

# SU(3)ゲージ理論における エントロピックC関数の測定

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collaboration with K. Nagata (KEK), Y.Nakagawa, A. Nakamura  
(Hiroshima U., RCNP) and V.I.Zakharov (Max Planck Inst.)

cf. arXiv:0911.2596 and 1104.1011:

Y.Nakagawa, A.Nakamura, S.Motoki and V.I.Zakharov



# メインテーマ

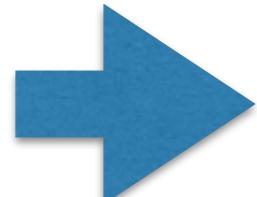
エンタングルメント・エントロピー(エントロピックC関数)から  
閉じ込めに対する新しい知見を得る

4次元ゲージ理論における  
量子エンタングルメントの非摂動論的な振る舞いは？

cf) arXiv:1508.07132[hep-lat] **today!**

Lattice study on QCD-like theory with exact center symmetry

T.Iritani, E.I and T.Misumi



関連した研究: 9月2日 (水) 河野さん

# Basic properties of E.E.

## entanglement entropy for quantum system

- how much a quantum state is entangled quantum mechanically
- d.o.f of the system
- quantum properties of the ground state for the system
- in finite T system, it gives a thermal entropy

## QCD theory ( $T=0$ )

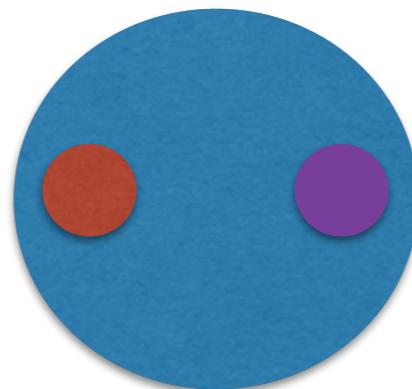
A color confinement changes the d.o.f of the system

microscopically

colorful  
(gluons)

$$\sim O(N_c^2)$$

$$\Lambda_{\text{QCD}}$$



macroscopically

colorless  
(singlet)

$$\sim O(1)$$

# Definition of the entanglement entropy

# Entanglement entropy (E.E.)

von Neumann entropy

$$S_{tot} = -\text{Tr} \rho_{tot} \log \rho_{tot}$$

density matrix

$$\rho_{tot} = |\Psi\rangle\langle\Psi| \quad |\Psi\rangle : \text{pure ground state}$$

decompose total Hilbert space into two subsystems

$$\mathcal{H}_{tot} = \mathcal{H}_A \otimes \mathcal{H}_B$$

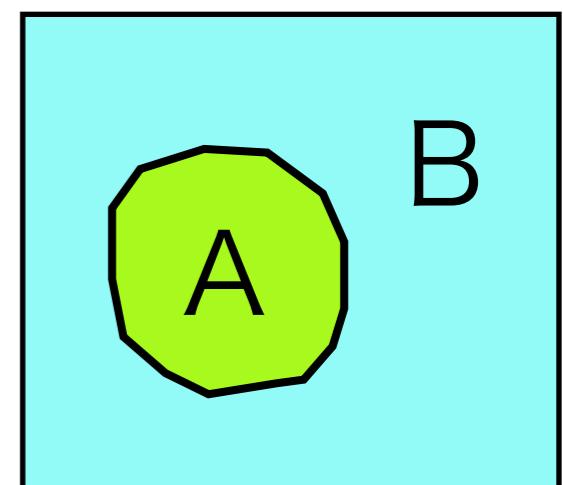
reduced density matrix

$$\rho_A = -\text{Tr}_{\mathcal{H}_B} \rho_{tot}$$

entanglement entropy

$$S_A = -\text{Tr}_A \rho_A \log \rho_A$$

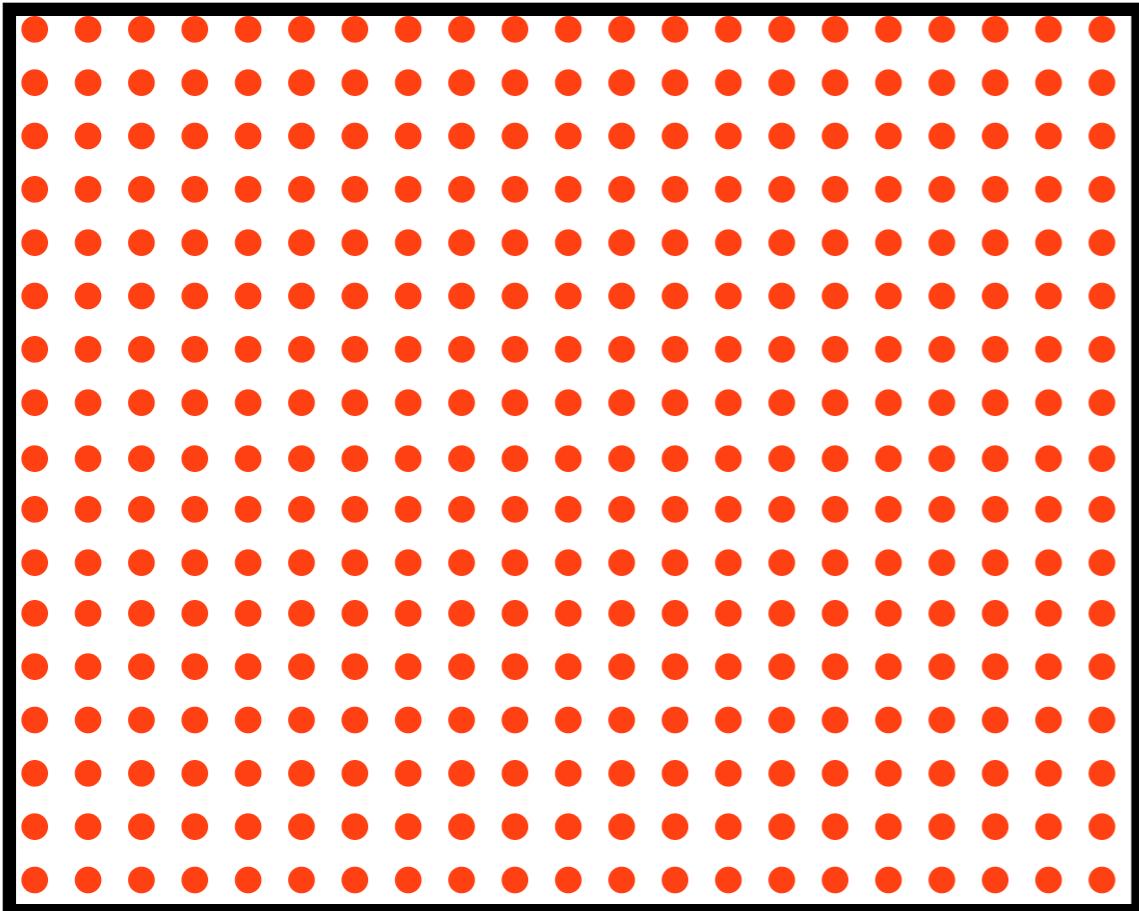
At finite T, it is equivalent to the thermal entropy.



