

Keldysh形式による量子輸送における揺らぎの定理

内海裕洋

三重大工学部



斎藤圭司

東京大学理学部

Dmitry Golubev

Karlsruhe Inst. of Technology

Michael Marthaler

University Karlsruhe

Gerd Schön

University Karlsruhe

藤澤利正

東京工業大学、NTT物性基礎研究所

小林研介

京都大学化学研究所

R. Leturcq

IEMN-CNRS

...

...

31 Aug 2010

Fluctuation theorem (FT)

Evans, Cohen, Morris, PRL 1993, Gallavotti, Cohen, PRL 1995

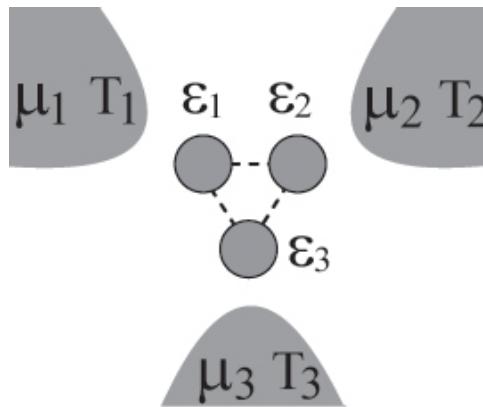
$$\frac{P(-\Delta S)}{P(+\Delta S)} = \exp(-\Delta S) \quad \Delta S \text{ entropy production during } \tau$$

An exact relation in *far from equilibrium* --- **microscopic reversibility**

- Second law of thermodynamics [Jarzynski PRL 1997, Crooks PRE 1999, etc.]
- Landauer principle [Piechocinska, PRA 2000]
- Fluctuation-dissipation theorem & Onsager relations [Gallavotti PRL 1996]

Full counting statistics (FCS)

[Tobiska Nazarov PRB 2005, Saito YU, PRB 2008, Förster Büttiker PRL 2008, Andrieux Gaspardt Monnai Tasaki NJP 2009]

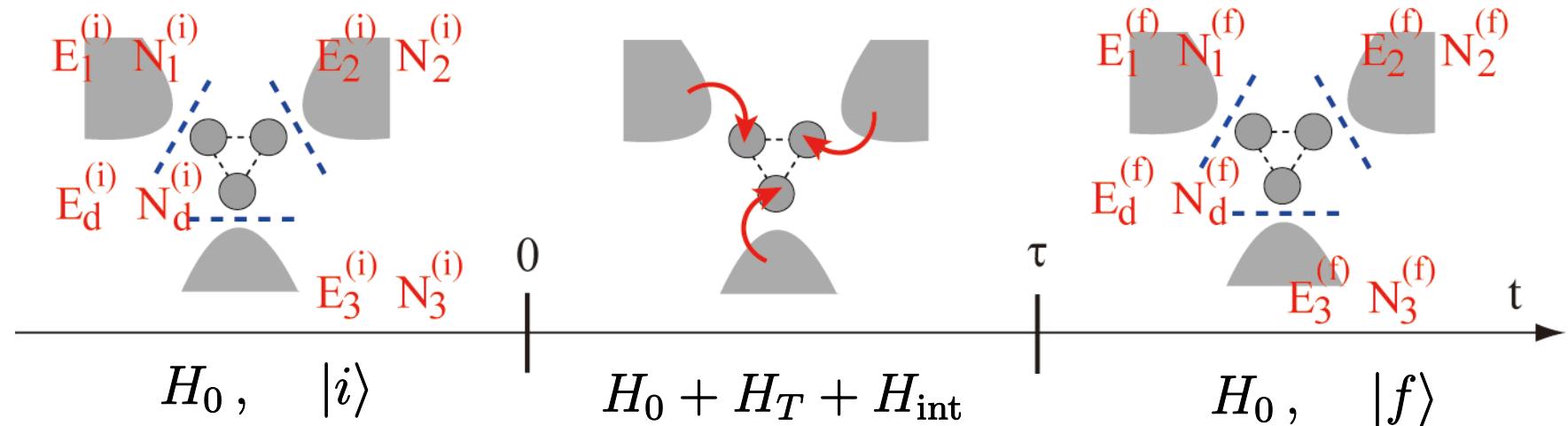


$$H_d = \sum_{i,j,\sigma} |t_{ij}| e^{i\phi_{ij}(B)} d_{i\sigma}^\dagger d_{j\sigma}$$

