

原子核における α 粒子の Bose-Einstein 凝縮

大久保茂男 S. Ohkubo
(高知女子大・環境理学科)



2010.9.1

熱場の量子論

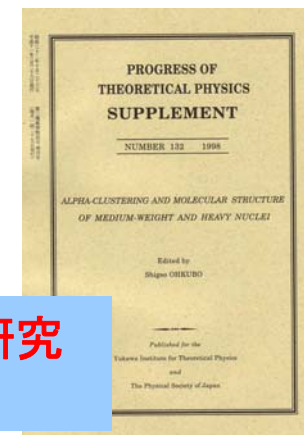
1999

クラスター模型 軽い領域だけでなく重い領域 ^{40}Ca - ^{44}Ti 領域での
成立 理論、実験

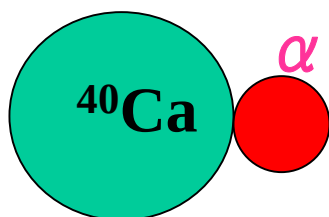
1998 PTP Supplement 132 (山屋 弁 追悼記念)

fp殻領域

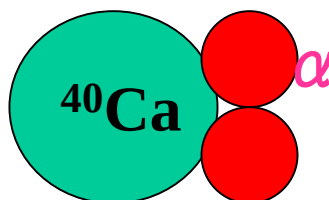
重い核の領域へのクラスター研究
の道が切り開かれた



^{44}Ti



^{48}Cr



クラスター模型の歴史と展開

2005 基礎研シンポジウム

学問の系譜—アインシュタインから湯川・朝永へ

素粒子論研究2006年3月号 Vol 112 No.6

物性研究 2006年6月号 Vol 86 No.3

(ダウンロード可能)

Schuck の質問 1999 Rab (Croatia) cluster conference

2010.9.1 熱場の量子論
 ^{48}Cr の $^{40}\text{Ca} + \alpha + \alpha$ クラスター構造 の講演

1983

α 粒子の多体系でBose-Einstein 凝縮はあるか？

Gambhir, Ring and Schuck:PRL 51, 1235 (1983):

A superfluid condensate of alpha particles ?

In conclusion we can say that our BCS-like theory for bosons in conjunction with the IBM2 Hamiltonian suggests that open-shell nuclei (open in protons and neutrons) may form a superfluid condensate of α particles rather than separate superfluid phases of proton and neutron pairs.

VOLUME 51, NUMBER 14

PHYSICAL REVIEW LETTERS

3 OCTOBER 1983

Nuclei: A Superfluid Condensate of α Particles? A Study within the Interacting-Boson Model

Y. K. Gambhir^(a) and P. Ring

Physik-Department, Technische Universität München, D-8046 Garching, West Germany

and

P. Schuck

Institut des Sciences Nucléaires, F-38042 Grenoble, France

(Received 11 July 1983)

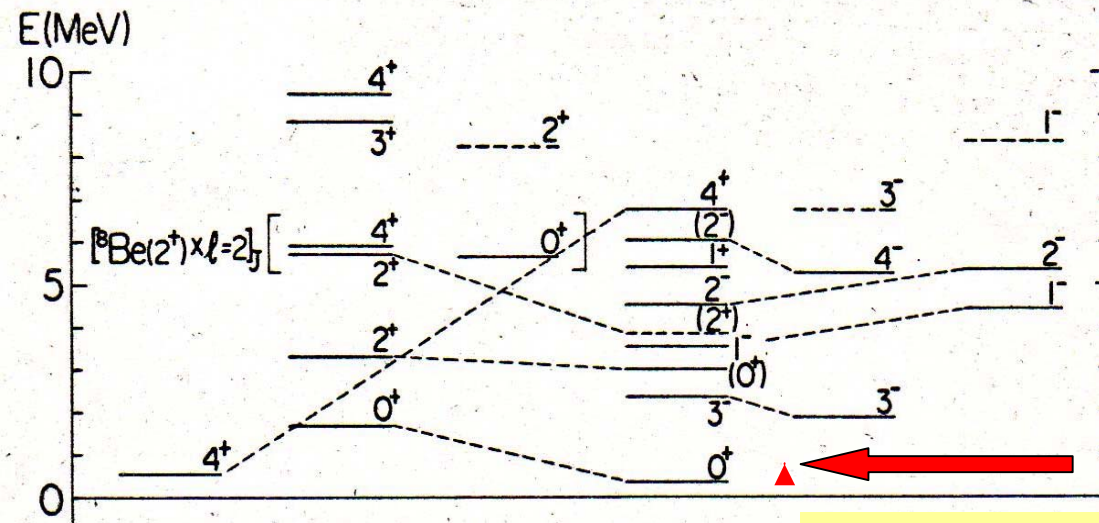
The authors have studied the question of whether pairs of neutrons and pairs of protons of the usual superfluid phases form a bound state to give rise to a superfluid condensate of " α particles." They indeed find indications for this to be the case from a BCS-like study for bosons using the proton-neutron interacting-boson model as well as from an even-odd effect in the number of pairs using experimental binding energies.

1977

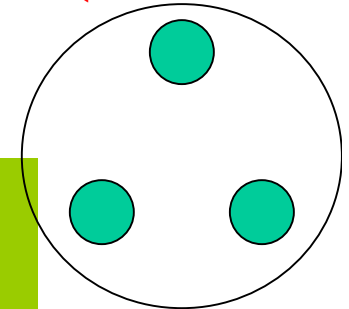
Uegaki 3 alpha cluster model

:Energy level ^{12}C (PTP 57,1262(1977))

GCM



3個の α がゆるやか結合(気体)



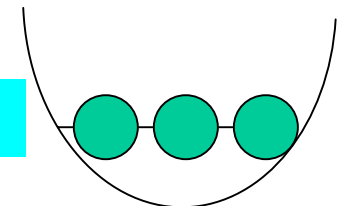
3α Hoyle state (7.65MeV)

閾値エネルギー (7.27MeV)

閾値からのエネルギー (0.38MeV)

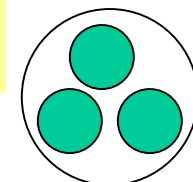
崩壊幅 8.7 eV

0S state



基底状態

3個の α がコンパクトに結合(固体)



1977

Kamimura

3 α RGM

INT. CONF. NUCLEAR STRUCTURE, TOKYO, 1977
 J. PHYS. SOC. JAPAN 44 (1978) SUPPL. p. 225-231

**V-B-a The Resonating Group Method for 3-Cluster Systems Based
 on the Use of Generator Coordinate Kernels**
 —Application to the 3α Structure of ^{12}C and the 3α Reaction in Stellar Evolution—

Y. FUKUSHIMA and M. KAMIMURA*

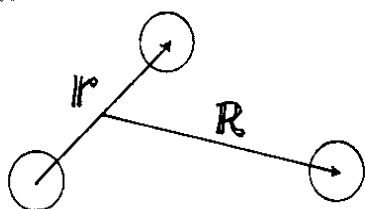
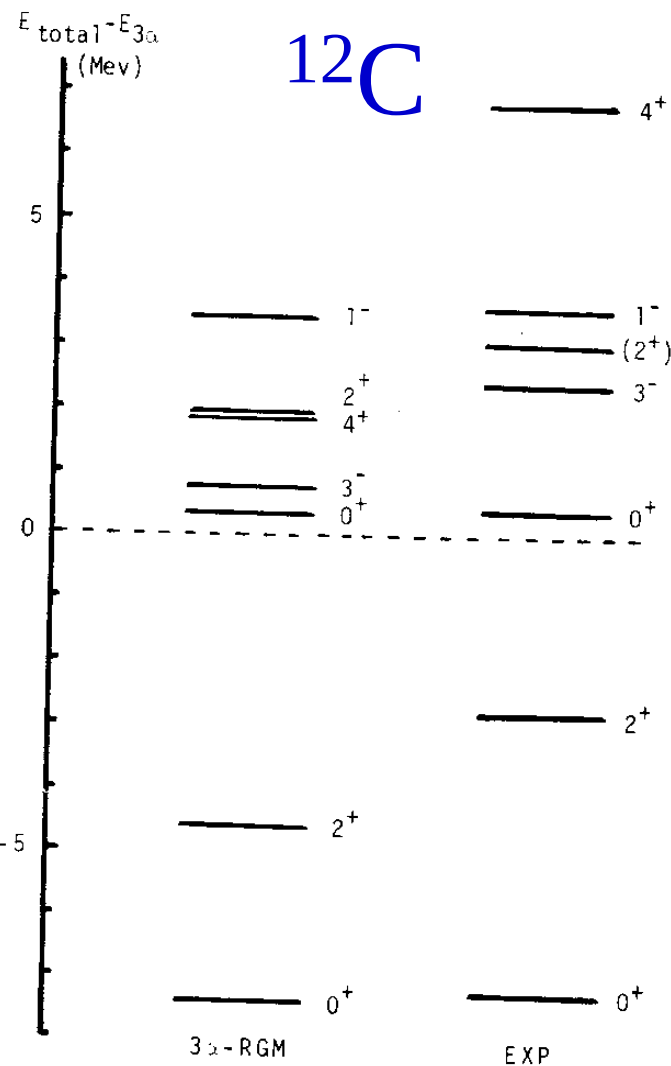


Fig. 1. Relative coordinates of the 3α system.