# Schematic picture of E.E.

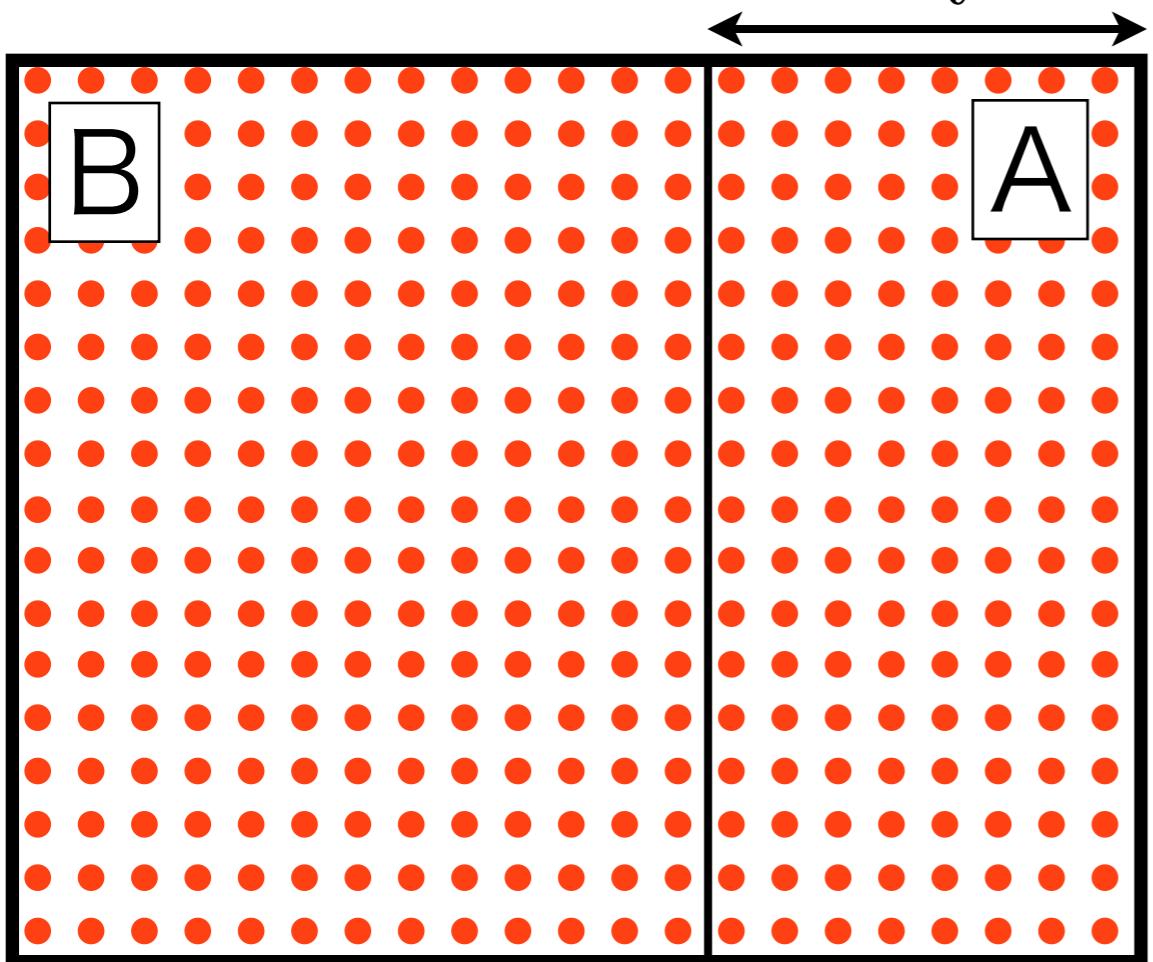
(1+1)-dim. model

Holzhey,Larsen and Wilczek: NPB424 (1994) 443  
Calabrese and Cardy: J.S.M.0406(2004)P06002  
Calabrese and Cardy, arXiv:0905.4013



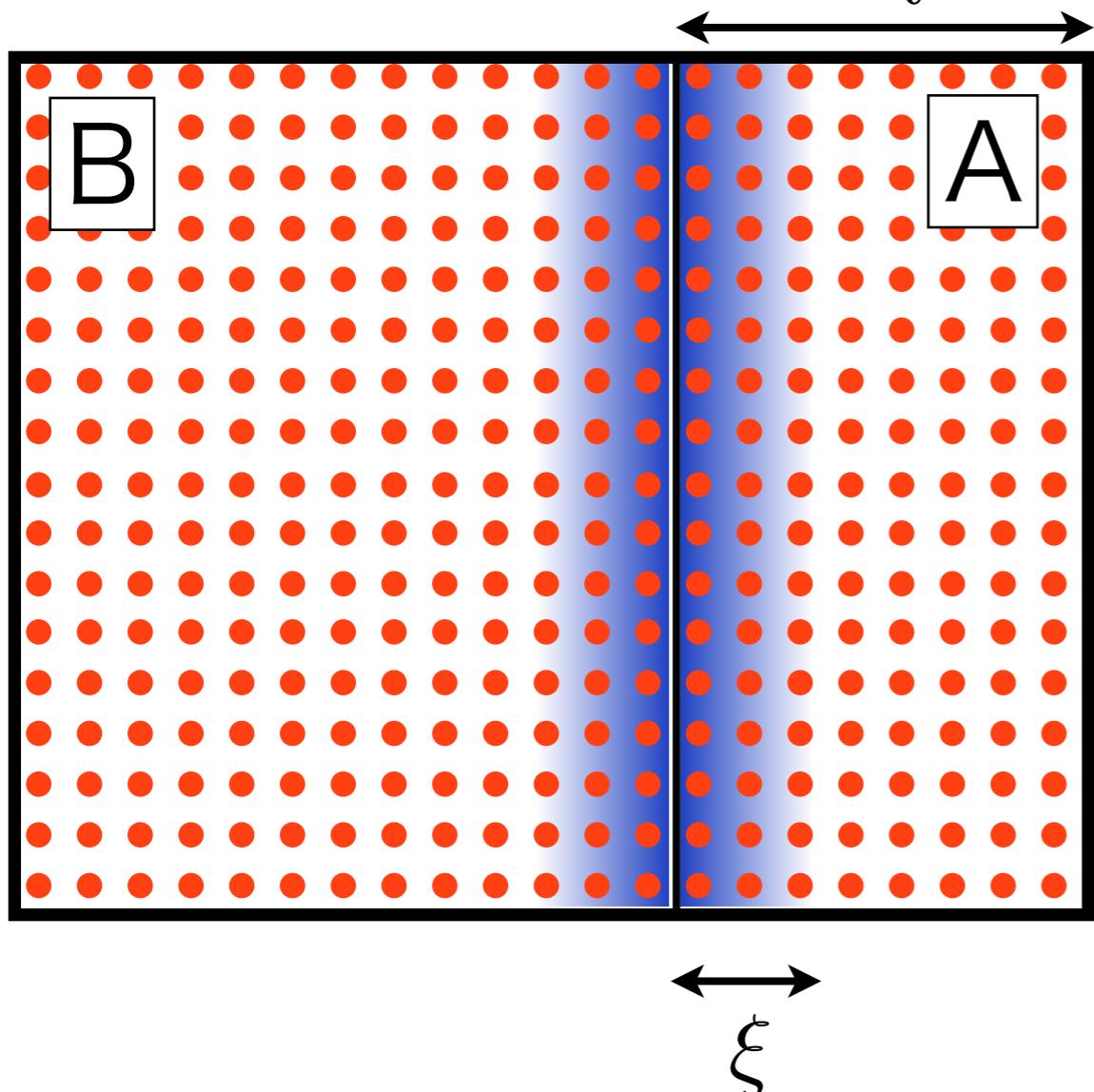
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(1+1)-dim. model



