基研研究会 熱場の量子論とその応用, 2010年8月30日~9月1日

Bose-Einstein condensation in cold atoms and Hawking radiation

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Bose-Einstein condensation in cold atoms



Black Hole





BHは真っ黒ではない BHの蒸発!?

$$n(E) = \frac{1}{\exp\left(\frac{E}{k_B T_H}\right) - 1}$$



(量子)流体と曲がった時空のアナロジー





♥ W.G. Unruh, PRL 46, 1351 (1981).

● 坂上雅昭,大橋憲,流体でのHawking輻射,物性研究 76, 328 (2001).

Outline

- 1. Bose-Einstein condensation
- 2. BEC in cold atoms
- 3. Black hole and Hawking radiation
- 4. "Hawking radiation" in BEC
 - Basic formulation
 - "Horizon" creation?
 - "Hawking radiation"

References

Y. Kurita and TM, Formation of a sonic horizon in isotropically expanding Bose-Einstein condensates, Phys. Rev. A 76, 053603 (2007).
Y. Kurita, M. Kobayashi, TM, M. Tsubota, and H. Ishihara, Spacetime analog of Bose-Einstein condensates: Bogoliubov-de Gennes formulation, Phys. Rev. A 79, 43616 (2009).
栗田泰生,小林未知数,森成隆夫,坪田誠,石原秀樹、 ボース・アインシュタイン凝縮体における粒子生成-曲った時空 とのアナロジー 日本物理学会誌第65巻第3号, 187 (2010).

Bose-Einstein condensation



BEC in cold atoms



原子をたくさん集めて冷やす







BEC in cold atoms



調和振動子ポテンシャル中の粒子のBEC

Advantages for using cold atoms

★系のパラメータを制御可能
★相互作用(Feshbach共鳴)
★閉じ込めポテンシャル
★格子(optical lattice)
★内部自由度

様々なtopological excitations

☆boson-fermion 混合系



*測定手段が限られている

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重力があるときの運動方程式 $m_I \frac{d^2 r}{dt^2} = \frac{GM}{r^2} m_G = g(r) m_G$ 加速系 $r' = r - \frac{1}{2}g(r)t^2$ $m_I \frac{d^2 r'}{dt^2} = \left(m_G - m_I \right) g(r) \longrightarrow$ Einstein equation $R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi G}{c^4}T_{\mu\nu}$

"Spacetime tells matter how to move; matter tells spacetime how to curve."

Black Hole

球対称な解: Schwarzschild solution

$$ds^{2} = \left(1 - \frac{r_{g}}{r}\right)c^{2}dt^{2} - \frac{dr^{2}}{1 - \frac{r_{g}}{r}} - r^{2}d\theta^{2} - r^{2}\sin^{2}\theta d\phi^{2}$$

Horizon半径:
$$r_g = \frac{2GM}{c^2}$$

$$r_g^{Sun} = 3 \text{km}$$

$$r_g^{Earth} = 1 \text{cm}$$

$$r_g^{TM} = 1 \times 10^{-15} \text{ \AA} = 6 \times 10^9 \ell_P$$





Black hole 近傍を通過する光



Schwarzschild時空での平面波









Hawking radiation



Schwarzschild spacetimeの書き換え

1

$$ds^{2} = \left(1 - \frac{r_{g}}{r}\right)c^{2}dt^{2} - \frac{dr^{2}}{1 - \frac{r_{g}}{r}} - r^{2}d\theta^{2} - r^{2}\sin^{2}\theta d\phi^{2}$$
$$\int cd\tilde{t} = cdt - \frac{v/c}{1 - v^{2}/c^{2}}dr, \quad v = -c\sqrt{\frac{r}{r}}$$

$$ds^{2} = c^{2}d\tilde{t}^{2} - \left(dr - vd\tilde{t}\right)^{2} - r^{2}d\theta^{2} - r^{2}\sin^{2}\theta d\phi^{2}$$



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"Hawking radiation" in BEC
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1) Basic formulation

2) "Horizon" creation?

3) "Hawking radiation"

Basic formulation

□ BECの記述: Gross-Pitaevskii方程式

BECにおけるゆらぎ:Bogoliubov-de Gennes 方程式

BECの記述:Gross-Pitaevskii方程式



ゆらぎの記述: Bogoliubov - de Gennes equation

boson field operator

$$\phi = \sum_{\alpha} \phi_{\alpha} b_{\alpha} = \phi_{0} b_{0} + \sum_{\alpha \neq 0} \phi_{\alpha} b_{\alpha}$$
$$b_{0} | g.s. \rangle = \sqrt{N_{0}} | g.s. \rangle$$
$$b_{0}^{\dagger} | g.s. \rangle = \sqrt{N_{0} + 1} | g.s. \rangle \approx \sqrt{N_{0}} | g.s. \rangle$$
$$\phi \longrightarrow \Phi + \phi$$

$$\left(i\hbar\partial_t\phi = \left(-\frac{\hbar^2}{2m}\nabla^2 + \frac{1}{2}m\omega_{ho}^2r^2 + 2U\Phi^*\Phi\right)\phi + U\Phi^2\phi^\dagger\right)$$