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熱場の量

2004

Hoyle stateのdilute propertyの実験的検証

Rainbow, prerainbow

Bose-Einstein凝縮と 核虹のエアリー構造

RAPID COMMUNICATIONS

PHYSICAL REVIEW C **70**, 041602(R) (2004)

Bose-Einstein condensation of α particles and Airy structure in nuclear rainbow scattering

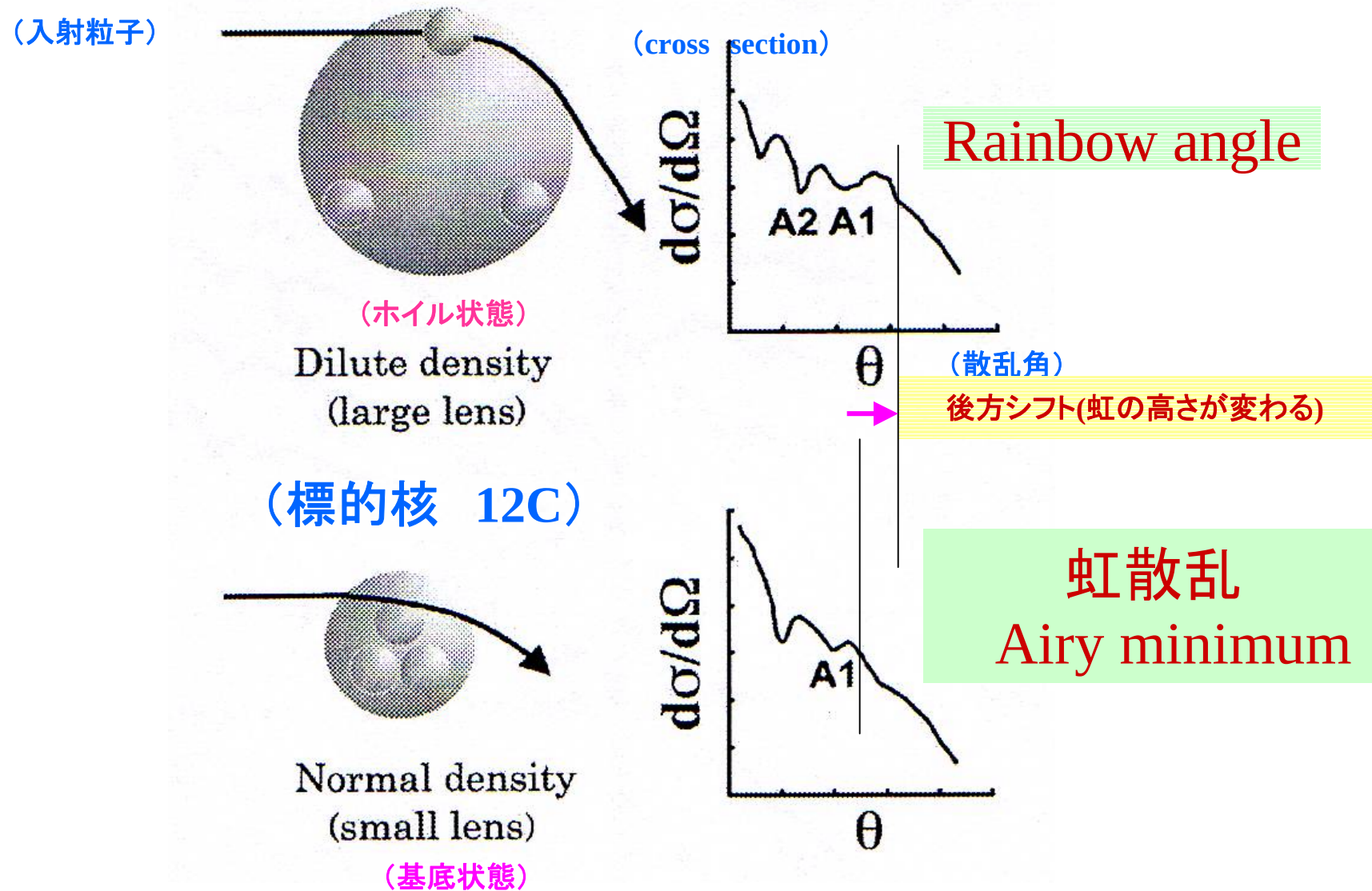
S. Ohkubo¹ and Y. Hirabayashi²

¹*Department of Applied Science and Environment, Kochi Women's University, Kochi 780-8515, Japan*

²*Information Initiative Center, Hokkaido University, Sapporo 060-0810, Japan*

(Received 4 December 2003; published 25 October 2004)

予想: Illustrative Figure of Airy structure and scattering



解析： Hoyle 状態 の構造とエアリー構造

$\alpha + {}^{12}\text{C}$ 散乱のチャネル結合法による解析

- ${}^{12}\text{C}$ の波動関数 3α RGM~ボーズ凝縮の波動関数
- 密度依存DDM3Y有効相互作用による畳込み

2004

$\alpha + {}^{12}\text{C}$ scattering

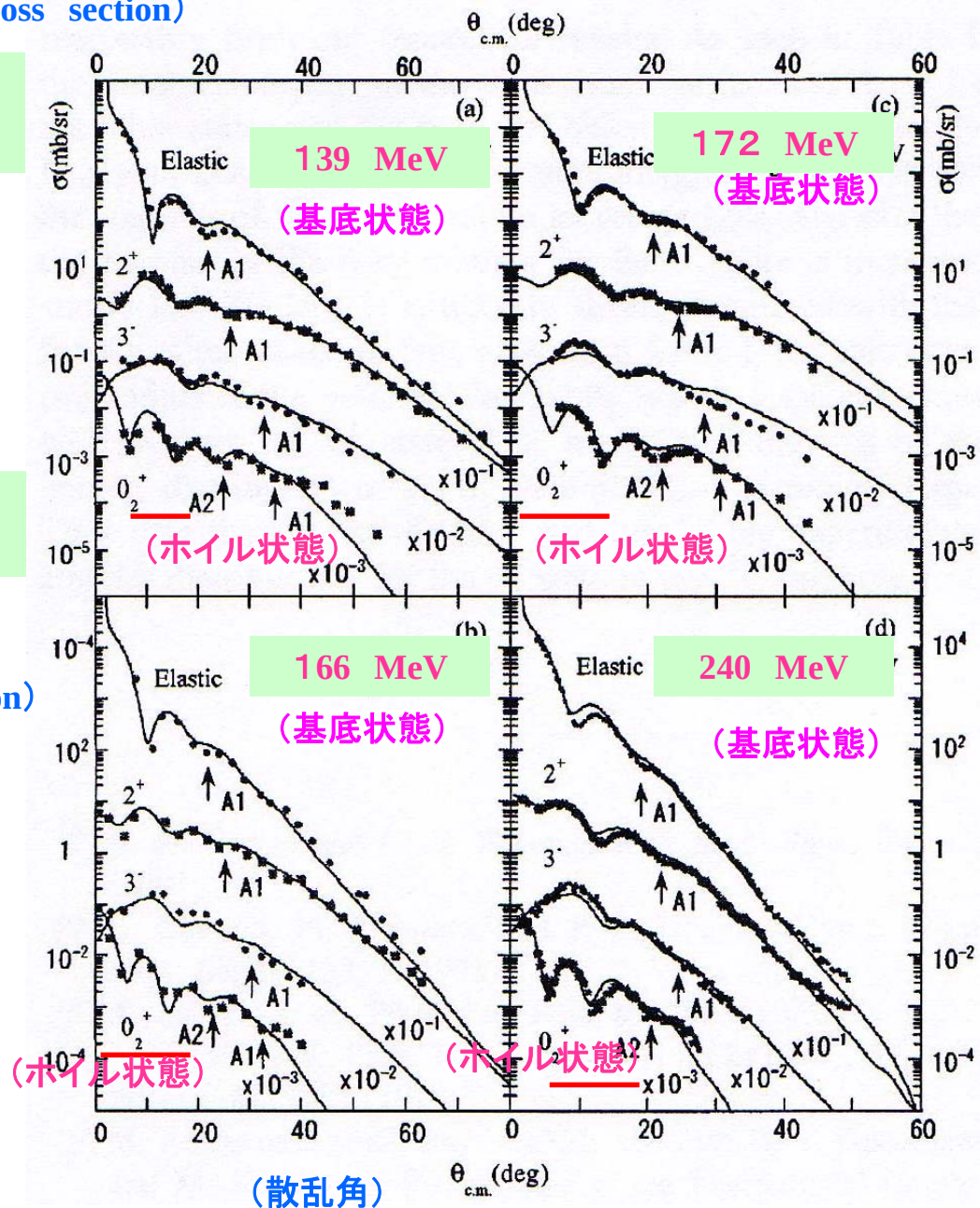
rainbow

Airy minimum

(cross section)

BOSE-EINSTEIN CONDENSATION OF α PARTICLES AND AIRY

(cross section)



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Potential parameters, rms, and volume integrals

TABLE I. The normalization factor N_R , volume integral per nucleon pair J_V , rms radius $\langle R^2 \rangle^{1/2}$, of the folding potential, and the strength of the imaginary potential. The radius $R_W=4.7$ fm and diffuseness $a_W=0.7$ fm of the imaginary potentials are used.

