eff}} \sim 10^4 \text{ K} \left(\frac{f}{10^{-6}} \right)^{1/4} \left(\frac{v}{0.3c} \right)^{-1/8} \left(\frac{M}{0.01 M_{\text{solar}}} \right)^{-1/8} \left(\frac{\kappa}{0.1 \text{ cm}^2 / \text{g}} \right)^{-3/8}$$



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 - ▶ The r-process is very sensitive to how much neutrons are there, that is, to the electron fraction Y_e ($= Y_p = 1 - Y_n$) : we need microphysics !
- ▶ Recent critical update : Opacities are dominated by lanthanoids : orders of magnitude (~ 100) larger (Kasen et al. 2013; Tanaka & Hotokezaka 2013)

$$t_{\text{peak}} \sim 10 \text{ days} \left(\frac{v}{0.3c} \right)^{-1/2} \left(\frac{M}{0.01M_{\text{solar}}} \right)^{1/2} \left(\frac{\kappa}{10 \text{ cm}^2 / \text{g}} \right)^{1/2}$$

1 day \Rightarrow 10 days

$$L_{\text{peak}} \sim 10^{41} \text{ erg/s} \left(\frac{f}{10^{-6}} \right) \left(\frac{v}{0.3c} \right)^{1/2} \left(\frac{M}{0.01M_{\text{solar}}} \right)^{1/2} \left(\frac{\kappa}{10 \text{ cm}^2 / \text{g}} \right)^{-1/2}$$

1/10 dimmer

$$T_{\text{peak}}^{\text{eff}} \sim 2 \times 10^3 \text{ K} \left(\frac{f}{10^{-6}} \right)^{1/4} \left(\frac{v}{0.3c} \right)^{-1/8} \left(\frac{M}{0.01M_{\text{solar}}} \right)^{-1/8} \left(\frac{\kappa}{10 \text{ cm}^2 / \text{g}} \right)^{-3/8}$$

Opt-UV \Rightarrow NIR



NS-NS vs. BH-NS

- ▶ kilonova and total mass of r-process element => **Mej > 0.01 Msun**
- ▶ **NS-NS** : **Soft EOS is necessary** (shocks play a role)
 - ▶ **Small diversity** in conditions before merger, $Mej \sim 0.01$ Msun will be universal within the typical mass range of NS-NS for soft EOS
- ▶ **BH-NS** : **Stiffer EOS is preferable** (tidal component is dominant)
 - ▶ **some diversity** is expected, because mass ejection (mostly tidal-driven) depends further on *mass and spin of BH* in addition to EOS