$$H_r = \sum_{k\sigma} \varepsilon_{rk} a_{rk\sigma}^\dagger a_{rk\sigma} \quad (r = 1, \dots, m)$$

$$H_{\text{int}} = \sum_{ij\sigma\sigma'} U_{i\sigma j\sigma'} d_{i\sigma}^\dagger d_{j\sigma'}^\dagger d_{j\sigma'} d_{i\sigma} / 2$$

$$H_T = \sum_{rki\sigma} |t_{rki}| e^{i\phi_{ri}(B)} d_{i\sigma}^\dagger a_{rk\sigma} + \text{H.c.} \quad \phi(-B) = -\phi(B)$$



$$H_0 = \sum_{s=1, \dots, m, d} H_s$$

At initial and final states, reservoirs and dots are decoupled
 \Rightarrow Energy & number of electrons are good quantum numbers

Flowing out charge $q_{cs} = N_s^{(f)} - N_s^{(i)}$ Flowing out energy $q_{hs} = E_s^{(f)} - E_s^{(i)}$

Joint probability distribution

$$P(\{q_{cs}, q_{hs}\}; B) = \sum_{if} |\langle f | e^{-iH\tau} | i \rangle|^2 \langle i | \rho_0 | i \rangle \prod_{s'} \delta_{q_{cs'}, N_{s'}^{(f)} - N_{s'}^{(i)}} \delta(q_{hs'} - E_{s'}^{(f)} + E_{s'}^{(i)})$$

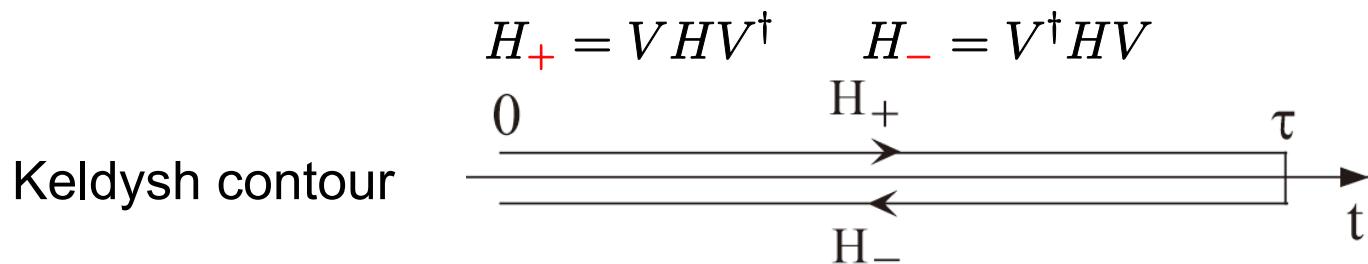
FCS & FT

Fourier Transform \Rightarrow Characteristic Function

$$\mathcal{Z}(\{\lambda_{cs}, \lambda_{hs}\}; B) = \text{Tr} [\rho_0 V^\dagger e^{iH\tau} V^2 e^{-iH\tau} V^\dagger] = \text{Tr} [e^{iH_- \tau} e^{-iH_+ \tau} \rho_0]$$

$$V(\{\lambda_{cs}, \lambda_{hs}\}) = \prod_s \exp[-i(\lambda_{hs} H_s + \lambda_{cs} N_s)/2]$$

↑ ↑
Time translation + Gauge transformation



Microscopic reversibility + $\rho_0 \propto \prod_{s=1, \dots, m, d} \exp[-(H_s - \mu_s N_s)/T_s]$

$$\mathcal{F}(\{\chi_{\alpha s} \equiv \lambda_{\alpha s} - \lambda_{\alpha m}\}; B) \equiv \lim_{\tau \rightarrow \infty} \ln \mathcal{Z}(\{\lambda_{\alpha s}\}; B)/\tau$$

$$\mathcal{F}(\{\chi_{\alpha r}\}; B) = \mathcal{F}(\{-\chi_{\alpha r} + i\mathcal{A}_{\alpha r}\}; -B)$$

$$\mathcal{A}_{cr} = \mu_r/T_r - \mu_m/T_m \quad \mathcal{A}_{hr} = T_m^{-1} - T_r^{-1} \quad \text{Affinities}$$