	E_L (MeV)	N_R	J^π	J_V (MeV fm ³)	$\langle R^2 \rangle^{1/2}$ (fm)	W (MeV)
(基底状態)	139	1.23	0_1^+	294	<u>3.484</u>	8.0
			2^+	291	3.469	10.0
			3^-	327	3.742	14.0
			(ホイル状態) 0_2^+	<u>366</u>	<u>4.304</u>	28.0
(基底状態)	166	1.26	0_1^+	286	<u>3.499</u>	7.0
			2^+	283	3.485	10.0
			3^-	320	3.752	14.0
			(ホイル状態) 0_2^+	<u>358</u>	<u>4.310</u>	27.0
(基底状態)	172.5	1.26	0_1^+	292	<u>3.506</u>	7.5
			2^+	289	3.493	9.0
			3^-	327	3.758	12.0
			(ホイル状態) 0_2^+	<u>367</u>	<u>4.315</u>	25.0
(基底状態)	240	1.42	0_1^+	264	<u>3.532</u>	7.5
			2^+	261	3.519	12.0
			3^-	295	3.777	17.0
			(ホイル状態) 0_2^+	<u>330</u>	<u>4.328</u>	30.0

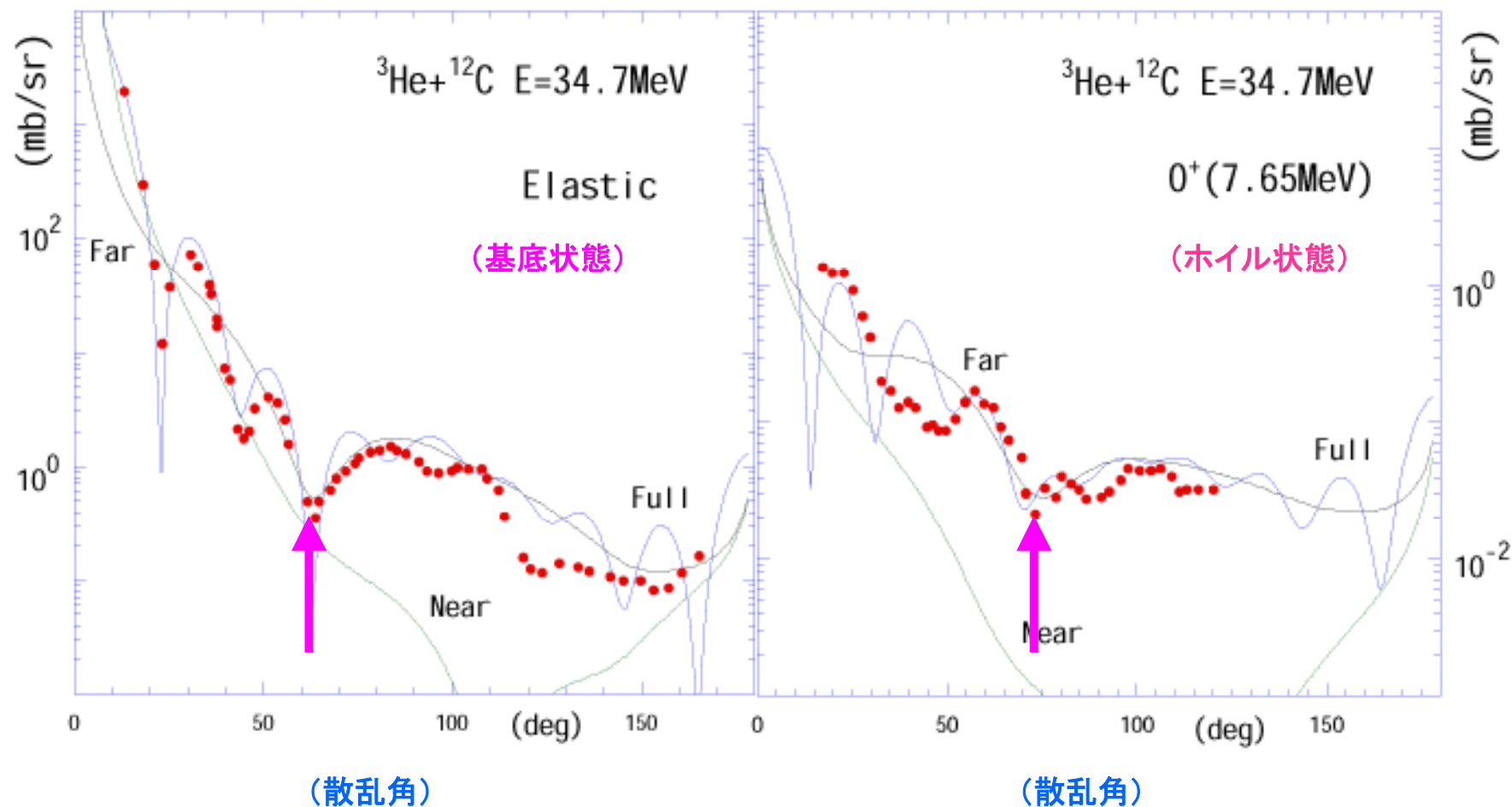
2007

$^3\text{He} + ^{12}\text{C}$ scattering $E=34.7$ MeV

prerainbow N/F decomposition

(cross section)

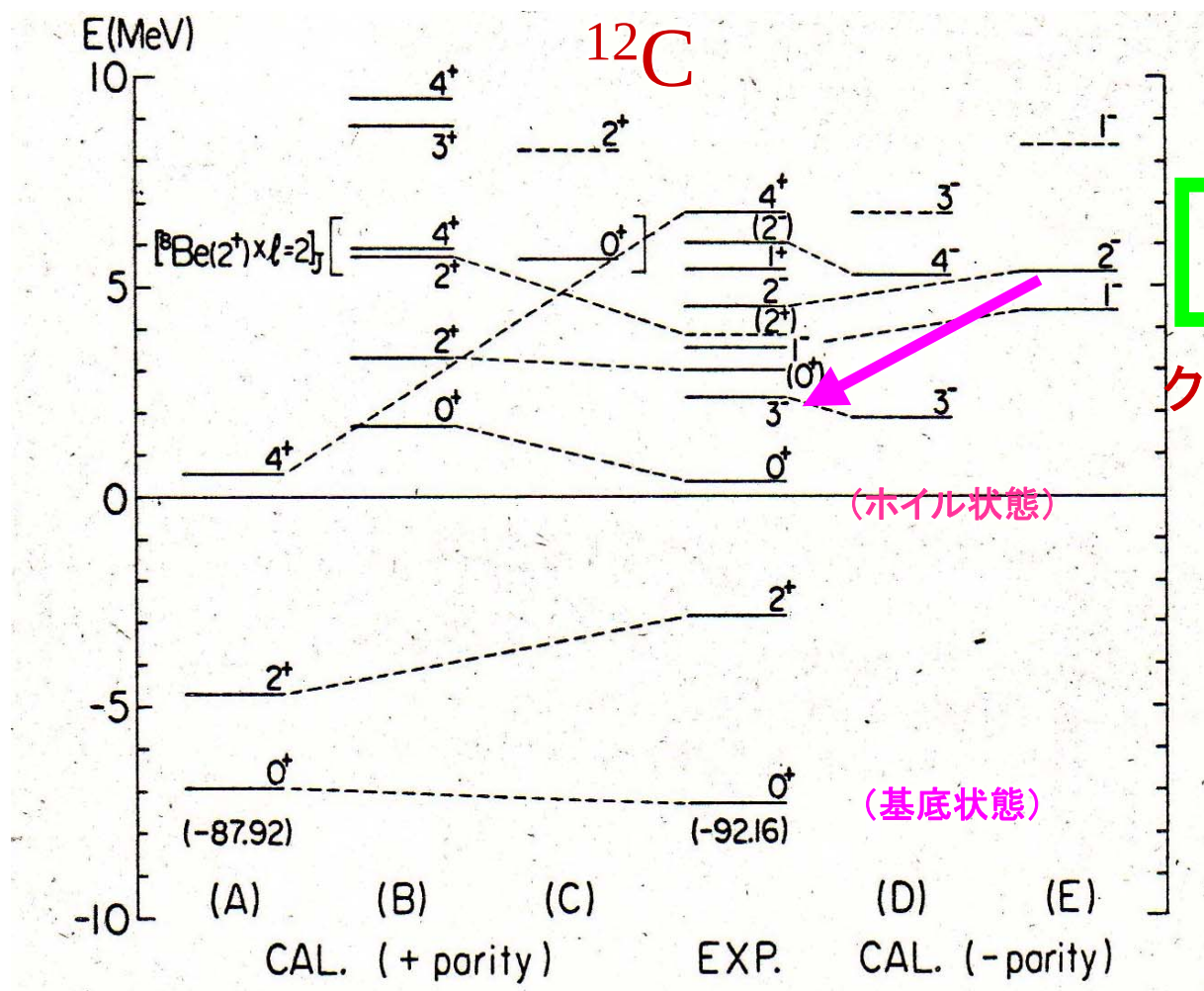
(cross section)



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熱場の量子論

^{12}C には量子渦状態は存在しない



3-は渦状態ではない

クラスターと殻構造の中間的構造をもつ

of angular momentum on rotational fre-
id lines indicate the variational result, and
resent the fully numerical results for the
ne). 10 (dashed line), and 8 (dotted line).