ゆらぎの記述: Bogoliubov - de Gennes equation

$$i\hbar\partial_t\phi = \left(-\frac{\hbar^2}{2m}\nabla^2 + \frac{1}{2}m\omega_0^2r^2 + 2U\Phi^*\Phi\right)\phi + U\Phi^2\phi^\dagger$$

Bogoliubov transformation

$$\phi(\mathbf{r},t) = \sum_{\alpha} \left[A_{\alpha}(\mathbf{r},t) b_{\alpha} + B_{\alpha}^{*}(\mathbf{r},t) b_{\alpha}^{\dagger} \right]$$

$$i\hbar\partial_{t} \begin{pmatrix} A_{\alpha}(\mathbf{r},t) \\ B_{\alpha}(\mathbf{r},t) \end{pmatrix} = \begin{pmatrix} K & M \\ -M^{*} & -K^{*} \end{pmatrix} \begin{pmatrix} A_{\alpha}(\mathbf{r},t) \\ B_{\alpha}(\mathbf{r},t) \end{pmatrix}$$
$$K = -\frac{\hbar^{2}}{2m}\nabla^{2} + \frac{1}{2}m\omega_{0}^{2}r^{2} + 2U\Phi^{*}\Phi, \quad M = U\Phi^{2}$$





位相と振幅による記述 (BEC成分)

$$\Phi = \rho_0^{1/2} \exp(i\theta_0)$$

規格化したGP方程式より

$$\partial_{t}\theta_{0} = -\frac{1}{2} (\nabla \theta_{0})^{2} - \frac{1}{2} r^{2} - g\rho_{0} - \frac{1}{8\rho_{0}^{2}} (\nabla \rho_{0})^{2} + \frac{1}{4\rho_{0}} \nabla^{2} \rho_{0}$$
$$\partial_{t}\rho_{0} + \nabla \cdot (\rho_{0} \nabla \theta_{0}) = 0$$

短波長のゆらぎを無視

$$\partial_t \theta_0 \simeq -\frac{1}{2} \left(\nabla \theta_0 \right)^2 - \frac{1}{2} r^2 - g \rho_0$$
$$\partial_t \rho_0 + \nabla \cdot \left(\rho_0 \nabla \theta_0 \right) = 0$$

 $\left(\lambda \gg \xi = \frac{1}{\sqrt{g\rho_0}}\right)$

位相と振幅による記述(ゆらぎ)

ゆらぎを考える:
$$egin{cases} heta_0+ heta\
ho_0+
ho \end{cases}$$

$$\partial_t \theta \simeq -(\nabla \theta_0)(\nabla \theta) - g\rho$$
$$\partial_t \rho + \nabla \cdot (\rho \nabla \theta_0) + \nabla \cdot (\rho_0 \nabla \theta) = 0$$

$$(\partial_t + \nabla \cdot \mathbf{v}_0)(\partial_t + \mathbf{v}_0 \cdot \nabla)\boldsymbol{\theta} \simeq \nabla \cdot (c_s^2 \nabla \boldsymbol{\theta})$$

$$\mathbf{v}_0 = \nabla \boldsymbol{\theta}_0$$

m

0

$$(\partial_t + \nabla \cdot \mathbf{v}_0)(\partial_t + \mathbf{v}_0 \cdot \nabla)\theta \simeq \nabla \cdot (c_s^2 \nabla \theta)$$

曲がった時空上での場の方程式

$$\frac{1}{\sqrt{-g}}\frac{\partial}{\partial x^{\mu}}\left(\sqrt{-g}g^{\mu\nu}\frac{\partial}{\partial x^{\nu}}\psi\right) = 0$$

$$g_{\mu\nu} = c_s \begin{pmatrix} c_s^2 - v_0^2 & v_0^x & v_0^y & v_0^z \\ v_0^x & -1 & 0 & 0 \\ v_0^y & 0 & -1 & 0 \\ v_0^z & 0 & 0 & -1 \end{pmatrix}_{\mu\nu}$$

BECのmetric

$$ds^{2} = c_{s}^{2} dt^{2} - (v_{0}^{r} dt - dr)^{2} - r^{2} d\theta^{2} - r^{2} \sin^{2} \theta d\phi^{2}$$

Schwarzschild metric

$$ds^{2} = c^{2}d\tilde{t}^{2} - \left(vd\tilde{t} - dr\right)^{2} - r^{2}d\theta^{2} - r^{2}\sin^{2}\theta d\phi^{2}$$



$$\Phi_{0} + \phi = (\rho_{0} + \rho)^{1/4} \exp(i\theta_{0} + i\theta)(\rho_{0} + \rho)^{1/4}$$

$$|\theta| \ll |\theta_0|, \quad \rho \ll \rho_0$$

$$\theta = \frac{1}{2i\rho_0} \left(\Phi_0^* \phi - \Phi_0 \phi^\dagger \right)$$
$$\rho = \Phi_0^* \phi + \Phi_0 \phi^\dagger$$