Universal Relations among Nonlinear Transport Coefficients

Two terminal isothermal system

$$\mathcal{F}(\lambda; B) = \mathcal{F}(-\lambda + i\mathcal{A}; -B) \quad \mathcal{A} = (\mu_1 - \mu_2)/T$$

$$\frac{P(q; B)}{P(-q; -B)} \approx e^{q\mathcal{A}}$$

Nonlinear transport coefficient $\langle\langle I^n \rangle\rangle = \left. \frac{\partial^n \mathcal{F}(\lambda; B)}{\partial(i\lambda)^n} \right|_{\lambda=0} = \sum_{m=0}^{\infty} L_m^n(B) \frac{\mathcal{A}^m}{m!}$

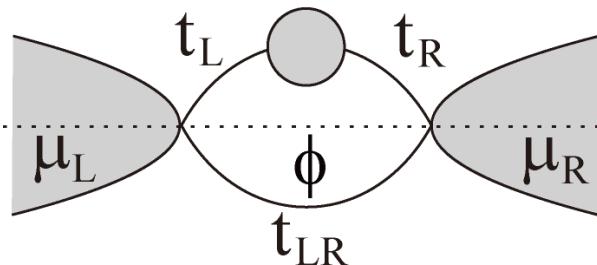
$$L_0^2 = 2L_1^1 \quad L_{1,-}^1 = 0 \quad \text{Fluctuation-dissipation theorem & Onsager relation}$$

$$L_{2,+}^1 = L_{1,+}^2 \quad L_{0,+}^3 = 0 \quad L_{2,-}^1 = L_{1,-}^2 / 3 = L_{0,-}^3 / 6$$

[Saito YU PRB 2008]

$$L_{\pm} = L(B) \pm L(-B)$$

QD Aharonov-Bohm interferometer



$$H = \sum_{r=L,R} H_r + H_D + H_T + H_{\text{ref}}$$

$$\left. \begin{array}{l} H_r = \sum_{k\sigma} \varepsilon_{rk\sigma} a_{rk\sigma}^\dagger a_{rk\sigma} \\ H_D = \sum_{\sigma=\uparrow,\downarrow} \epsilon_D d_\sigma^\dagger d_\sigma + U n_\uparrow n_\downarrow \\ H_T = \sum_{rk\sigma} t_r d_\sigma^\dagger a_{rk\sigma} + \text{H.c.} \\ H_{\text{ref}} = \sum_{kk'\sigma} t_{LR} e^{i\phi} a_{Rk\sigma}^\dagger a_{Lk'\sigma} + \text{H.c.} \end{array} \right\} \text{Anderson model}$$

$$\phi(-B) = -\phi(B)$$

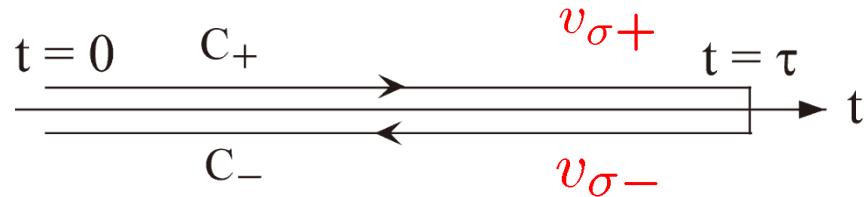
[YU Saito PRB 2009]

Keldysh path-integral

[Kamenev, in “Nanophysics: Coherence and Transport” (Elsevier 2005)]

$$\mathcal{Z} = \int \mathcal{D}[\mathbf{v}_\sigma, a_{rk\sigma}^*, d_\sigma^*, a_{rk\sigma}, d_\sigma] \exp \left[i \int_C dt \left(\mathcal{L}_0 + \frac{i}{U} \mathbf{v}_\sigma(t) v_\sigma(t) \right) \right]$$

Stratonovich-Hubbard transformation $U d_\uparrow^* d_\uparrow d_\downarrow^* d_\downarrow \rightarrow \sum_\sigma v_\sigma d_\sigma^* d_\sigma - v_\uparrow v_\downarrow / U$