# Schematic picture of E.E.

(1+1)-dim. model



At the critical point,

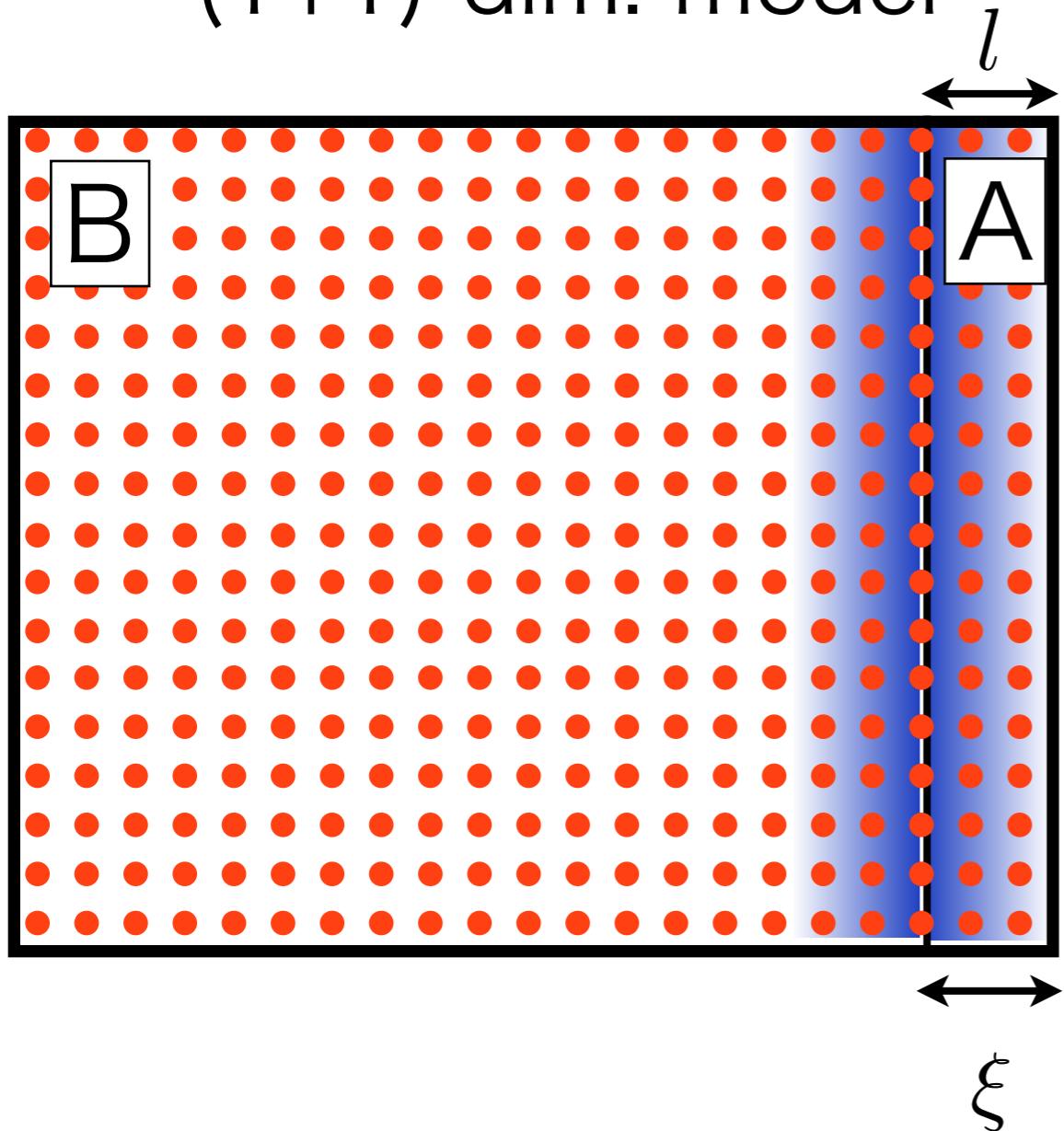
$$S_A(l) = \frac{c}{3} \log \frac{l}{a} + c_1$$

c is the central charge  
in 2d CFT.