α 粒子はBose・Einstein凝縮しているか？

実体的な物の論理 ？

- ^{12}C のHoyleの3 α 系の密度が希薄であることは実験的に確認
- しかし、超流動や渦は確認されず α 粒子Bose・Einstein凝縮している実験的な直接的証拠はない。

4 α 系: ^{16}O における4 α ボーズ凝縮

確かな実験的証拠？ ？

束縛状態でなく散乱状態として解くべき

閾値近傍の状態でクラスターの発達した状態は

- 1) 散乱条件を正しく取り入れて、
- 2) 散乱状態で適切な有効相互作用を用いて解かれるべき。

(有効相互作用のstarting energy dependence

有効相互作用と模型空間の相互規定性)

Bose-Einstein α 粒子凝縮にともなう 実体的な固有のモードは見られないのか？

3 α 系では dilute property は実験的にも確認されるが、Bose-Einstein α 粒子凝縮による実体的な物の論理(超流動、量子渦など)は実験的には確認できてない。

^{16}O の4 α 構造

有限系の原子核の α 粒子ではどうか？

1. 0^+ のみでなく 励起状態とともに理解されるべき。

(部分のみでなく全体を見るべき: 部分と全体の弁証法)

2. 束縛状態近似ではなく共鳴状態として正確に解くべき。

1967

¹⁶Oの4α構造の研究状況:

Chevallier et. al. 4α linear chain model

VOLUME 160, NUMBER 4

Breakup of O¹⁶ into Be⁸+Be⁸†

P. CHEVALLIER AND F. SCHEIBLING

Institut de Recherches Nucléaires, Strasbourg-Cronenbourg, France

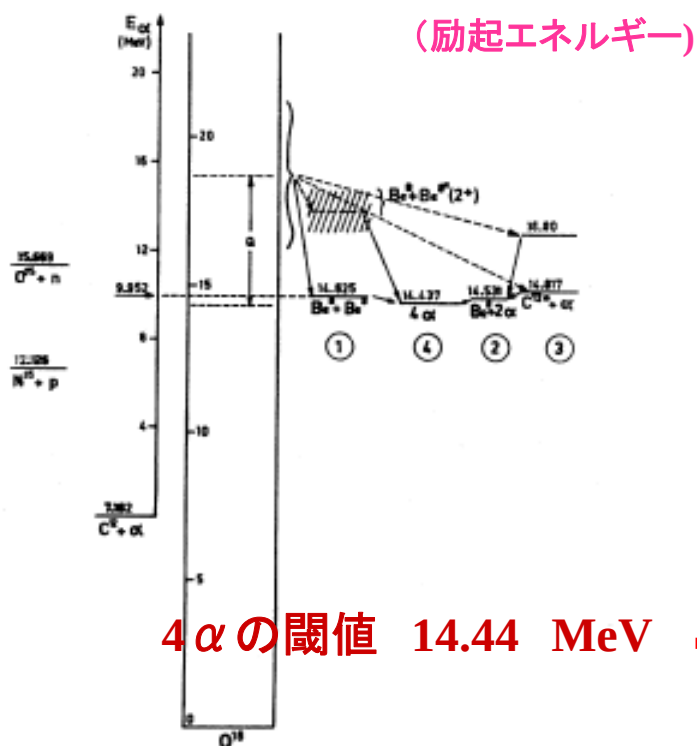
AND

G. GOLDRING, I. PLESSER, AND M. W. SACHS*

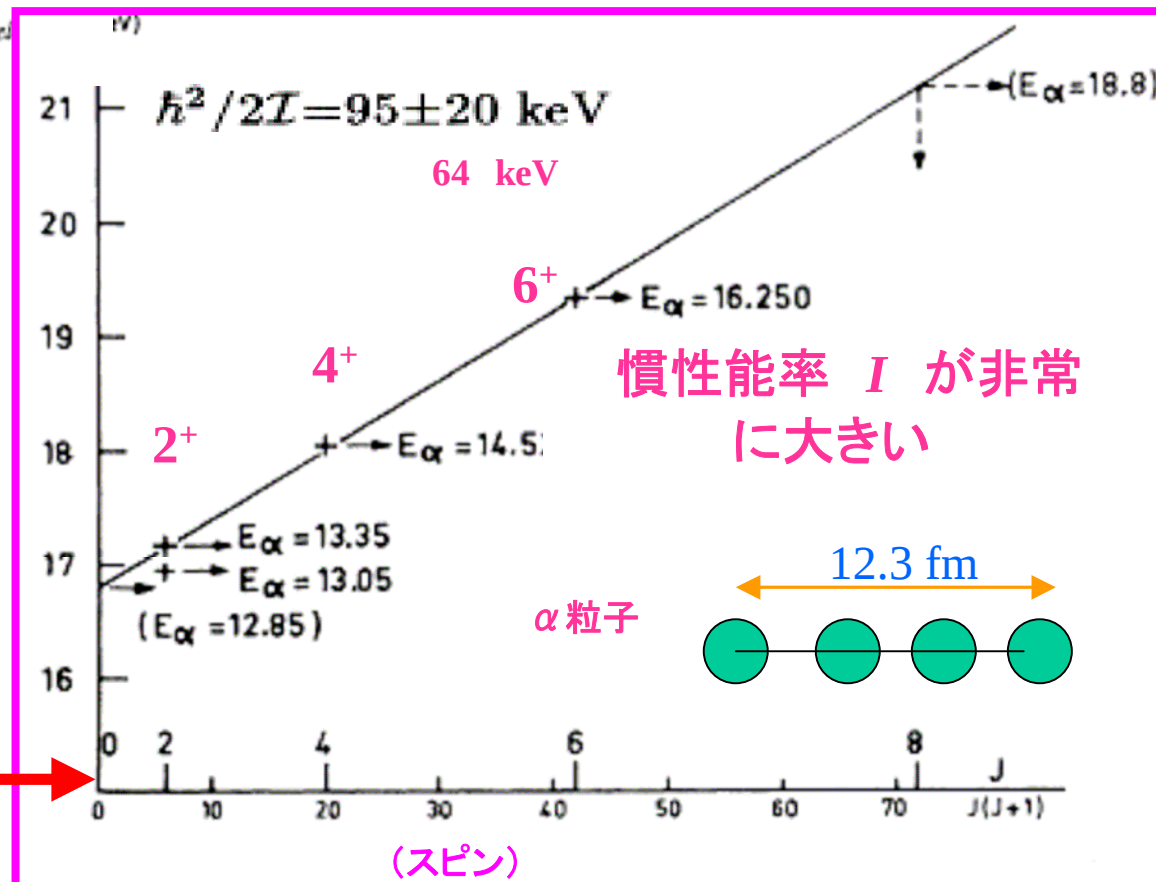
The Weizmann Institute of Science, Rehovoth, Israel

(Received 27 January 1967)

1. O¹⁶ → Be⁸+Be⁸
Be⁸ → 2α
Be⁸ → 2α
2. O¹⁶ → Be⁸+2α
Be⁸ → 2α
3. O¹⁶ → C¹²*(0⁺; 7.66 MeV)+α (and higher states)
C¹²* → Be⁸+α
Be⁸ → 2α
4. O¹⁶ → 4α.