Saddle-point approximation $v_{\sigma,\pm}(t) = v_\pm$

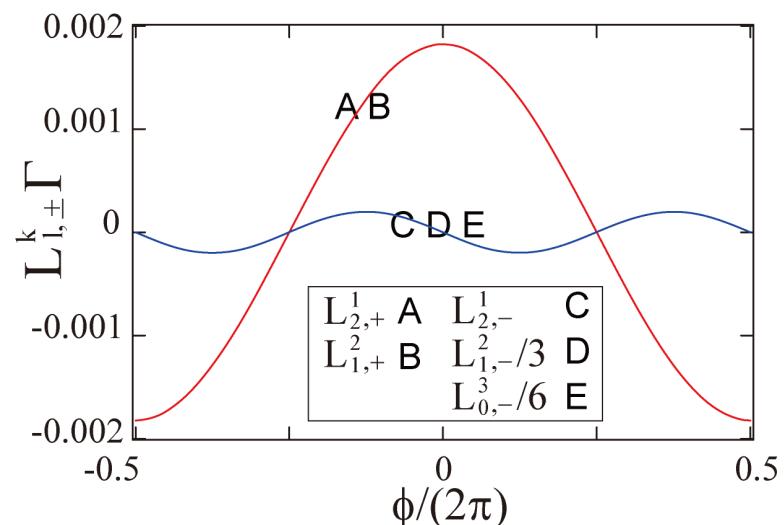
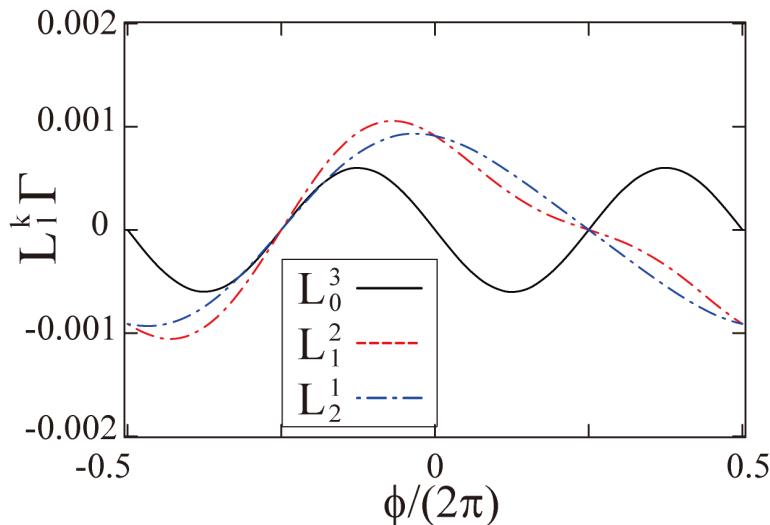
$$\frac{v_+ + v_-}{2} = v_c \quad \underline{v_+ - v_- = i v_q}$$

$$\mathcal{F} = \mathcal{F}_0 - 2v_c v_q / U \quad v_q = \frac{U}{2} \frac{\partial \mathcal{F}_0}{\partial v_c} \quad v_c = \frac{U}{2} \frac{\partial \mathcal{F}_0}{\partial v_q}$$