$\xi$ : correlation length of the system

# Schematic picture of E.E.

(1+1)-dim. model



In the non critical system,

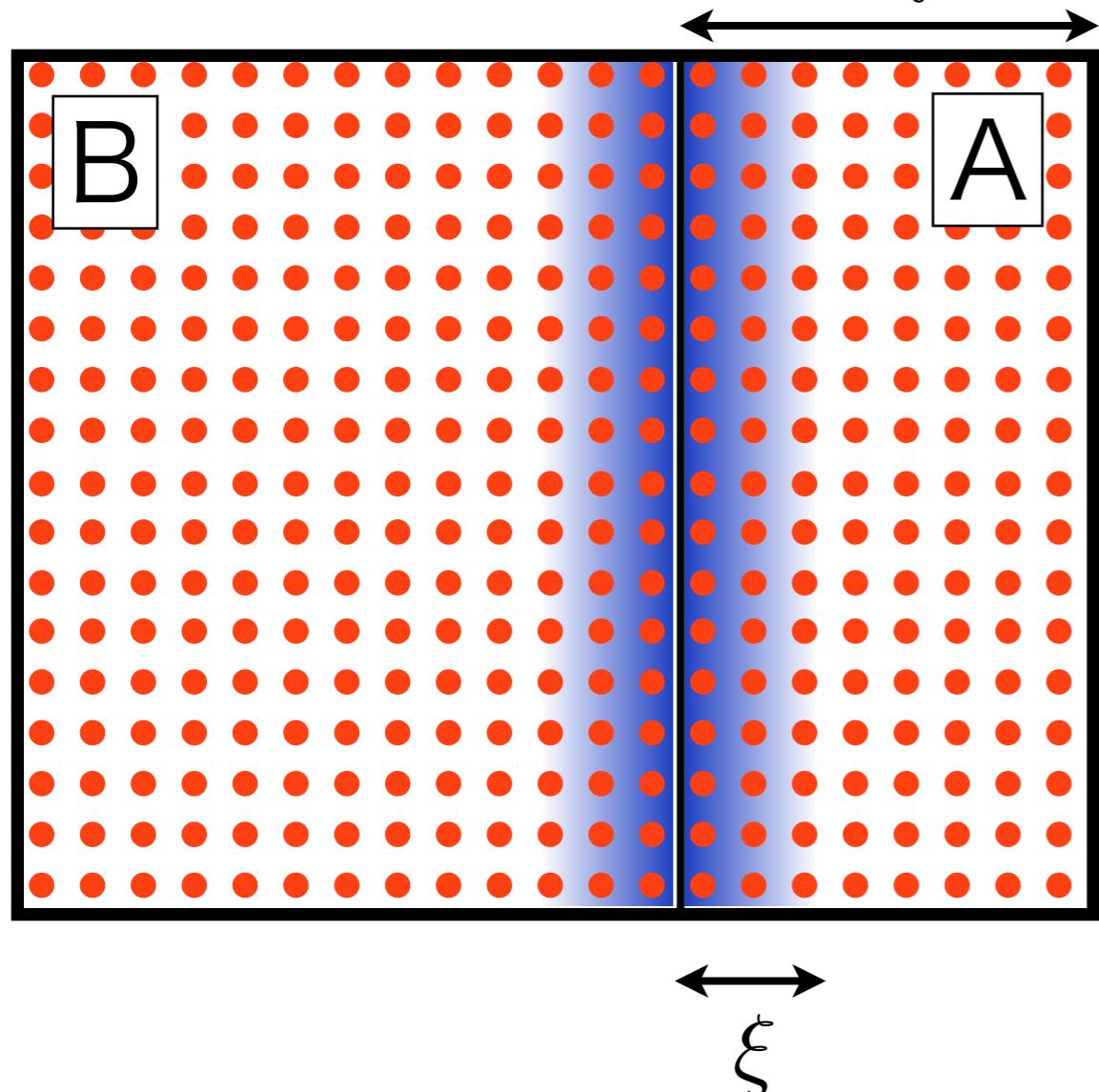
$$l \ll \xi$$

$$S_A(l) \sim \frac{c}{3} \log \frac{l}{a}$$

$\xi$ : correlation length of the system

# Schematic picture of E.E.

(1+1)-dim. model



In the non critical case,

$$l \ll \xi$$

$$S_A(l) \sim \frac{c}{3} \log \frac{l}{a}$$

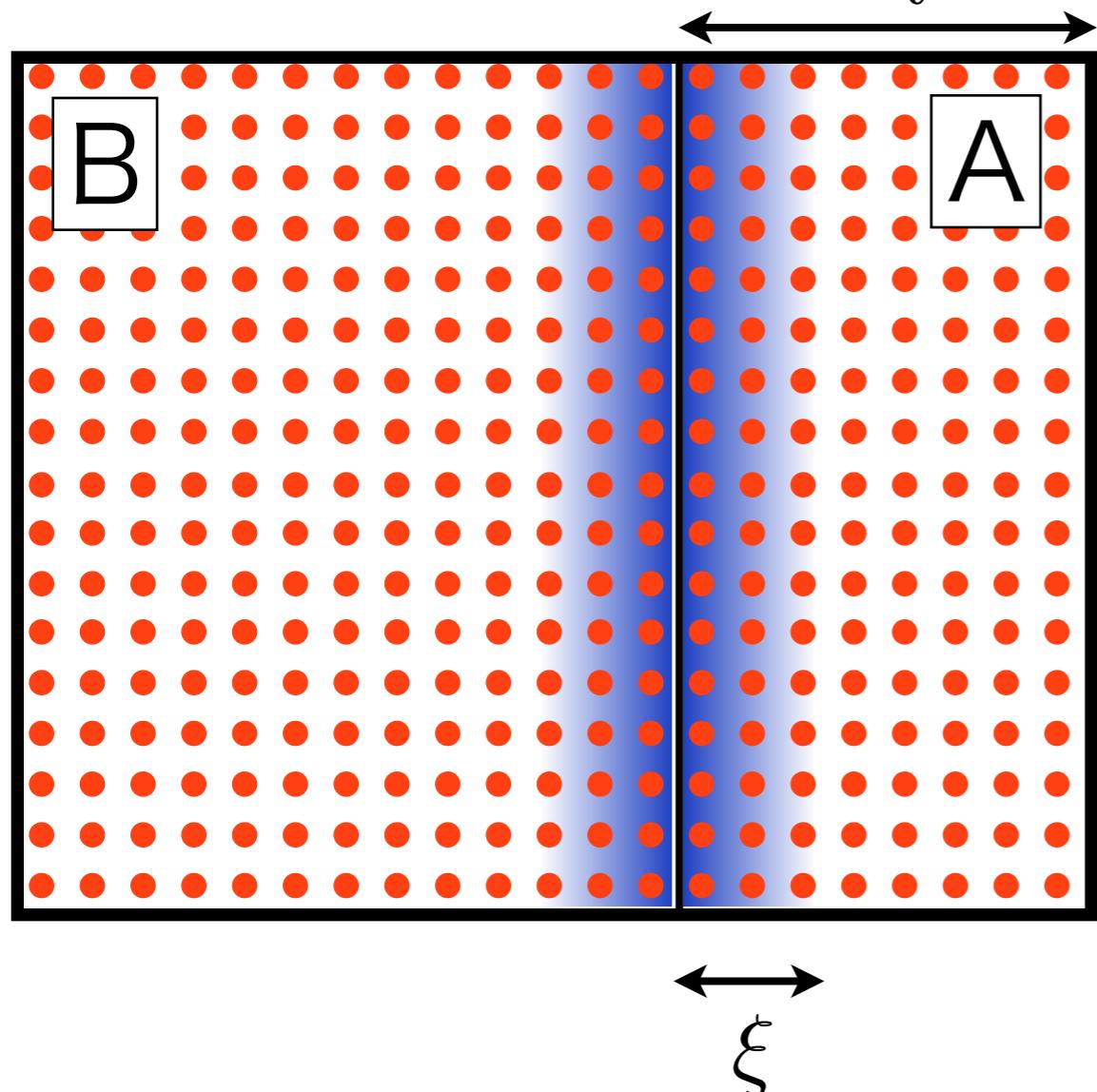
$$l \gg \xi$$

$$S_A(l) \rightarrow \frac{c}{3} \log \frac{\xi}{a}$$

$\xi$ : correlation length of the system

# Schematic picture of E.E.

(1+1)-dim. model



At the critical point or  $l \ll \xi$  in the noncritical system

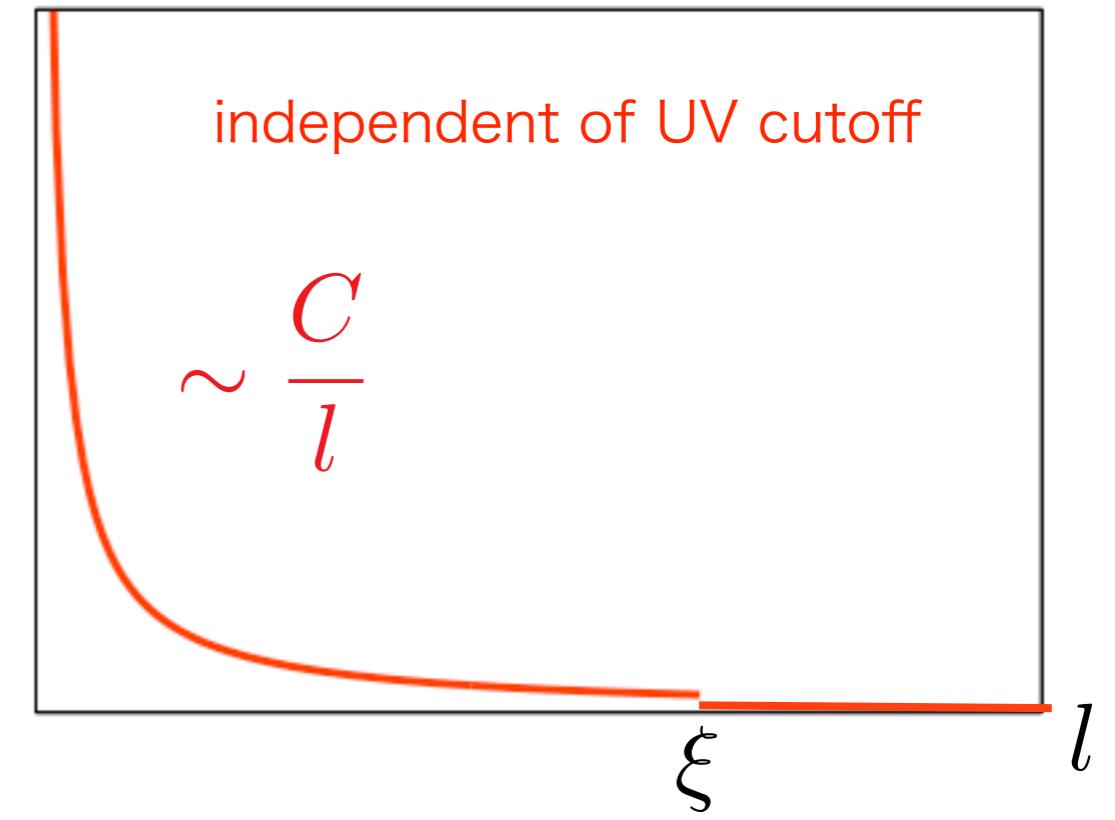
$$S_A(l) = \frac{c}{3} \log \frac{l}{a} + c_1$$

In the non-critical system,

$$S_A(l) \underset{l \gg \xi}{\rightarrow} \frac{c}{3} \log \frac{\xi}{a}$$

$\xi$ : correlation length of the system

$$\frac{\partial S_A(l)}{\partial l}$$



# Difficulties to obtain E.E. in 4d gauge theory

- UV cutoff dependence of 4d E.E.