4αの閾値 14.44 MeV



^{16}O の $\alpha + ^{12}\text{C}$ (Hoyle) クラスタ構造

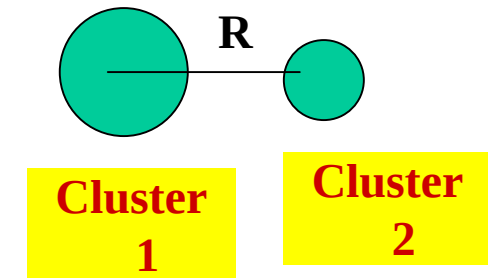
チャネル結合法による $\alpha + ^{12}\text{C}$ 散乱の解析

- ^{12}C の波動関数 3α RGM~ボーズ凝縮の波動関数
 - 密度依存DDM3Y有効相互作用による畳込み

Coupled channel equations

$$\Psi = \sum_i \varphi(\alpha) \varphi_i(^{12}\text{C}) \chi_i(\mathbf{R}), \quad (1)$$

$$\left[-\frac{\hbar^2}{2\mu} \nabla^2 + U_{ii}(\mathbf{R}) - (E - \epsilon_i) \right] \chi_i(\mathbf{R}) = - \sum_{j \neq i} U_{ij}(\mathbf{R}) \chi_j(\mathbf{R}). \quad (2)$$



Double folding
potential

$$V_{ij}(\mathbf{R}) = \int \rho_\alpha(\mathbf{r}_\alpha) \rho_{ij}(\mathbf{r}_T) v(\mathbf{R} - \mathbf{r}_T + \mathbf{r}_\alpha) d\mathbf{r}_T d\mathbf{r}_\alpha, \quad (3)$$

renormalization
factor $N_R + iN_I$

$$U_{ij}(\mathbf{R}) = N_R V_{ij}(\mathbf{R}) + iW(\mathbf{R}) \delta_{ij}, \quad (4)$$

DDM3Y effective interaction

$$v_{NN}(E, \rho; s) = g(E, s) f(E, \rho). \quad f(E, \rho) = C(E) [1 + \alpha(E) e^{-\beta(E)\rho}],$$

取り入れる ^{12}C のchannel

g.s., 2_1^+ (4.44 MeV), 3^- (9.65 MeV), 0_2^+ (7.65 MeV) and 2_2^+ (10.3 MeV). The absorption due to the coupling to all the other

Imaginary potential

$E_L = 18 \text{ MeV}$

open channels, i.e., $p + ^{15}\text{N}$, $n + ^{15}\text{O}$ and $d + ^{14}\text{N}$ channels

$$N_I = 0.045$$

Real potential

$$E_\alpha = 139, 166 \text{ and } 172.5 \text{ MeV}$$

$$N_R = 1.23 - 1.26$$

$E_L = 18 \text{ MeV}$

$$N_R = 1.398$$

High energy alpha + scattering

$E_L = 172.5 \text{ MeV}$ 166 MeV

(cross section)

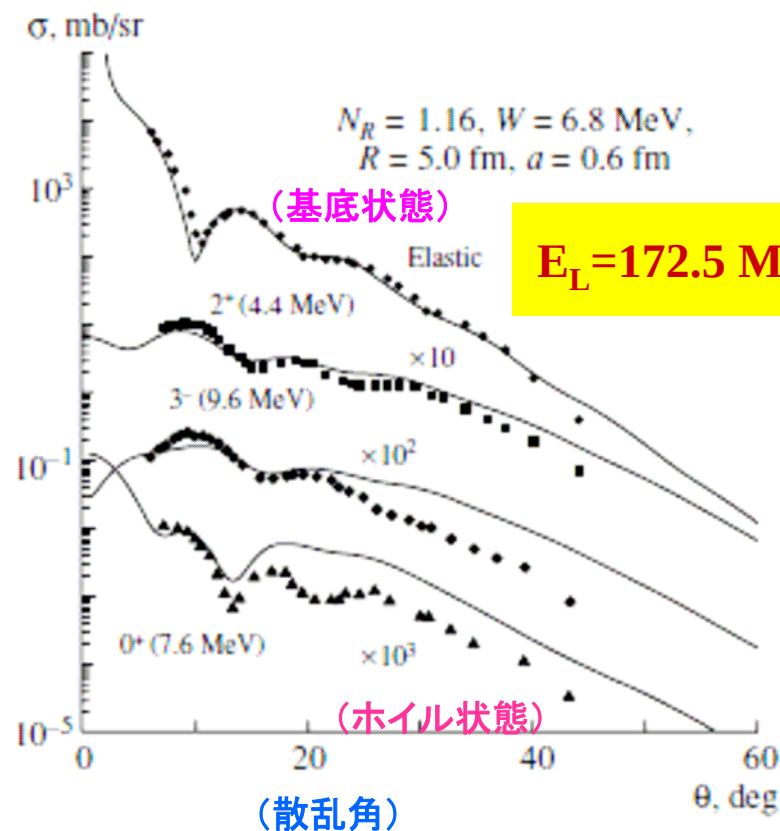


Fig. 1. The experimental data [13] of elastic and inelastic α scattering from ^{12}C at $E_\alpha = 172.5 \text{ MeV}$ are compared with the coupled-channel calculations (solid curves).

2010.3.1

(cross section)

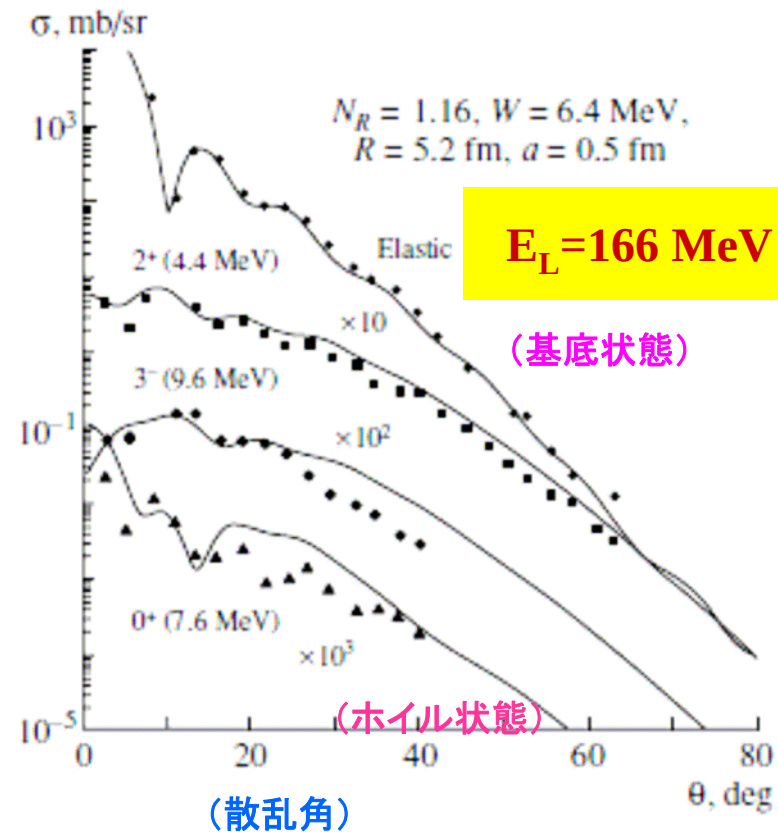


Fig. 2. The experimental data [14] of elastic and inelastic α scattering from ^{12}C at $E_\alpha = 166 \text{ MeV}$ are compared with the coupled-channel calculations (solid curves).

High energy alpha + scattering

$E_L = 139 \text{ MeV}$ 106 MeV

(cross section)

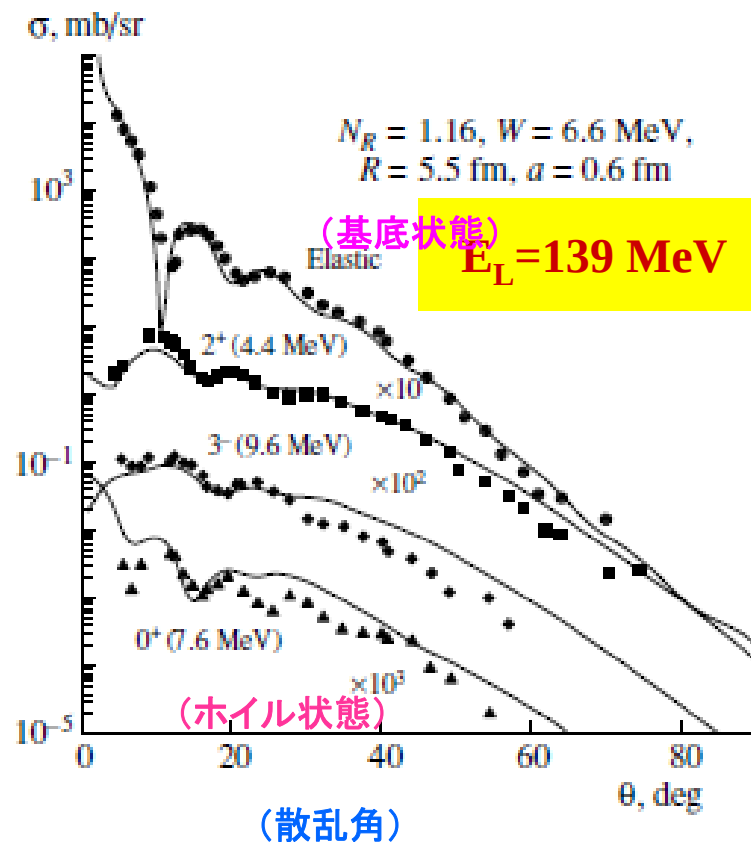


Fig. 3. The experimental data [15] of elastic and inelastic α scattering from ^{12}C at $E_\alpha = 139 \text{ MeV}$ are compared with the coupled-channel calculations (solid curves).