Extension of the Onsager relation

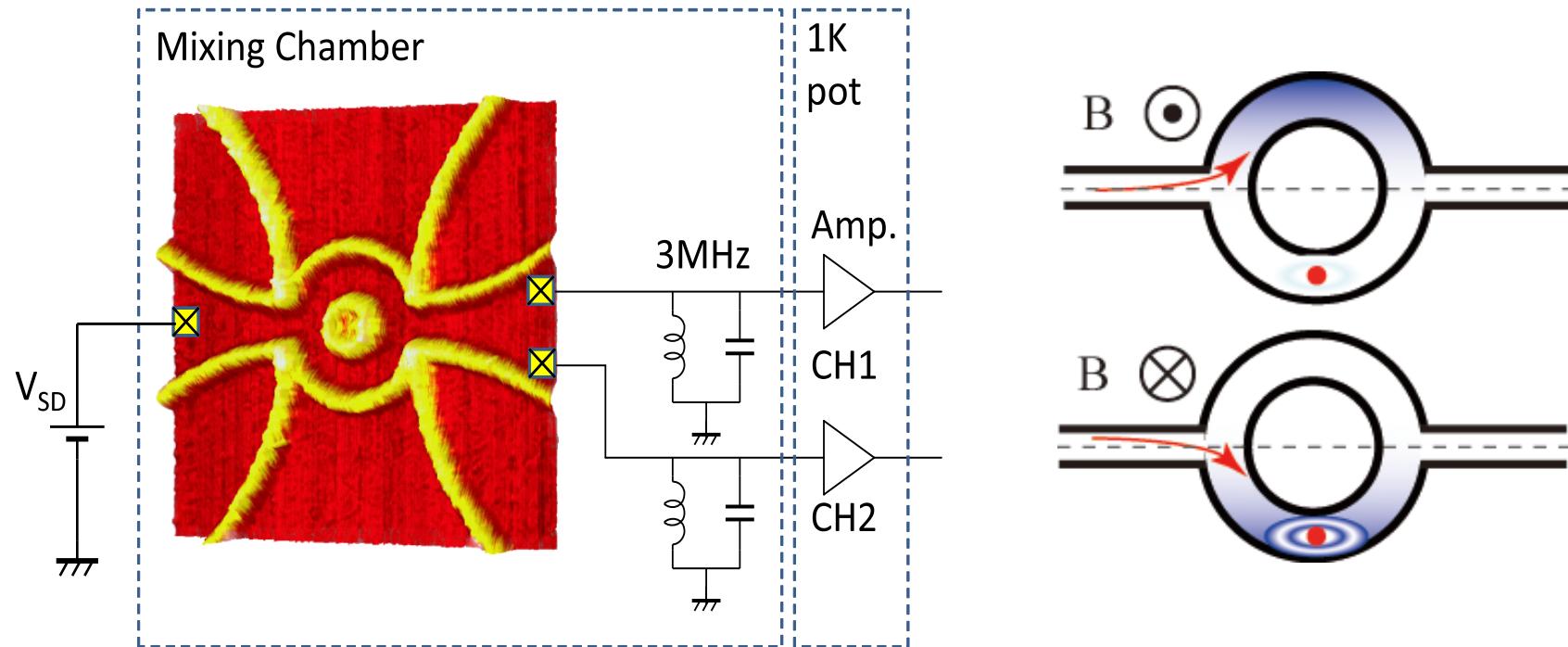
$$\left\{ \begin{array}{l} L_0^3 = \langle\langle I^3 \rangle\rangle|_{\mathcal{A}=0} = 6 \tilde{U}^{eq.} S_{IN}^{eq.} \chi_{II,N}^{eq.} \\ L_1^2 = d\langle\langle I^2 \rangle\rangle/d\mathcal{A}|_{\mathcal{A}=0} = 2 \tilde{U}^{eq.} \chi_{NI}^{eq.} \chi_{II,N}^{eq.} + 4 \tilde{U}^{eq.} S_{IN}^{eq.} \chi_{I,IN}^{eq.} \\ L_2^1 = d^2\langle\langle I \rangle\rangle/d\mathcal{A}^2|_{\mathcal{A}=0} = 4 \tilde{U}^{eq.} \chi_{NI}^{eq.} \chi_{I,IN}^{eq.} \end{array} \right. \quad \begin{array}{l} \chi_{II,N} = \frac{\partial S_{II}}{\partial \epsilon_D} \\ \chi_{I,IN} = \frac{\partial \chi_{IN}}{\partial \mathcal{A}} \\ \chi_{IN} = \frac{1}{2} \frac{\partial \langle\langle I \rangle\rangle}{\partial \epsilon_D} \end{array}$$