2d

$$S_A(l) = c_0 \log(l/a)$$

Entropic C-function  $C(l) = l \frac{dS}{dl}$

4d

$$S_A(l) = c_0 \frac{\text{Area}}{a^2} - c'_0 \frac{\text{Area}}{l^2} + c_1 \log(l/a) + (\text{regular terms})$$

Entropic C-function

$$C(l) = l^3 \frac{1}{|\partial A|} \frac{dS}{dl}$$

Ryu and Takayanagi:PRL96(2006)181602  
JHEP 0608(2006)045

Nishioka and Takayanagi JHEP 0701(2007) 090

cf) Holographic approach

C is obtained by AdS and QFT

- (local) gauge invariance

# E.E. for gauge theory

- P.V.Buividovich and M.I.Polikarpov PLB670(2008)141  
extended Hilbert space
- H.Casini, M.Muerta and J.A.Rosabal arXiv:1312.1183  
electric b.c.(electric center), magnetic center, trivial center
- D.Radicevic arXiv:1404.1391  
magnetic center
- W.Donnelly PRD85 (2012) 085004  
extended lattice construction
- S.Ghosh, R.M.Soni,S.P.Trivedi arXiv:1501.02593
- S.Aoki, T.Iritani, M.Nozaki et.al. arXiv:1502.04267  
maximally gauge invariant reduced density matrix

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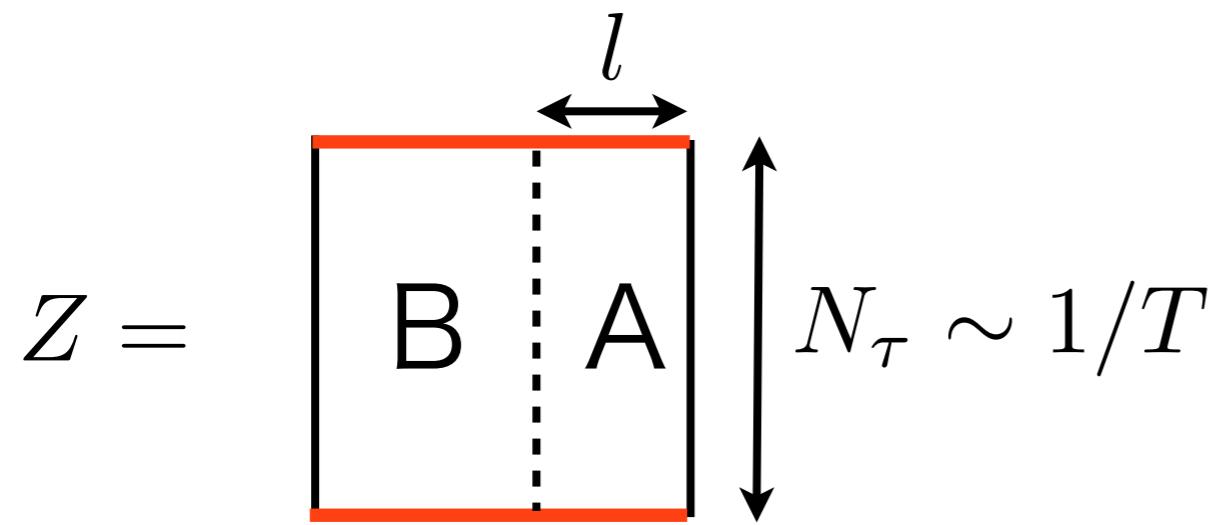
Our work

red definitions are inadequate  
for E.E. or rhoA

# Replica method

Calabrese and Cardy: J.S.M.0406(2004)P06002

# Replica method



entanglement entropy:

$$S_A = - \lim_{n \rightarrow 1} \frac{\partial}{\partial n} \ln \text{Tr}_A \rho_A^n$$

$$S_A(l) = - \lim_{n \rightarrow 1} \frac{\partial}{\partial n} \ln \left( \frac{Z(l, n)}{Z^n} \right)$$

# Replica method

$$Z = \text{[Diagram of a system A+B with a horizontal dashed line at height } l\text{, divided into regions A and B. The width of the system is labeled } l\text{. The width of region A is also labeled } l\text{. The total width of the system is labeled } N_\tau \sim 1/T\text{.]}$$

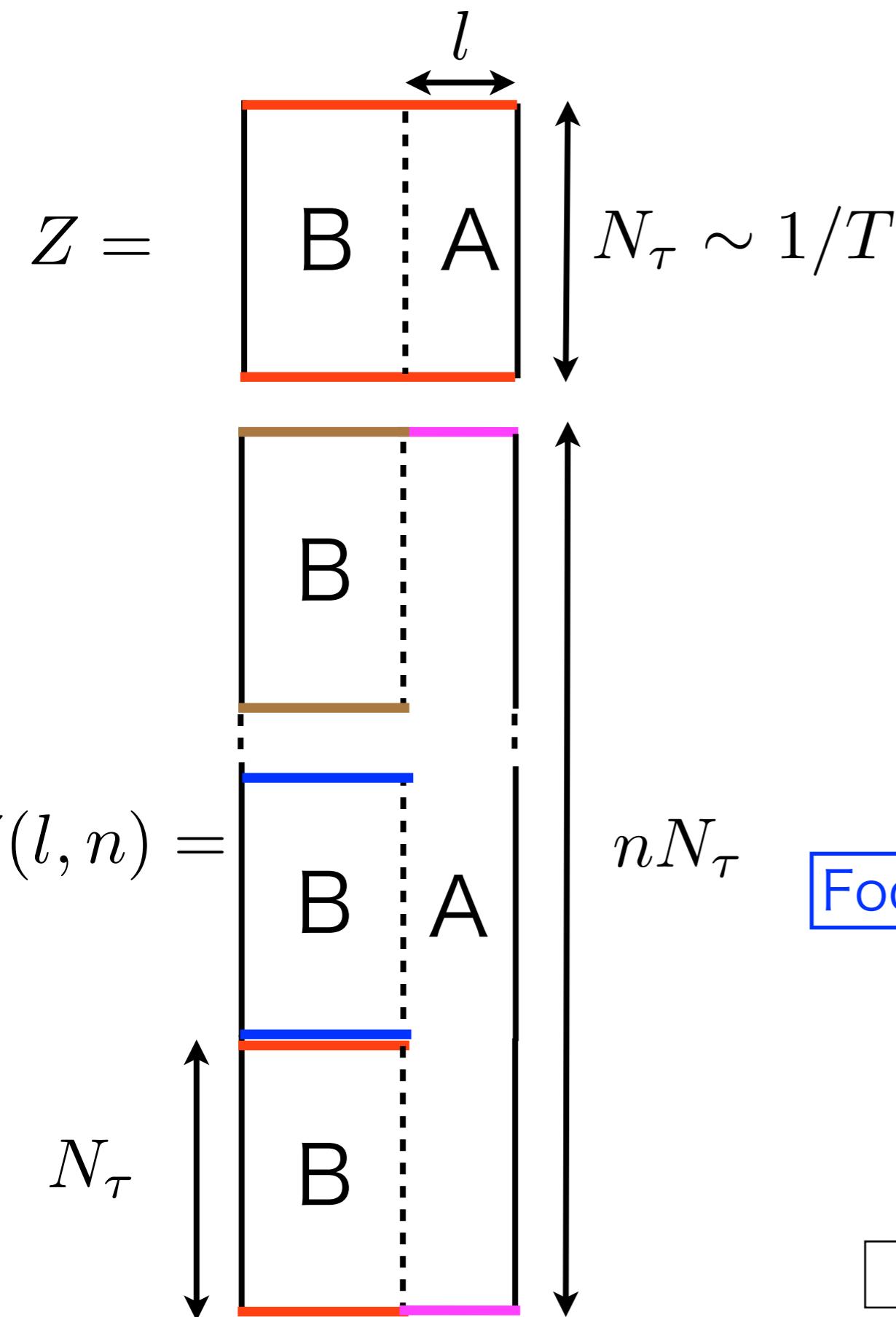
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$$Z(l, n) = \text{[Diagram of a system A+B with a horizontal dashed line at height } l\text{, divided into regions A and B. The width of the system is labeled } nN_\tau\text{. The width of region A is labeled } l\text{. The total width of the system is labeled } N_\tau\text{. The diagram shows three horizontal bars: a blue bar at the top, an orange bar in the middle, and a red bar at the bottom. The blue bar is at height } l\text{, the orange bar is at height } l/2\text{, and the red bar is at height } 0\text{. The regions A and B are indicated by vertical dashed lines at the boundaries of the bars. The width of region A is labeled } l\text{. The width of the system is labeled } nN_\tau\text{.]}}$$