2010.3.1

(cross section)

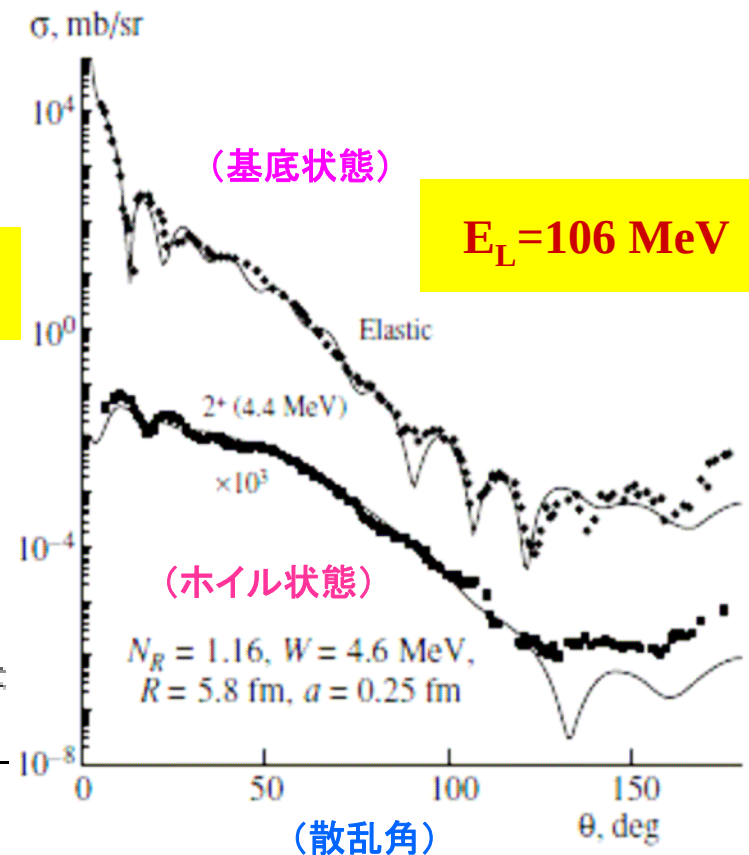


Fig. 4. The experimental data [17] of elastic and inelastic α scattering from ^{12}C at $E_\alpha = 104 \text{ MeV}$ are compared with the coupled-channel calculations (solid curves).

核物理の基礎

High energy alpha + scattering

$E_L = 82 \text{ MeV}$ 65 MeV

(cross section)

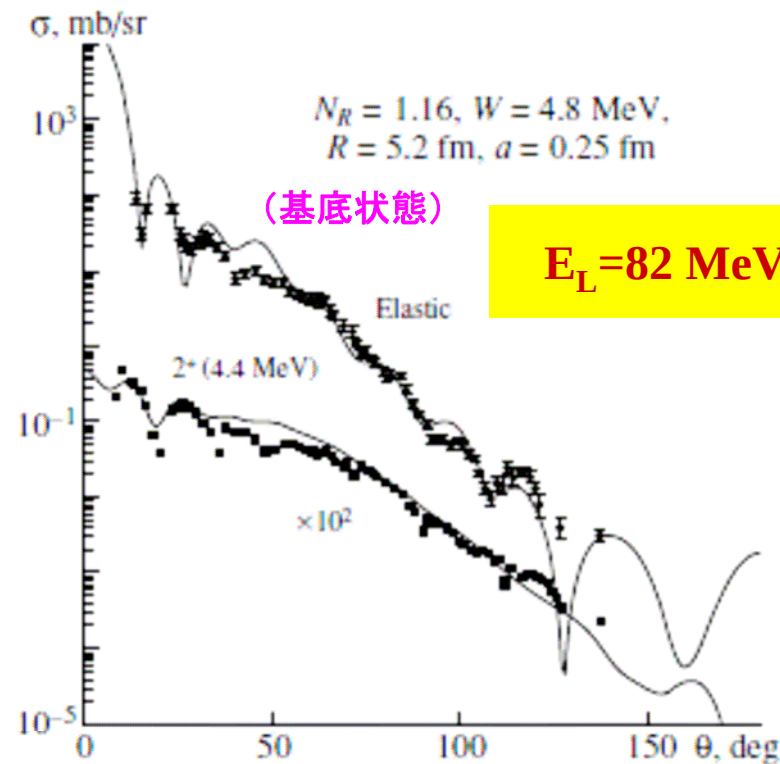


Fig. 5. The experimental data [19] of elastic and inelastic α scattering from ^{12}C at $E_\alpha = 82 \text{ MeV}$ are compared with the coupled-channel calculations (solid curves).

(cross section)

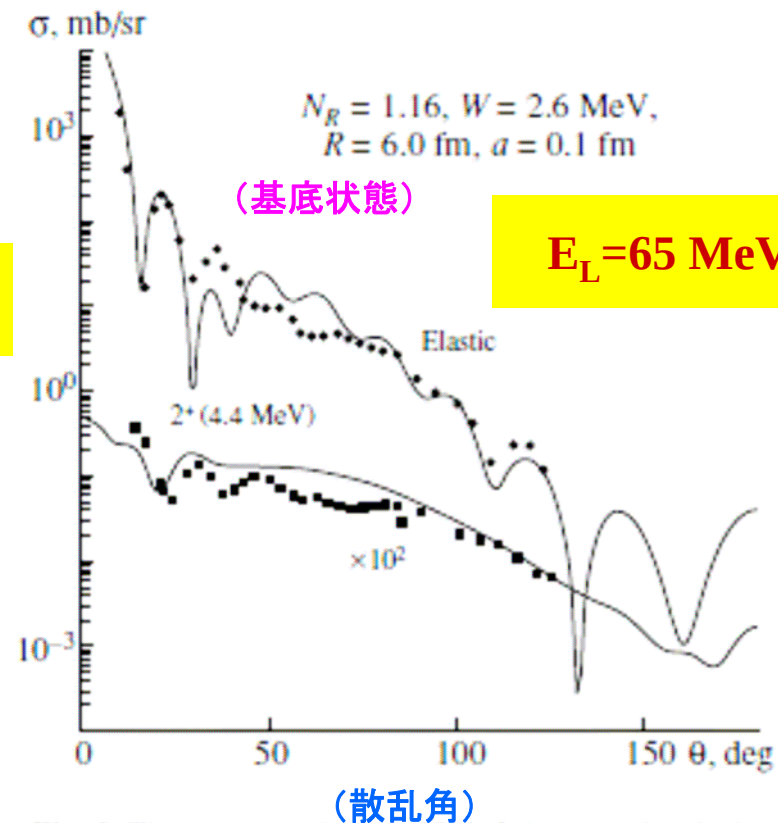
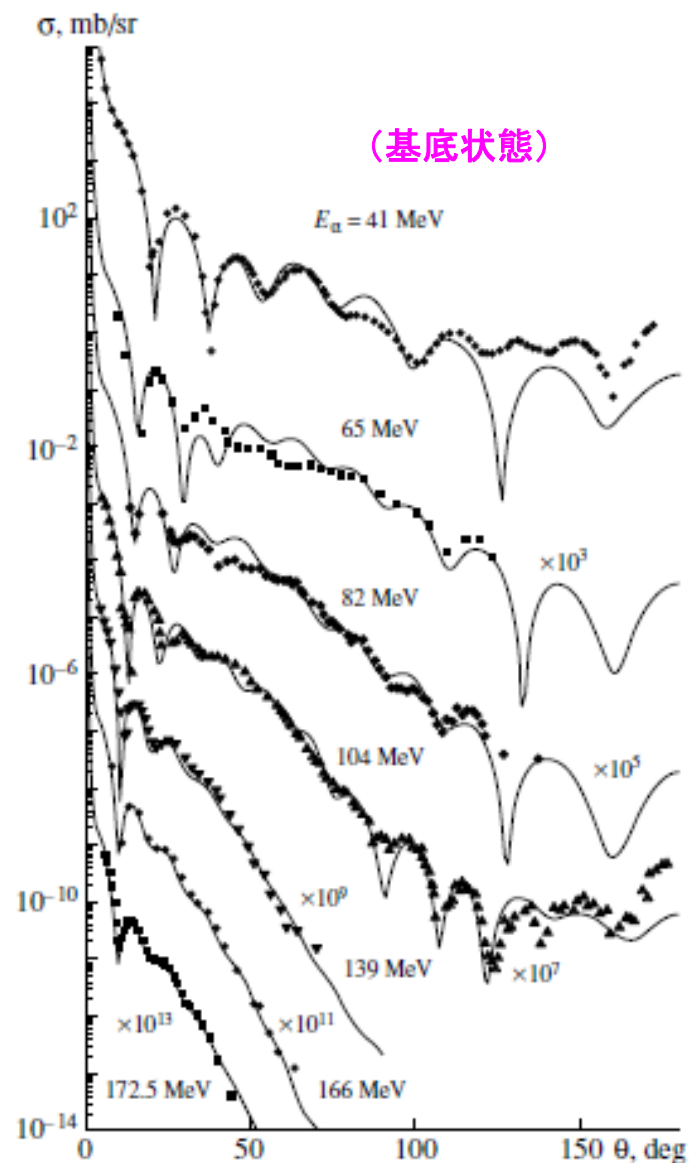


Fig. 6. The experimental data [20] of elastic and inelastic α scattering from ^{12}C at $E_\alpha = 65 \text{ MeV}$ are compared with the coupled-channel calculations (solid curves).