内積
$$4\pi \int_0^\infty dr r^2 \left[A_\alpha^*(r) A_\beta(r) - B_\alpha^*(r) B_\beta(r) \right] = \delta_{\alpha\beta}$$

$$\phi(\mathbf{r},t) = \sum_{\alpha} \left[A_{\alpha}(\mathbf{r},t) b_{\alpha} + B_{\alpha}^{*}(\mathbf{r},t) b_{\alpha}^{\dagger} \right]$$
$$\theta = \sum_{\alpha} \left(f_{\alpha} b_{\alpha} + f_{\alpha}^{*} b_{\alpha}^{\dagger} \right)$$
$$\rho = \sum_{\alpha} \left(g_{\alpha} b_{\alpha} + g_{\alpha}^{*} b_{\alpha}^{\dagger} \right)$$

$$\begin{pmatrix} f_{\alpha}, f_{\beta} \end{pmatrix} = \delta_{\alpha\beta}$$

$$(p,q) = \frac{4\pi i}{g} \int_{0}^{\infty} dr r^{2} \Big[p^{*}(r) (\partial_{t} + \mathbf{v}_{0} \cdot \nabla) q(r) - \Big[(\partial_{t} + \mathbf{v}_{0} \cdot \nabla) p^{*}(r) \Big] q(r) \Big]$$

WKB analysis



$$(\partial_t + \nabla \cdot \mathbf{v}_0)(\partial_t + \mathbf{v}_0 \cdot \nabla)\theta \simeq \nabla \cdot (c_s^2 \nabla \theta)$$

Horizon近傍を考える
$$v = c + \kappa (r - r_H) + ...$$

 $\theta = \exp(-i\omega t + iK(r))$



WKB analysis: Hawking Temperature

$$\left|\theta_{k}\right|^{2} = \frac{1}{k^{2}} \frac{2\pi\omega}{\kappa} \frac{1}{\exp\left(\frac{2\pi\omega}{\kappa}\right) - 1}$$

Hawking temperature

$$k_{B}T_{H} = \frac{\hbar}{2\pi}\kappa \qquad \kappa = \frac{\partial}{\partial r}(v-c)\Big|_{r=r_{H}}$$

$$v = c + \kappa \left(r - r_H \right) + \dots$$















Horizon and approximate Hawking temperature







"Hawking radiation"

Dynamical evolution of "spacetime"

Gross-Pitaevskii equation

時空のdynamics

$$i\hbar\partial_t \Phi = \left(-\frac{\hbar^2}{2m}\nabla^2 + \frac{1}{2}m\omega_{ho}^2r^2\right)\Phi + U\Phi^*\Phi\Phi$$

Bogoliubov-de Gennes equation 量子

量子場のdynamics

$$i\hbar\partial_{t} \begin{pmatrix} A_{\alpha}(\mathbf{r},t) \\ B_{\alpha}(\mathbf{r},t) \end{pmatrix} = \begin{pmatrix} K-E_{0} & M \\ -M^{*} & -K^{*}-E_{0} \end{pmatrix} \begin{pmatrix} A_{\alpha}(\mathbf{r},t) \\ B_{\alpha}(\mathbf{r},t) \end{pmatrix}$$
$$K = -\frac{\hbar^{2}}{2m}\nabla^{2} + \frac{1}{2}m\omega_{0}^{2}r^{2} + 2U\Phi^{*}\Phi, \quad M = U\Phi^{2}$$

Particle creation spectrum: calculation

$$b_{\alpha}^{(1)} = \sum_{\beta} \left[A_{\alpha\beta} b_{\beta}^{(2)} + B_{\alpha\beta} b_{\beta}^{(2)\dagger} \right]$$

$$B_{\beta\alpha}^{*} = \left(f_{\beta}^{(2)*}, f_{\alpha}^{(1)}\right)$$
$$= \frac{i}{g} \int_{0}^{\infty} dr \left[f_{\beta}^{(2)}(r) \left(\partial_{t} + \mathbf{v}_{0} \cdot \nabla\right) f_{\alpha}^{(1)}(r) - \left[\left(\partial_{t} + \mathbf{v}_{0} \cdot \nabla\right) f_{\beta}^{(2)}(r)\right] f_{\alpha}^{(1)}(r)\right]$$

$$f^{(2)}_{eta}$$
時間発展してきた Φ を用いて計算

$$B_{\beta\alpha}^{*} = -4\pi \int_{0}^{\infty} dr r^{2} \left[A_{\beta}^{(2)} B_{\alpha}^{(1)} - B_{\beta}^{(2)} A_{\alpha}^{(1)} \right]$$

Time-evolution



Particle creation spectrum



Origin?



How particles are created?



Conclusion

Cold atoms are useful for investigating Hawking radiation physics!

Appendix

Black hole entropy

Hawking temperature

$$k_B T_H = \frac{\hbar c^3}{8\pi GM}$$

Thermodynamical relation

$$dU = c^2 dM = T_H dS$$

$$\bigvee$$

$$S / k_B = \frac{1}{4} \times \frac{4\pi r_g^2}{\ell_P^2}$$