# Replica method



entanglement entropy:

$$S_A = - \lim_{n \rightarrow 1} \frac{\partial}{\partial n} \ln \text{Tr}_A \rho_A^n$$

$$S_A(l) = - \lim_{n \rightarrow 1} \frac{\partial}{\partial n} \ln \left( \frac{Z(l, n)}{Z^n} \right)$$

observable:

$$\frac{\partial S_A(l)}{\partial l} = \lim_{n \rightarrow 1} \frac{\partial}{\partial l} \frac{\partial}{\partial n} F[l, n] \quad : \text{free energy}$$

$$\begin{aligned}
 & \rightarrow \frac{F[l+a, n=2] - F[l, n=2]}{a} \\
 & = \int_0^1 d\alpha \langle S_{l+a}[U] - S_l[U] \rangle_\alpha
 \end{aligned}$$

using the interpolation action

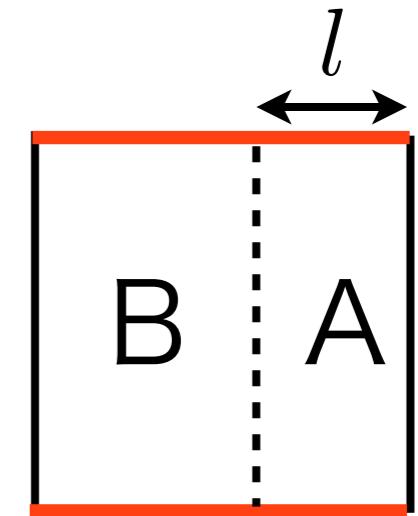
$$S_{int} = (1 - \alpha) S_l[U] + \alpha S_{l+a}[U]$$

we measure the diff. of the action density

# Simulation results

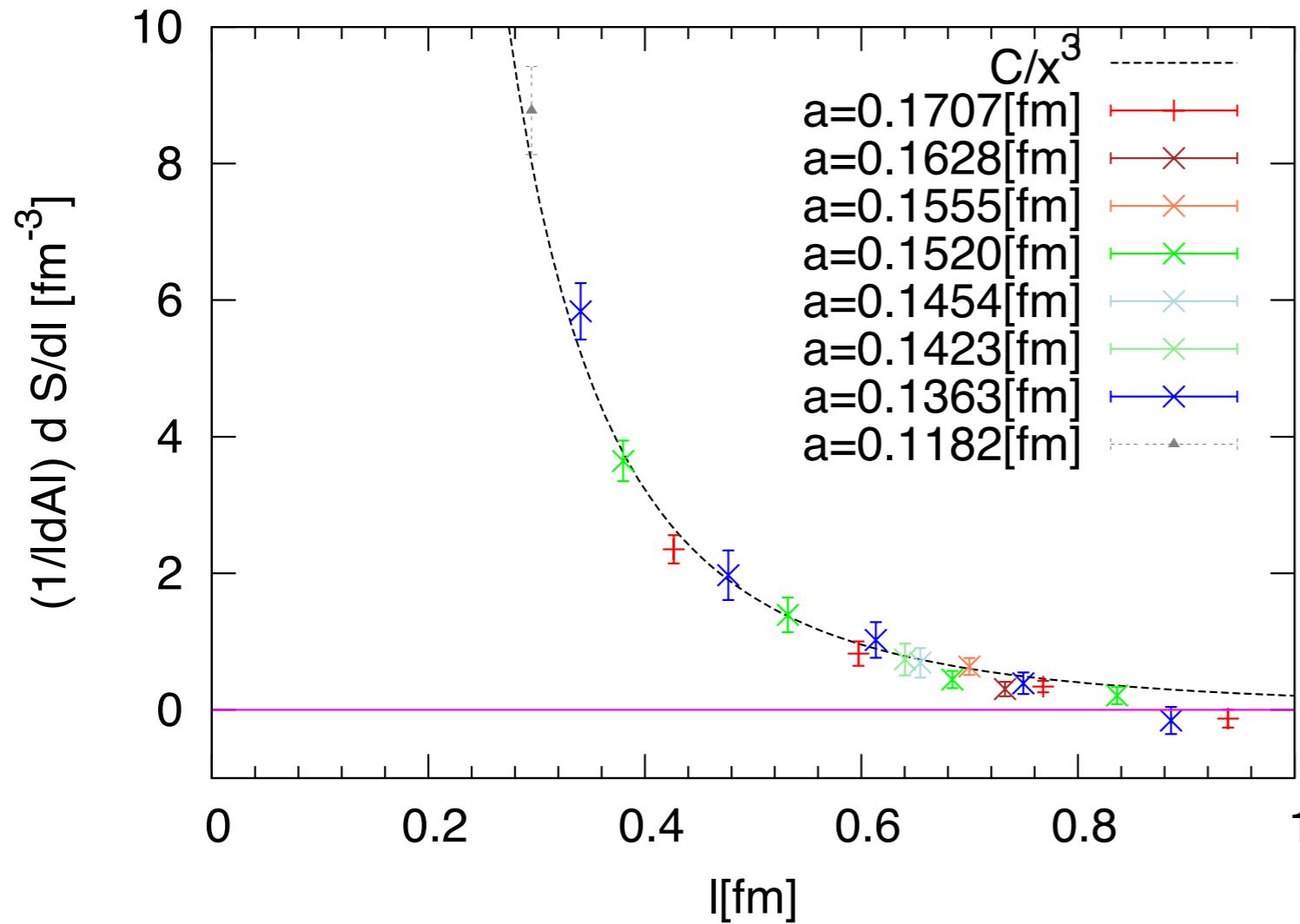
# Simulation setup

- Wilson plaquette gauge action
- $N_s = N_t = 16, 32$
- $l/a = 2, 3, 4, 5, (6)$
- $\beta = 5.70 - 5.87$
- # of configuration 12,000~84,000
- scale setting  $r_0 = 0.5$  fm and ALPHA coll.



# Lattice results for quenched SU(3)

T=0, quenched QCD



We measure  $\sim 84,000$  configs.

- in short range,  $1/l^3$  scaling
- the coefficient is  $C=0.2064(73)$

cf.)  $C \sim 0.09$  in SU(2)