High energy alpha + scattering

$E_L = 172 \text{ MeV} - 41 \text{ MeV}$ (elastic)

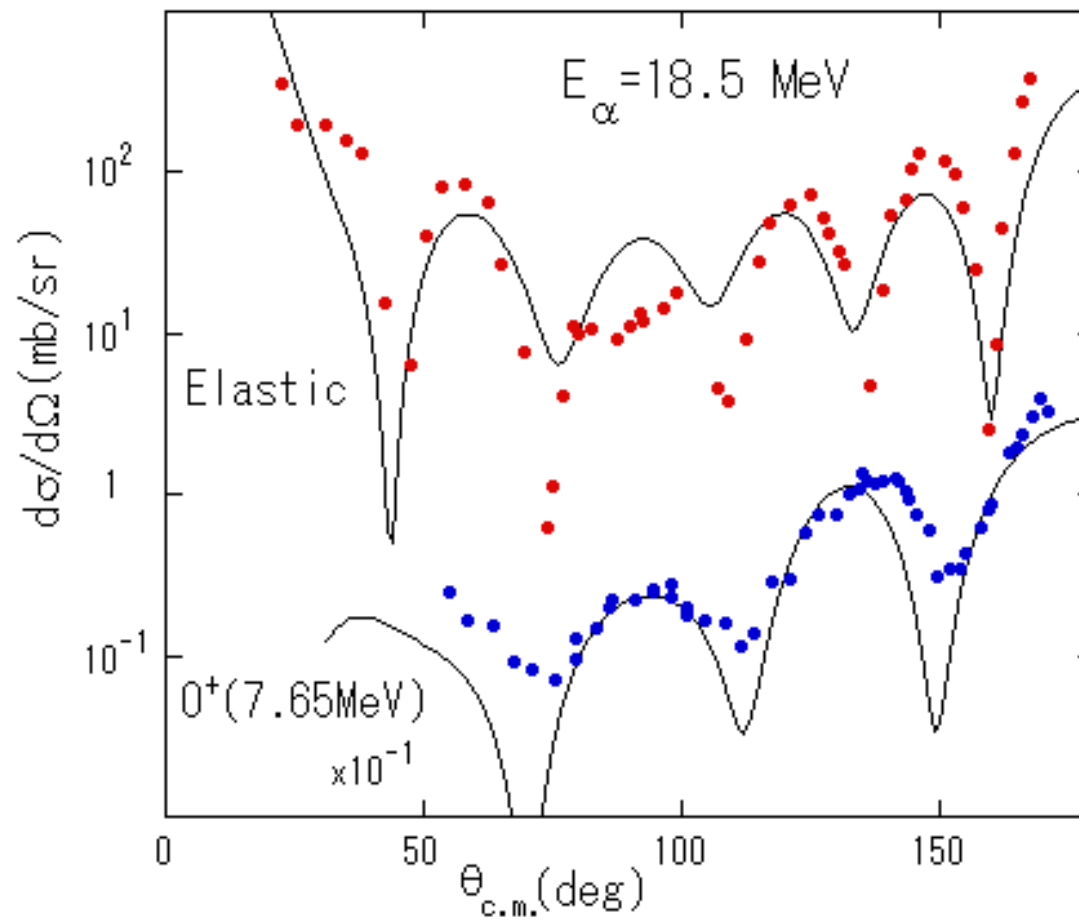


Normalization factor N_R and volume integral per nucleon pair J_V (in MeV fm^3)

$E_\alpha, \text{ MeV}$	N_R	J_V	J_V [28]	J_V [15]
41	0.90	246		
65	1.16	313		
82	1.16	305		
104	1.16	298	393	319
139	1.16	286	353	278
166	1.16	279	326	
172.5	1.16	286		277

Fig. 7. The experimental angular distributions of elastic α scattering from ^{12}C are compared with the coupled-channel calculations (solid curves).

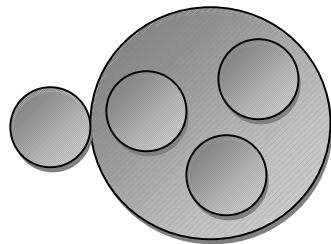
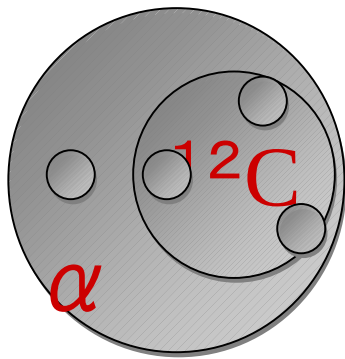
$E_L=18.5$ MeV (lowest energy) $\alpha + {}^{12}\text{C}$ 散乱の解析



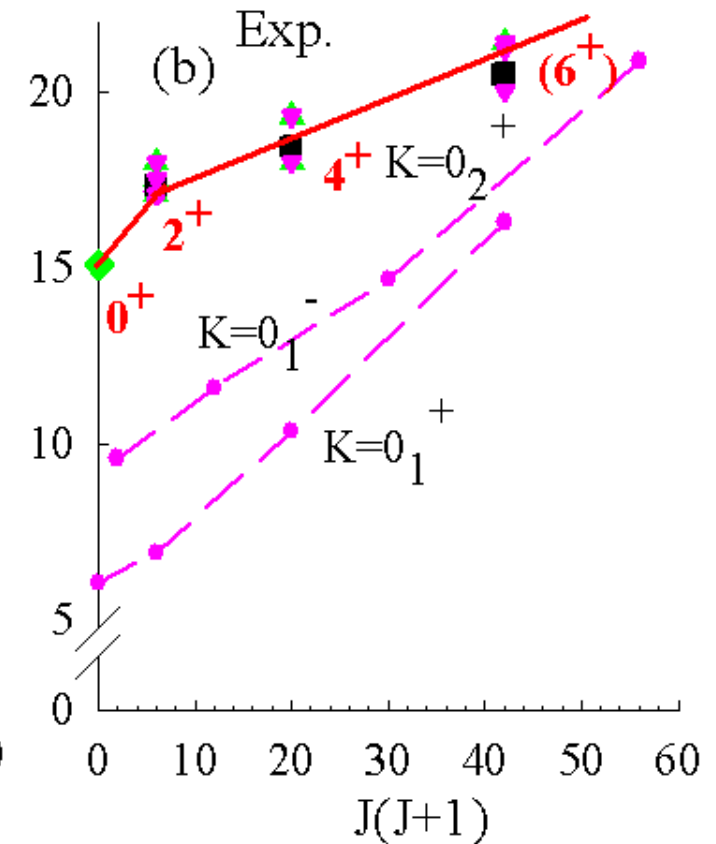
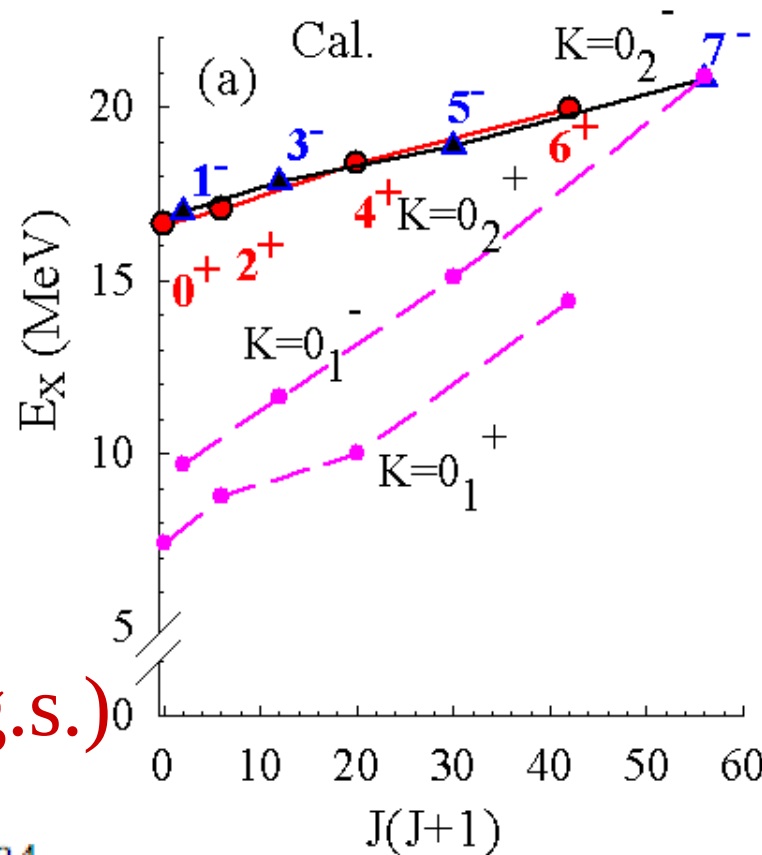
^{16}O の $\alpha + ^{12}\text{C}(\text{Hoyle})$ クラスター構造 をもつエネルギー準位