Coulomb interaction induces non-zero nonlinear transport coefficients



Magnetic field asymmetry in nonlinear conductance

Nakamura, Kobayashi, Leturcq, YU, et al. PRL 2010



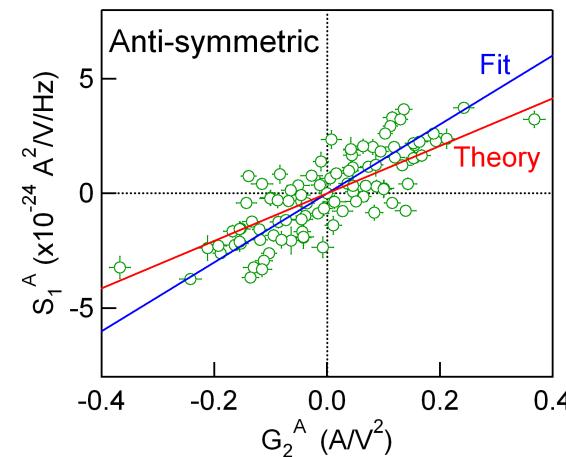
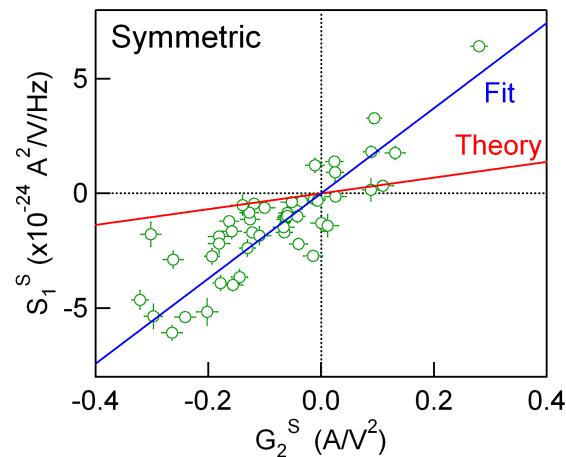
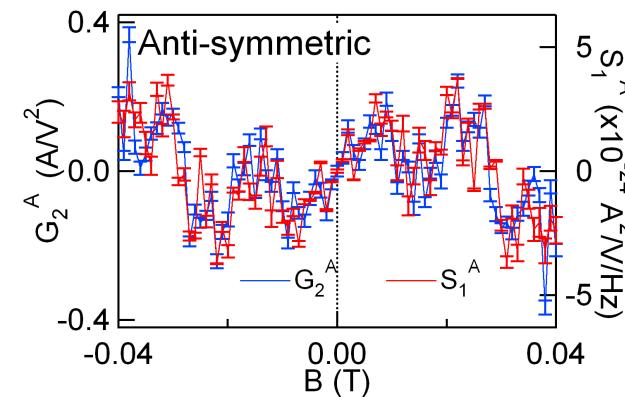
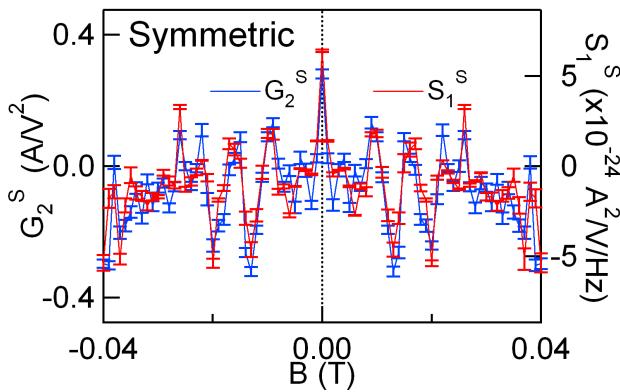
FT for lower order cumulants

$$G_2^S(B) \equiv G_2(B) + G_2(-B)$$

$$S_1^S(B) \equiv S_1(B) + S_1(-B)$$

$$G_2^A(B) \equiv G_2(B) - G_2(-B)$$

$$S_1^A(B) \equiv S_1(B) - S_1(-B)$$



$$S_1^S(B) = 2k_B T G_2^S(B) \alpha^S$$

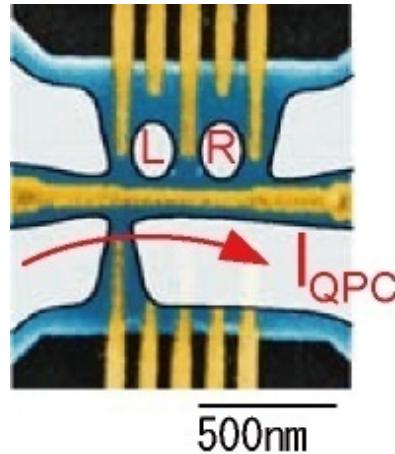
$$\alpha^S = 5.38^{+1.18}_{-0.70} \quad 1(\text{theory})$$