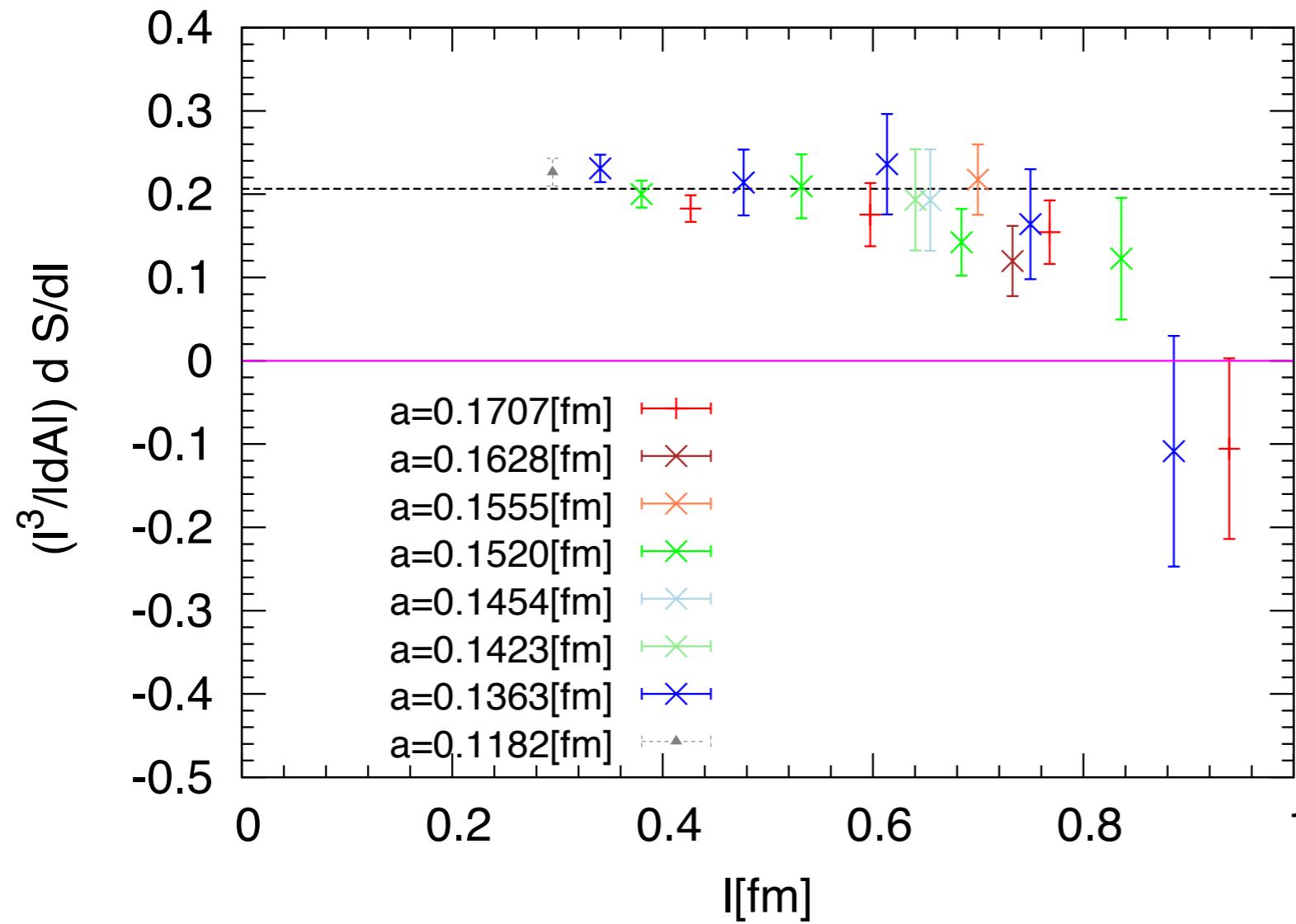
The  $N_c$  dependence is

$$\frac{0.2064}{0.09} = 2.29 \sim \frac{3^2}{2^2}$$

# Entropic C-function

independent of UV cutoff

$$C(l) = l^3 \frac{1}{|\partial A|} \frac{dS}{dl}$$



In short  $l$  region,  $C$  is constant.

The discontinuity is not clear.

cf.)  $\frac{1}{\Lambda_{QCD}} \sim 0.7[\text{fm}]$

# Comparison with Ryu-Takayanagi results

Ryu and Takayanagi:PRL96(2006)181602  
JHEP 0608(2006)045

Holographic (or field theoretical) approach

$$(3+1)\text{-dim. CFT} \quad \frac{1}{|\partial A|} S_A(l) = c \frac{N_c^2}{a^2} - c' \frac{N_c^2}{l^2}$$

$c'$  is obtained by AdS and QFT

$$c' \sim 0.0049 \quad \text{for free real scalar theory}$$

Estimation for non-abelian gauge theory  $A_\mu^a \quad \{ \begin{array}{l} a = 1, \dots, 8 \\ \mu = 1, \dots, 4 \end{array}$

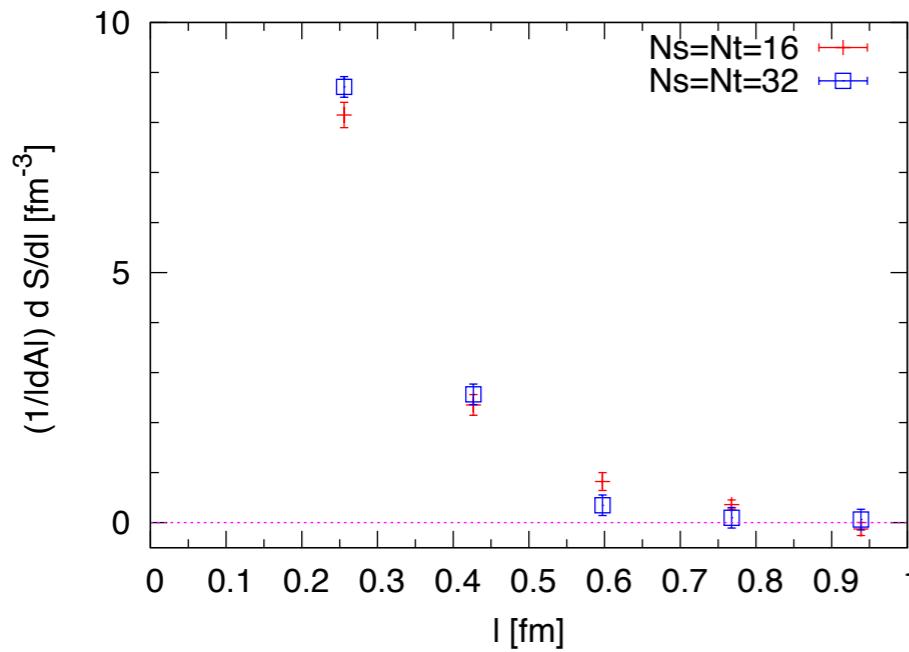
$$C_{\text{gauge}} \sim 2c' \cdot 2 \cdot 8 \sim 0.1568$$