$\alpha + ^{12}\text{C}(\text{Hoyle})$

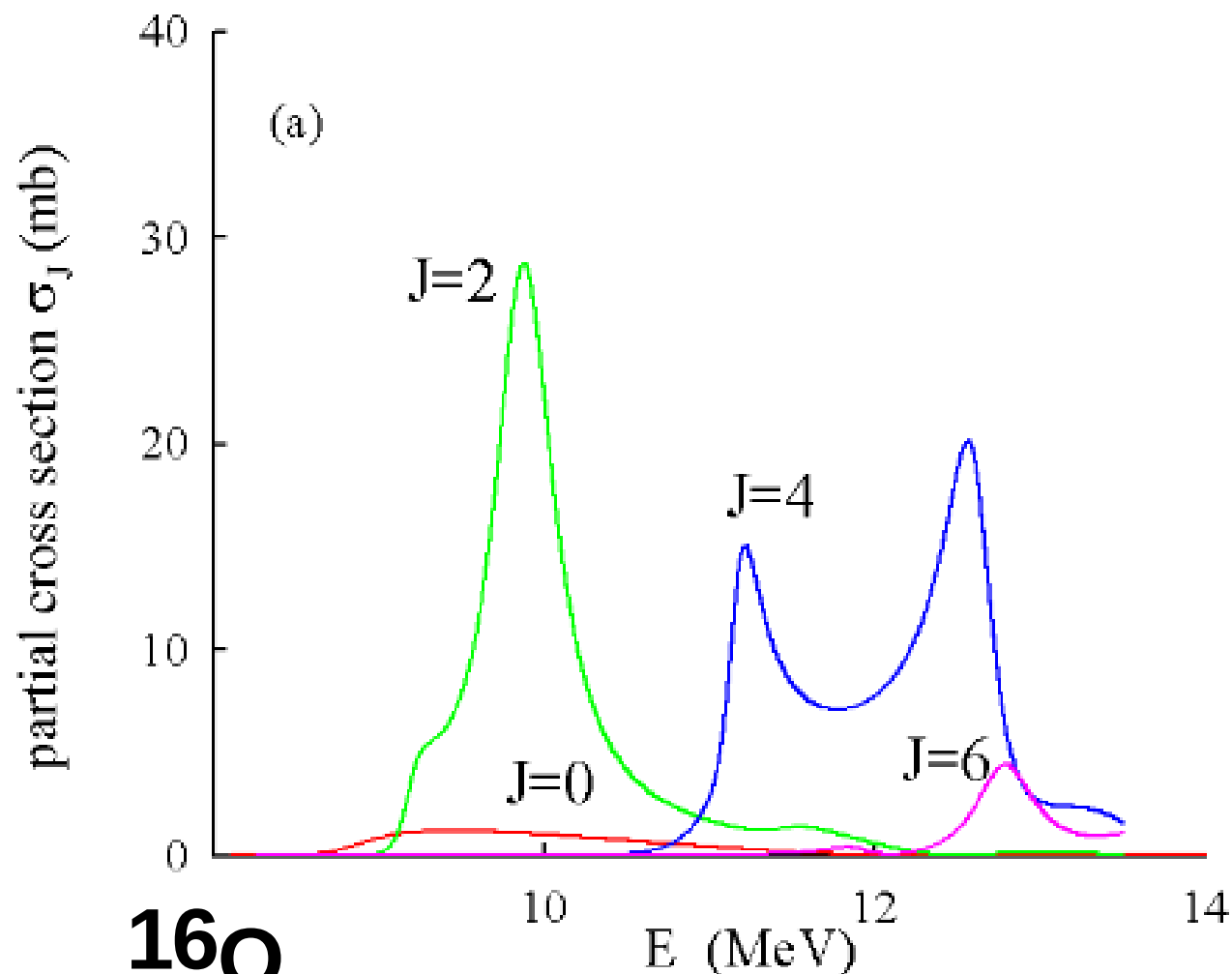
^{16}O



$\alpha + ^{12}\text{C}(\text{g.s.})$



partial cross sections to the Hoyle state



$\alpha + {}^{12}\text{C}$ の閾値(7.16MeV)からのエネルギー

K=0₁⁺, K=0₁⁻ band widths

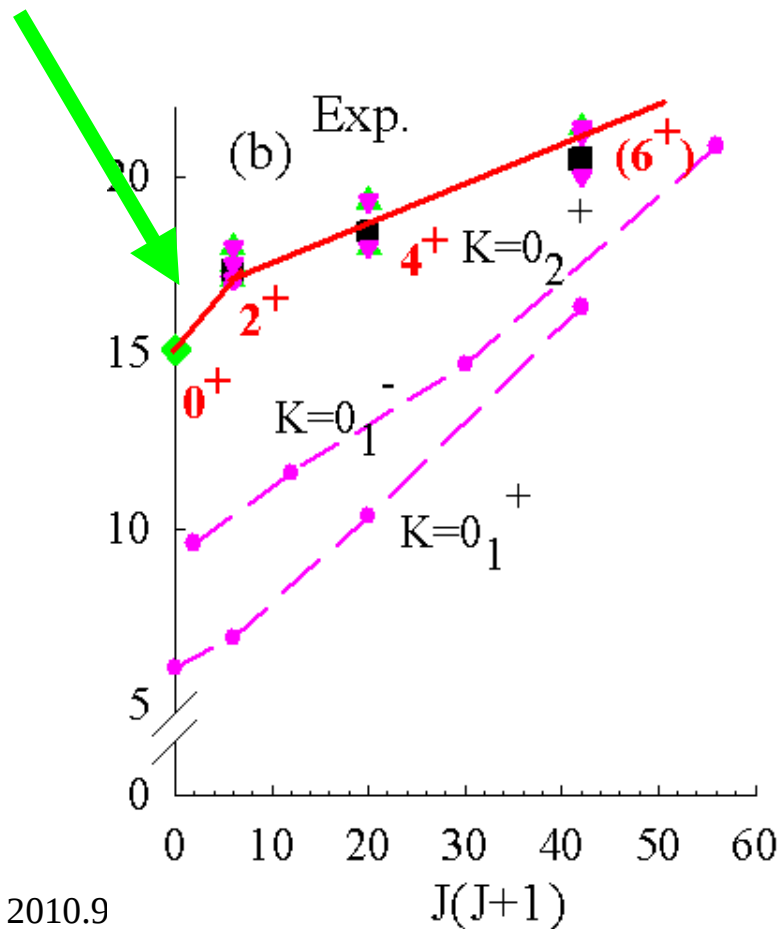
		cal.					exp.				
J^π	E_x	Γ_α	$\theta_\alpha^2(a)$				E_x	Γ_α	$\theta_\alpha^2(a)$		
	(MeV)	(keV)	a (fm)				(MeV)	(keV)	a (fm)		
			5.2	5.6	6.0				5.2	5.6	6.0
K=0 ₁ ⁺	4 ⁺	10.00	5	0.15	0.09	0.06	10.36	26±3	0.35	0.22	0.14
	6 ⁺	14.37	34	0.10	0.06	0.04	16.28	420±20	0.36	0.26	0.20
K=0 ₁ ⁻	1 ⁻	9.67	778	1.02	0.87	0.79	9.59	480±20	0.71	0.61	0.54
	3 ⁻	11.61	822	0.64	0.54	0.48	11.60	800±100	0.63	0.53	0.47
	5 ⁻	15.09	440	0.26	0.21	0.18	14.66	672±11	0.49	0.38	0.31
	7 ⁻	20.86	1790	0.77	0.59	0.49	20.86	904±55	0.39	0.30	0.25

K=0₂⁺, K=0₂⁻ band widths

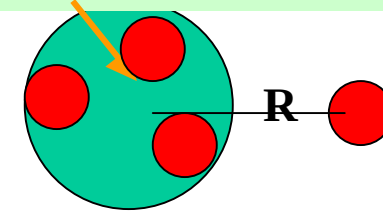
J^π		E_x	Γ	J^π		E_x	Γ
		(MeV)	(MeV)			(MeV)	(MeV)
$K = 0_2^+$	0^+	16.61	1.14	$K = 0_2^-$	1^-	16.98	0.57
	2^+	17.04	0.45		3^-	17.83	0.27
	4^+	18.38	0.23		5^-	18.86	0.24
	6^+	19.95	0.39		7^-	20.81	1.17

superfluidity

Reduction of moment of inertia



Local α condensation



$^{12}\text{C}(0_2^+)$

α

Large moment of inertia I
 $I = I(\alpha) + I(^{12}\text{C}(0_2^+)) + I(\text{rel})$

56% 40%

$\rightarrow R=5.9 \text{ fm}$
 cf. 4 α linear chain model
 $R=12.3 \text{ fm}$

15.1 MeV 0⁺ state ($\Gamma = 186 \text{ keV}$)

$\Gamma_{\alpha} / \Gamma = 0.35$

0⁺ states:

15.1 MeV 0⁺

17.6 MeV (0⁺, 1⁻)

18.1 MeV (0⁺, spin tentative)

量子論

まとめ

- 1) ^{16}O の4 α 閾値近傍に local condensed $\alpha + ^{12}\text{C}$ (Hoyle)クラスター構造をもつ状態が存在する
(4 α chainにあらず)
- 2) 大きい慣性能率: 大きい半径をもつ
- 3) 15.1MeV O^+ 状態: 慣性能率が4分の1に減少
 α 凝縮状態の可能性、超流動の特徴
- 4) Local α condensation が広く存在する可能性