$$S_1^A(B) = 6k_B T G_2^A(B) \alpha^A$$

$$\alpha^A = 1.45^{+0.21}_{-0.12} \quad 1(\text{theory})$$

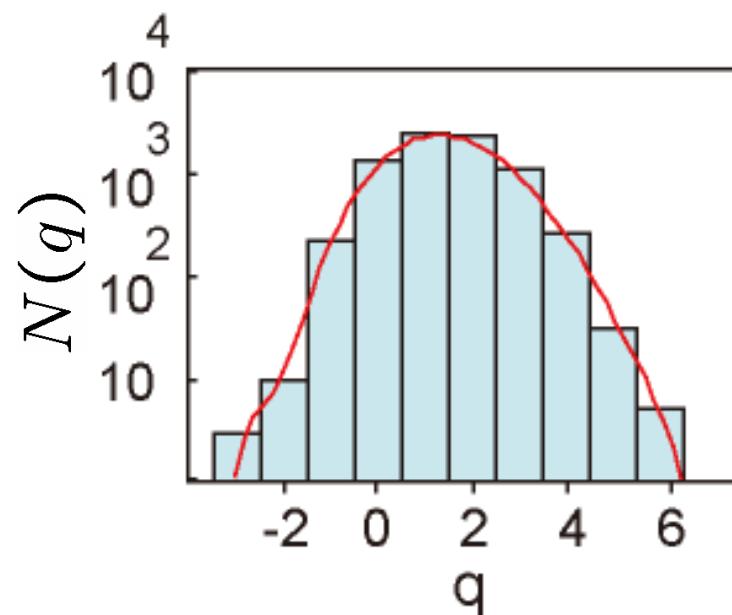
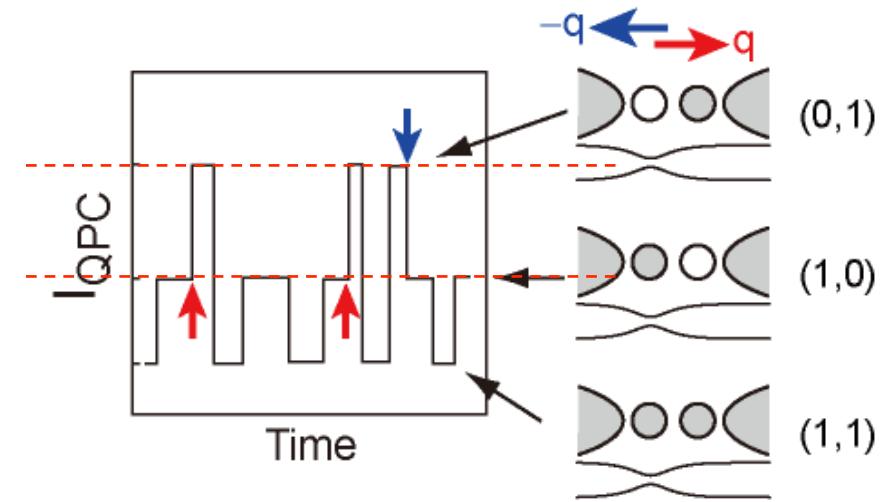
Bidirectional single electron counting

T. Fujisawa, T. Hayashi, R. Tomita, Y. Hirayama, Science 2006



Double quantum dots
(sample)