Our numerical result

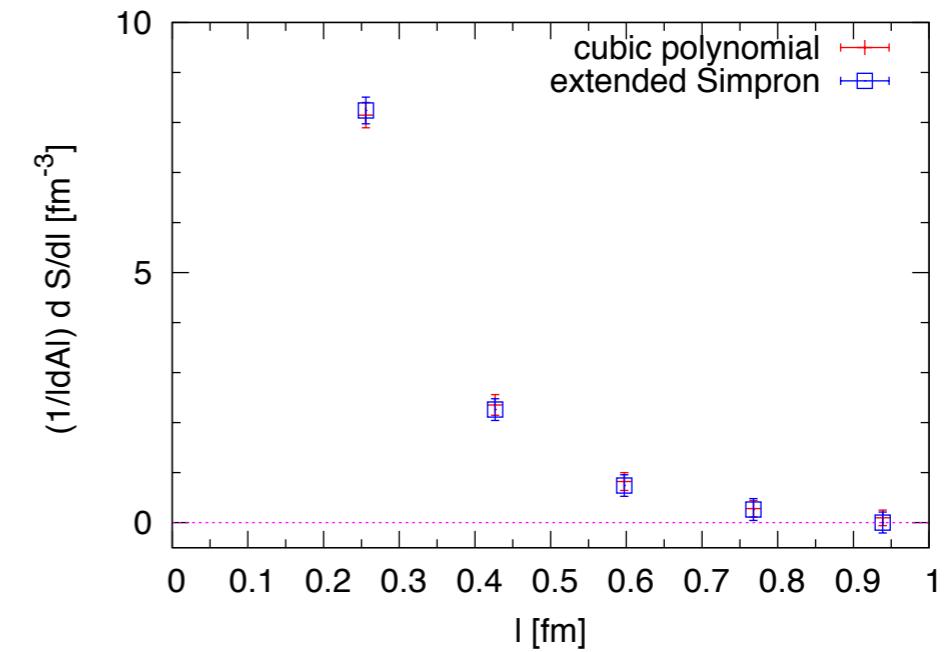
$$C_{\text{gauge}} \sim 0.2064$$

# Detail analyses

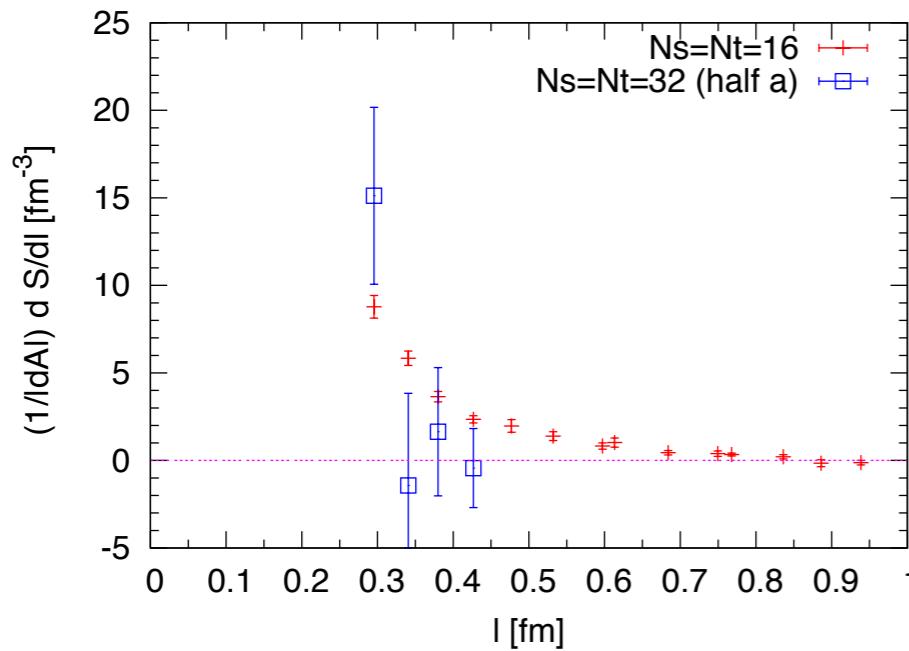
finite vol. effect



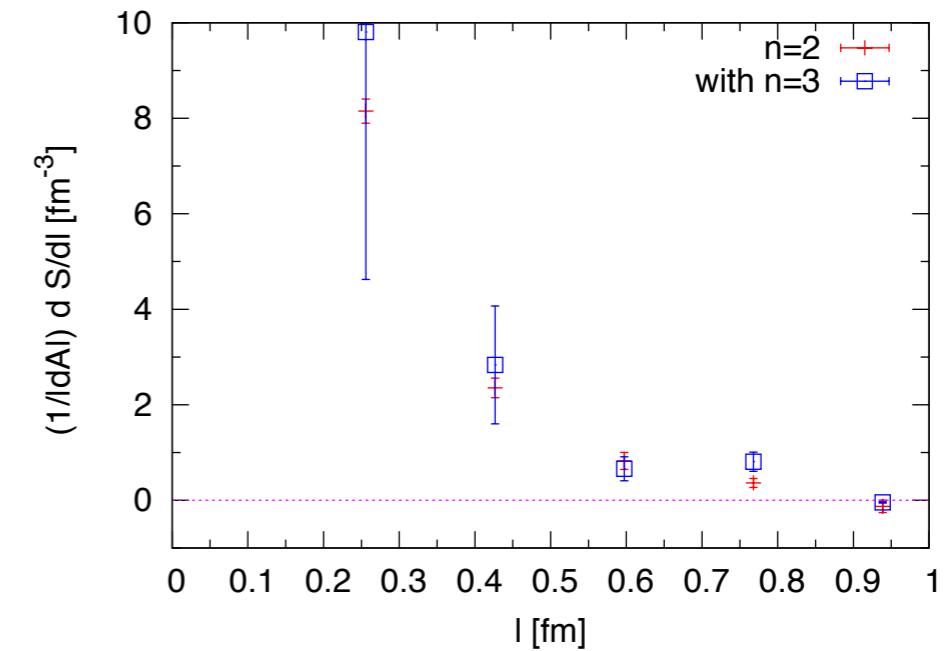
algorithm dependence of  
the numerical integration



UV cutoff dependence



replica number dependence

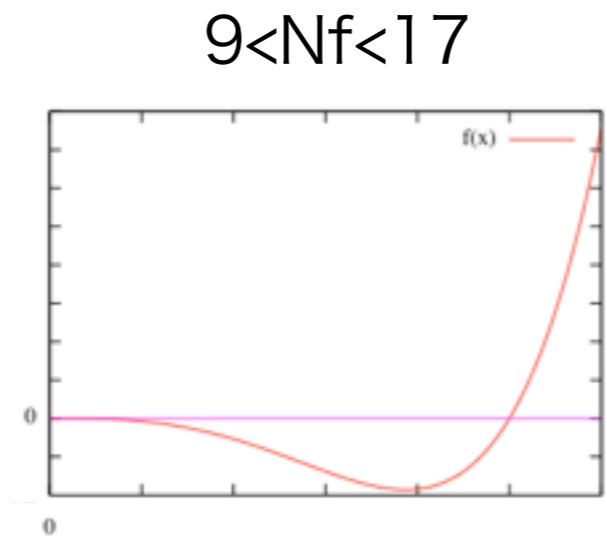
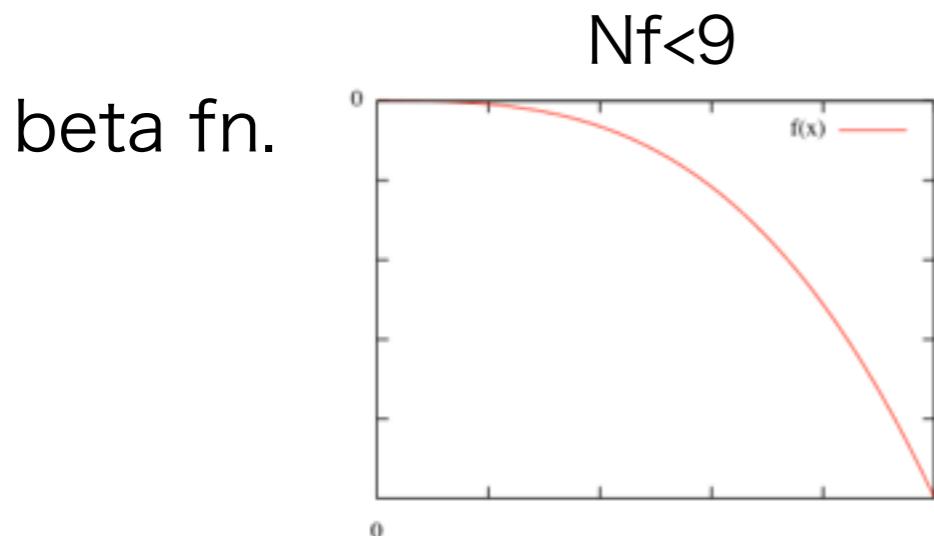


# Summary

- This is the first precise determination of E.E. for quenched QCD
- $N_c$  dependence in the short  $|t|$  region is  $N_c^2$  as expected by AdS/CFT and field theoretical insights
- No discontinuity exists as contrast with SU(2) results
- Entropic C-function shows UV cutoff independence
- Value of C-function agrees with Ryu-Takayanagi work
- replica number ( $n \rightarrow 1$ ) dependence

# Future directions for E.E. using the lattice

- QCD at zero  $T$ 
  - give a novel observation for confinement
  - even in full QCD case
- QCD at finite  $T$ 
  - gives the thermal entropy and the correlation length in QGP phase
- conformal window in 4dim  $N_f$  flavor QCD
  - would give the  $\alpha$ -function and central charge



nontrivial IR fixed point is found by lattice simulation  
cf) E.I. PTEP(2013)083B01