Quantum point-contact
electrometer



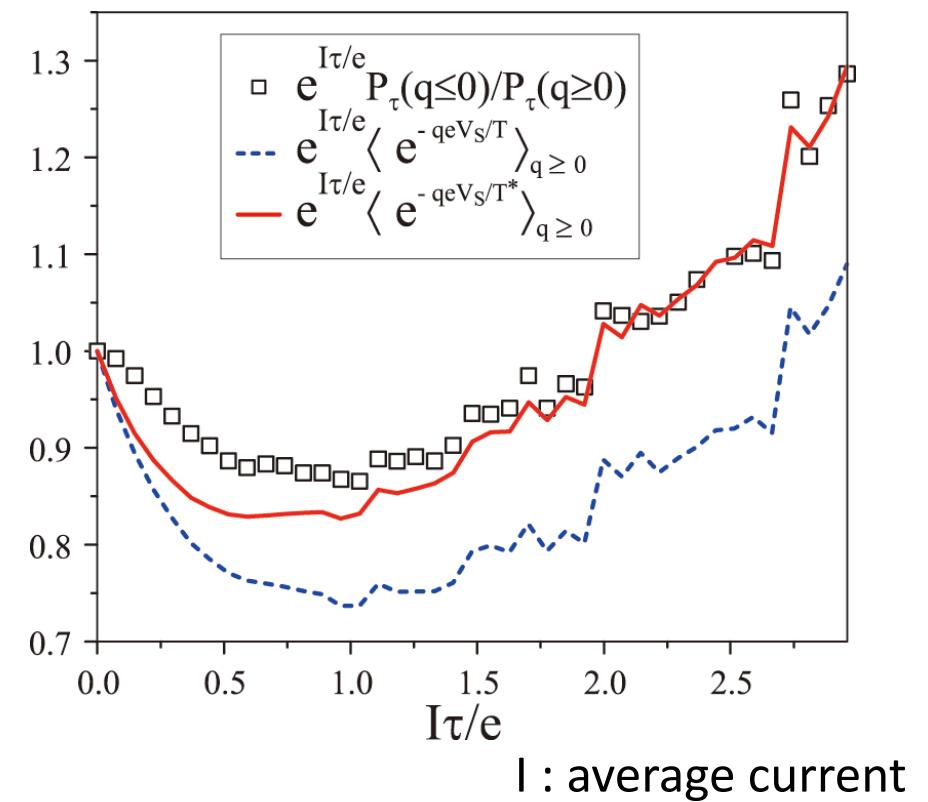
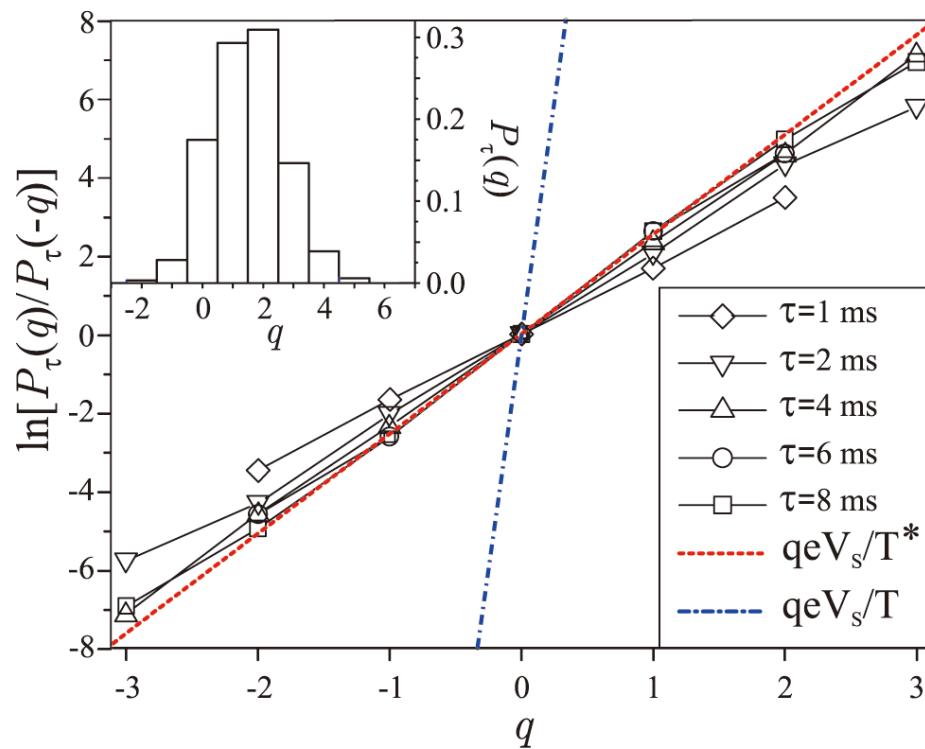
Short measurement time
& Low temperature
 $\tau = 4\text{ms}$

Experimental Test of FT

YU, D.Golubev, M.Marhtaler, K.Saito, T.Fujisawa, G.Schön, Phys. Rev. B **81**, 125331 (2010)

$$\frac{P_\tau(q)}{P_\tau(-q)} = \exp(qeV_s/T)$$

$$\frac{P_\tau(q \leq 0)}{P_\tau(q \geq 0)} = \langle \exp(-qeV_s/T) \rangle_{q \geq 0}$$



Fit to experiment with $T^* = 1.37\text{K} \gg T = 130\text{mK} \Leftrightarrow$ enhanced “temperature”

Summary

1. We combine the fluctuation theorem and the Keldysh formalism.
2. We experimentally confirmed the universal relations among nonlinear transport coefficients.

[1] K.Saito, YU, PRB **78**, 115429 (2008).

[2] YU, K. Saito, PRB 79, 23511 (2009).

[3] YU, D.Golubev, M.Marhtaler, K.Saito, T.Fujisawa, G.Schön, Phys. Rev. B 81, 125331 (2010)

[4] YU, D.Golubev, M.Marhtaler, T.Fujisawa, G..Schön, arXiv0908.0229.

[5] S.Nakamura, Y.Yamauchi, M.Hashisaka, K.Chiba, K.Kobayashi, T.Ono, R.Leturcq, K.Ensslin, K.Saito, YU, A.C.Gossard, arXiv:0911.3470, hys. Rev. Lett. 104, 080602 (2010).

Supported by Strategic international cooperative program „Nanoelectronics“

